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EVOLUTION OF THERMALLY PULSING ASYMPTOTIC GIANT BRANCH STARS. V. CONSTRAINING THE MASS LOSS AND LIFETIMES OF INTERMEDIATE-MASS, LOW-METALLICITY AGB STARS*

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ABSTRACT

Thermally pulsing asymptotic giant branch (TP-AGB) stars are relatively short lived (less than a few Myr), yet their cool effective temperatures, high luminosities, efficient mass loss, and dust production can dramatically affect the chemical enrichment histories and the spectral energy distributions of their host galaxies. The ability to accurately model TP-AGB stars is critical to the interpretation of the integrated light of distant galaxies, especially in redder wavelengths. We continue previous efforts to constrain the evolution and lifetimes of TP-AGB stars by modeling their underlying stellar populations. Using *Hubble Space Telescope* (*HST*) optical and near-infrared photometry taken of 12 fields of 10 nearby galaxies imaged via the Advanced Camera for Surveys Nearby Galaxy Survey Treasury and the near-infrared *HST*/SNAP follow-up campaign, we compare the model and observed TP-AGB luminosity functions as well as the ratio of TP-AGB to red giant branch stars. We confirm the best-fitting mass-loss prescription, introduced by Rosenfield et al., in which two different wind regimes are active during the TP-AGB, significantly improves models of many galaxies that show evidence of recent star formation. This study extends previous efforts to constrain TP-AGB lifetimes to metallicities ranging $-1.59 \lesssim [\text{Fe}/\text{H}] \lesssim -0.56$ and initial TP-AGB masses up to $\sim 4 M_{\odot}$, which include TP-AGB stars that undergo hot-bottom burning.

Key words: galaxies: dwarf – galaxies: stellar content – stars: AGB and post-AGB – stars: evolution – stars: general – stars: mass-loss

1. INTRODUCTION

Thermally pulsing asymptotic giant branch (TP-AGB) stars are fascinating in their structural complexity, their short lifetimes, and the dramatic impacts they have on their host galaxies. TP-AGB stars are known as important sources of dust (e.g., Gail & Sedlmayr 1985, 1999; Gehrz 1989; Boyer et al. 2012; Melbourne & Boyer 2013), they chemically enrich their ISM (e.g., Marigo 2001; Ventura et al. 2001; Karakas & Lattanzio 2007), and they are critical to the interpretation of the spectral energy distributions of high-redshift galaxies (e.g., Bertoldi et al. 2003; Conroy et al. 2009; Valiante et al. 2009; Melbourne et al. 2012).

A full comprehension of the TP-AGB phase and its effects on its environment is critical to understanding galaxies' near- and far-infrared light output. However, TP-AGB stellar evolution is complex, with double shell fusion, large convective events (third dredge up), fusion at the base of the convective envelope (hot-bottom burning, HBB, for initial TP-AGB masses $\gtrsim 3 M_{\odot}$), and pulsations that help to create dust and eventually drive the mass-loss rate. As a consequence, even its most basic evolutionary property—the TP-AGB lifetime—differs greatly from set to set of evolutionary calculations (which can be appreciated in Figure 3 of Girardi & Marigo 2007a).

The strongest observational constraints on TP-AGB lifetimes first came from resolved stellar populations in the Magellanic Clouds (MCs; Frogel et al. 1990; van Loon et al. 2005; Girardi & Marigo 2007a) and then nearby galaxies (Girardi & Marigo 2007a; Rosenfield et al. 2014). The MC clusters are natural laboratories for studying the TP-AGB phase as there are many clusters with well defined main-sequence turn-offs (e.g., Gallagher et al. 1996; Geha et al. 1998; Holtzman et al. 1999; Olsen 1999), known distances (e.g., Pietrzyński et al. 2013), deep optical and IR observations (Blum et al. 2006; Meixner et al. 2006), and known metallicities (e.g., Bica et al. 1998; Da Costa & Hatzidimitriou 1998; Dirsch et al. 2000; Hill et al. 2000; Geisler et al. 2003; Cole et al. 2005; Ferraro et al. 2006). However, individual MC clusters lack a statistically significant population of TP-AGB stars, which has led some researchers to combine information from MC clusters based on age and derive TP-AGB lifetime constraints from ensemble populations (e.g., Frogel et al. 1990; Girardi & Marigo 2007a; Noël et al. 2013). This method is now seen as problematic given many MC clusters happen to be of an age that temporarily boosts the number counts of TP-AGB stars, falsely signifying a longer lifetime. However, TP-AGB “boosting” does not eliminate MC clusters or fields from constraining the TP-AGB; it requires an extra calibration step to determine the amount of boosting when modeling stellar populations of ~ 1.6 Gyr (see Girardi et al. 2013).

Outside the MCs, star clusters and fields containing significant numbers of TP-AGB stars are rarely observed down to the old main sequence turn-off; however, modern color-magnitude diagram (CMD) fitting techniques allow their star

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Table 1
Basic Parameters of the Galaxy Sample

Galaxy	Optical Filters	A_V	$(m - M)_0$	D (Mpc)	$\frac{SF}{SF_{TOT}}$ <1 Gyr	$\langle[Fe/H]\rangle$ 1–3 Gyr	$\frac{SF}{SF_{TOT}}$ 1–3 Gyr	$\langle[Fe/H]\rangle$ 1–3 Gyr
NGC 300-WIDE1	<i>F606W, F814W</i>	0.01	26.50	2.00	$0.08^{+0.01}_{-0.01}$	$-0.19^{+0.12}_{-0.12}$	$0.29^{+0.02}_{-0.04}$	$-0.21^{+0.12}_{-0.12}$
UGC 8508	<i>F475W, F814W</i>	0.05	27.05	2.57	$0.07^{+0.01}_{-0.01}$	$-1.19^{+0.12}_{-0.12}$	$0.09^{+0.03}_{-0.02}$	$-1.22^{+0.12}_{-0.12}$
NGC 4163	<i>F606W, F814W</i>	0.06	27.29	2.87	$0.04^{+0.00}_{-0.00}$	$-0.99^{+0.12}_{-0.12}$	$0.04^{+0.00}_{-0.00}$	$-1.00^{+0.12}_{-0.12}$
NGC 2403-DEEP	<i>F606W, F814W*</i>	0.12	27.50	3.16	$0.07^{+0.01}_{-0.01}$	$-0.49^{+0.12}_{-0.12}$	$0.09^{+0.02}_{-0.04}$	$-0.53^{+0.12}_{-0.12}$
NGC 2403-HALO-6	<i>F606W, F814W</i>	0.12	27.50	3.16	$0.10^{+0.02}_{-0.03}$	$-0.78^{+0.13}_{-0.12}$	$0.30^{+0.07}_{-0.08}$	$-0.81^{+0.16}_{-0.13}$
UGC 4459	<i>F555W, F814W</i>	0.12	27.79	3.61	$0.08^{+0.01}_{-0.01}$	$-0.99^{+0.12}_{-0.35}$	$0.11^{+0.03}_{-0.01}$	$-1.02^{+0.12}_{-0.12}$
NGC 3741	<i>F475W, F814W</i>	0.08	27.55	3.24	$0.07^{+0.01}_{-0.01}$	$-1.19^{+0.12}_{-0.12}$	$0.07^{+0.02}_{-0.02}$	$-1.20^{+0.12}_{-0.12}$
UGC 5139	<i>F555W, F814W</i>	0.15	27.95	3.89	$0.07^{+0.01}_{-0.01}$	$-0.80^{+0.12}_{-0.12}$	$0.14^{+0.02}_{-0.03}$	$-0.89^{+0.12}_{-0.12}$
UGC 4305-1	<i>F555W, F814W</i>	0.01	27.65	3.39	$0.08^{+0.01}_{-0.01}$	$-1.18^{+0.12}_{-0.12}$	$0.06^{+0.01}_{-0.01}$	$-1.18^{+0.12}_{-0.12}$
UGC 4305-2	<i>F555W, F814W</i>	0.10	27.70	3.47	$0.06^{+0.00}_{-0.00}$	$-1.08^{+0.12}_{-0.12}$	$0.05^{+0.01}_{-0.01}$	$-1.09^{+0.12}_{-0.12}$
KDG 73	<i>F475W, F814W</i>	0.06	28.03	4.04	$0.04^{+0.01}_{-0.01}$	$-1.20^{+0.12}_{-0.12}$	$0.07^{+0.04}_{-0.04}$	$-1.28^{+0.12}_{-0.12}$
UGCA 292	<i>F606W, F814W</i>	0.00	28.64	5.35	$0.45^{+0.06}_{-0.05}$	$-1.20^{+0.12}_{-0.12}$	$0.55^{+0.07}_{-0.23}$	$-1.27^{+0.14}_{-0.12}$

Note. Properties of the galaxy sample. The first column lists the target consistent with the footprints named in Dalcanton et al. (2009, 2012) and available on the Mikulski Archive for Space Telescopes (MAST), followed by the optical filters used in the *MATCH* CMD fitting (*WFPC2 data). Columns 3 and 4 are the extinction, A_V , and distance modulus, $(m - M)_0$, (calculated from TRGB fitting; Dalcanton et al. 2009). Columns 5 and 7 show the fraction of total stars formed between 3.5 Myr and 1 Gyr and between 1 Gyr and 3 Gyr, respectively. Columns 6 and 8 show the average metallicity in the same time bins. The youngest PARSEC models currently available in *MATCH* are 3.5 Myr.

formation histories (SFHs) to be well constrained up to distances of about 4 Mpc for high-resolution *Hubble Space Telescope* (*HST*) imaging (e.g., Dolphin 2002; Weisz et al. 2011). Such imaging was obtained via the *HST*/Advanced Camera for Surveys (ACS) Nearby Galaxy Survey Treasury (ANGST; Dalcanton et al. 2009), which surveyed a volume-limited sample of ~ 70 nearby galaxies from 1 to 4 Mpc in optical ACS and WFPC2 filters. ANGST galaxies show a wide diversity of SFHs (see, e.g., Weisz et al. 2011; Williams et al. 2013) and present a multitude of candidate TP-AGB stars at a range of metallicities.

Girardi et al. (2010, hereafter G10) used ANGST galaxies that showed little-to-no recent SF to constrain lifetimes of low-mass, low-metallicity TP-AGB stars by suggesting a phase of mass loss before the onset of dust-driven winds (pre-dust). Rosenfield et al. (2014, hereafter R14) extended the G10 galaxy sample by incorporating near-infrared imaging from a follow-up *HST*/SNAP campaign (AGB-SNAP; Dalcanton et al. 2012) and presented a mass-loss prescription that best represents the optical and near-infrared (NIR) luminosity functions (LFs) of seven galaxies which also showed little-to-no recent SF. Both studies used the most likely SF history of each galaxy to produce stellar population synthesis models, which were used to compare the mean predicted TP-AGB lifetimes to the data using the ratio of TP-AGB stars to red giant branch (RGB) stars (N_{TP-AGB}/N_{RGB}). R14 further performed a calibration of the TP-AGB mass-loss prescriptions by comparing the shapes of the predicted and observed luminosity functions. However, due to the lack of recent SF in the selected galaxies in both programs, the TP-AGB stars studied were restricted to initial masses ~ 0.8 – $2.5 M_\odot$ and metallicities $-1.54 < [Fe/H] < -0.86$. In this paper, we extend the analysis of G10 and R14 to include 10 more galaxies of the AGB-SNAP sample that show evidence of recent SF and probe a metallicity range of $-1.59 \lesssim [Fe/H] \lesssim -0.56$ and an initial mass range of ~ 0.8 – $4 M_\odot$ of stars that may eventually populate the TP-AGB, importantly including TP-AGB stars that may undergo HBB. The paper is outlined as follows. In Section 2 we briefly

describe the observations, photometry, and data reduction used in this project and present the re-derivation of SFHs and a new method of identifying TP-AGB stars in galaxies with recent SF. In Section 3 we review the R14 TP-AGB model that we test against the observations. We model the data using population synthesis described in Section 4. We compare the population synthesis models to the observations in Section 5 and conclude in Section 6.

2. DATA

2.1. Galaxy Sample

The galaxy sample is derived from Melbourne et al. (2012, hereafter M12), who developed 23 resolved stellar source catalogs of 22 galaxies that matched the AGB-SNAP (*HST*/WFC3) NIR data with the ANGST (*HST*/ACS or WFPC2) optical data. We selected 12 fields of 10 galaxies from M12 that show evidence of recent SF (and were not one of the 7 galaxies studied in R14).

Detailed information on the galaxies in this study can be found in Dalcanton et al. (2009, 2012). Briefly, they are all dwarf spheroidal or dwarf irregular galaxies with the exceptions of NGC 2403 and NGC 300. In the case of NGC 2403, we study a deep (WFPC2) field in the disk of the galaxy as well as a shallower halo (ACS) field (see Figure 2 in Williams et al. 2013) and we use two fields of UGC 4305 (HoII) with similar SFHs. Some parameters of interest of these galaxies are presented in Table 1 below and typical optical–NIR CMDs from the M12 matched catalogs are presented in Figure 3.

We make use of the ANGST optical data, AGB-SNAP NIR data, the M12 matched catalogs, and the corresponding artificial star tests (ASTs) in this study. Following the methodology of R14, we re-calculate the SFHs using the same stellar evolutionary models that we use in our population synthesis code. In order to have the best leverage on the SFH, we use the optical ANGST data which reach several magnitudes below the tip of the RGB (TRGB), usually beyond the red clump. The AGB-SNAP NIR data (as part of the M12 matched catalogs) have a more complete census of TP-AGB

stars than the optical ANGST data so we use them to compare against our simulated stellar populations. Finally, artificial star tests are employed in the SFH derivation, applied to make completeness corrections to the observed LFs (see Figure 6), and used to simulate statistically similar uncertainties and crowding in the modeled stellar populations. We now briefly discuss the reduction and photometry of each catalog and the construction of the M12 four-filter catalogs.

2.2. Reduction and Photometry

The reduction and photometry of the ANGST and AGB-SNAP data as well as the AST strategy are fully described in Dalcanton et al. (2009, 2012). In short, all data were photometered and reduced as part of the University of Washington photometry pipeline, which uses the DOLPHOT package⁸ (Dolphin 2000) and the ACS, WFC3, or WFC2 module. Briefly we use the “good” (gst) star catalogs that apply conservative limits of the DOLPHOT flags: sharpness, $(\text{sharp}_1 + \text{sharp}_2)^2 \leq 0.075$; crowding, $\text{crowd}_1 + \text{crowd}_2 \leq 0.1$; signal-to-noise, $S/N > 4$; and $\text{flag} < 8$ (i.e., not extended, elongated, highly saturated, or significantly cut off by CCD chip edges).

We refer the reader to M12 for complete details on the construction of the four-filter NIR–optical matched catalog and only outline the methodology here. To create the matched catalogs, M12 transformed the NIR data set coordinate system to that of the optical using 150 bright, spatially dispersed red stars in the DOLPHOT output catalogs (which were also input to the ANGST gst catalogs). The transformation was calculated iteratively (with the MATCH routine developed by Michael Richmond); first they found a linear fit between the coordinate systems and then used the fit as the starting point for a quadratic solution. M12 reports that 90% of the stars in the NIR catalogs are matched to within $0''.07$ of a star in the optical catalogs.

2.3. Star Formation Histories

Star formation histories were calculated using the `calcsfh` routine in the MATCH package (Dolphin 2002). MATCH is a CMD-fitting package that finds the most likely SFH and metallicity evolution of an input CMD given an IMF, binary fraction, uncertainties and completeness of the photometry, and stellar evolution models. We adopted the Kroupa (2001) IMF and a binary fraction of 0.35 with a flat distribution of secondary masses. The faint limits of the photometry were set to the 50% completeness limit determined by the artificial star tests, listed in Table 5 of Dalcanton et al. (2009). For the stellar evolution models, we use a subset of the Padova TRIeste Stellar Evolution Code (PARSEC V1.2S) stellar evolution tracks (Bressan et al. 2012, 2013) (discussed below). Finally, we adopted a binning scheme in color–magnitude space of 0.05 and 0.1, respectively, and metallicity bins (as $\log Z$) of 0.15 dex. We further allowed MATCH to find the best-fitting distance and extinction. However, in the cases where `calcsfh` did not converge to within 0.05 dex of the distance modulus or A_V value derived from TRGB fitting (Dalcanton et al. 2009), we adopted their values to ensure the model CMDs agreed to within a pixel of the binned data.

PARSEC was also used in R14 to derive the SFHs since it calculates the first thermal pulse used in the COLIBRI TP-AGB models (Marigo et al. 2013). Since R14, PARSEC has undergone significant updates, the largest of which extended the upper mass limit of the models from $12 M_\odot$ to $350 M_\odot$ in certain metallicities (Tang et al. 2014) and presented new low-mass stellar evolution models (Chen et al. 2015). To import PARSEC for use in MATCH, we redefined PARSEC equivalent stellar evolutionary points (i.e., interpolation points along a stellar evolution track) such that each track has a uniform number of points and equidistant spacing between them. As MATCH requires a complete and regular grid in mass and metallicity, this resulted in PARSEC V1.2S in MATCH having a different parameter range than the full PARSEC V1.2S. In MATCH, the PARSEC V1.2S grid covers a mass range of $0.1 M_\odot \leq M \leq 120 M_\odot$ which translates to an age range of $3.55 \text{ Myr} \leq \text{Age} \leq 15.8 \text{ Gyr}$ and a metallicity range from $-2.18 \leq [\text{Fe}/\text{H}] \leq 0.42$ (the full grid of PARSEC V1.2S tracks reaches $[\text{Fe}/\text{H}] = 0.6$ and younger ages). Finally, the BC transformations of the stellar models occur within MATCH and use those specified in Girardi et al. (2008), which largely rely on the Castelli & Kurucz (2003) model atmospheres.

With the above parameter search space, we re-fit the program galaxies using the optical data from ANGST and ASTs described in Dalcanton et al. (2009) and Weisz et al. (2011). Following the method of G10 and R14, we excluded stars above the TRGB from the fitting (using the TRGB magnitudes calculated in Dalcanton et al. 2012). While G10 found very small or no differences in measured SFH including or not including TP-AGB stars, we excluded them from the data as we had not yet ported model TP-AGB stars into MATCH and did not want to introduce bias due to the fitting algorithm finding high S/N data where there were no model points.

The last four columns of Table 1 list the relative amounts of recent SF and corresponding average metallicities for the galaxy sample. Cumulative SFHs are shown in Figure 1 with random uncertainties calculated with the hybrid Monte Carlo algorithm implemented as the `hybridMC` routine in MATCH and described in Dolphin (2013). All SFH fits are consistent to less than 0.01% level except for the KDG 73 fit which is consistent to 0.12%.

The program targets already have measured SFH using MATCH (Weisz et al. 2011; Dalcanton et al. 2012; Melbourne et al. 2012) with the exceptions of the fields of NGC 2403 and NGC 300 which were separately studied in detail in Williams et al. (2013) and Gogarten et al. (2010). All of the above previous SFH determinations were calculated with slightly different stellar models (Padua; Marigo et al. 2008; Girardi et al. 2010). Therefore it is worth comparing the published results to our new SFH derivations.

To compare the SFHs calculated with either the PARSEC or the Padua models, we overplot the systematic uncertainties from Weisz et al. (2011) in Figure 1. The systematic uncertainties were estimated by deriving the SFH of a synthetic CMD with a different set of stellar evolution models than were inputted and applying the uncertainties to the galaxy’s SFH solution. For example, Weisz et al. (2011) produce a synthetic CMD using BaSTI models (Pietrinferni et al. 2004) and then derive the SFH of that synthetic CMD using Padua models. The offsets between the input SFH that created the synthetic CMD and the output SFH are related to systematic uncertainties in the stellar models (for full details of this method, see

⁸ <http://purcell.as.arizona.edu/dolphot/>

Figure 1. MATCH solutions for the cumulative SF rate as a function of galaxy age for the galaxies in the sample (black). Dark gray regions denote random uncertainties calculated by MATCH’s hybrid Monte Carlo algorithm (see Dolphin 2013), light gray regions are the systematic uncertainties calculated in Weisz et al. (2011, where available). SFH solutions using PARSEC agree to within systematic uncertainties except for NGC 4163 and UGC 5139, where PARSEC predicts higher SF in the oldest, most uncertain age bins (see text).

Figure 2. Different SFHs due to slight differences in the underlying stellar evolution models do not significantly affect the TP-AGB mass distribution. Top: MATCH-derived cumulative SF using optical data of UGC 5139 and stellar evolution models from Padua (black; derived in Weisz et al. 2011) and PARSEC (blue; see Figure 1). Significant inconsistencies (outside the bounds of the gray systematic uncertainties derived in Weisz et al. 2011) are present at the oldest ages. Bottom: derived mass distributions using each SFH in the TRILEGAL code (see Section 4). The Padua mass distribution is shown in gray with its TP-AGB in black, and the PARSEC mass distribution is shown in light blue with its TP-AGB in blue. Despite the major inconsistencies in the inferred mass distributions, they do not significantly affect the shape of the TP-AGB mass distribution.

Dolphin 2012). We include systematic uncertainties as a visual guide to the range of possible SFHs, though we do not expect differences from Padua to PARSEC to be nearly as great as are seen between BaSTI-Padua. Furthermore, it should be noted that for these data, using one stellar evolution model instead of another will not produce totally incomparable SFHs (assuming the age, mass, and metallicity ranges of the models are comparable; compare Dolphin 2012).

The cumulative SFHs are within uncertainties with the exceptions of NGC 4163 and UGC 5139, where Weisz et al.

(2011) calculate relatively suppressed SF at ages older than ~ 10 Gyr compared to our solutions. The discrepancies between PARSEC and Padua SFH solutions for these two galaxies clearly go beyond the systematic uncertainties calculated by Weisz et al. (2011) at the oldest ages (and therefore lowest masses, where some of the largest upgrades to the Padua models occurred; Chen et al. 2014). However, as we now show, the apparently large deviations at old ages do not affect the relative mass distribution of the TP-AGB and can be neglected for this study.

Figure 3. Selecting TP-AGB and RGB stars. Top left: optical–NIR absolute magnitude CMD of UGC 5139 with the results of double Gaussian fitting method discussed in the text (thin black lines) to test the separation line between RHeB and TP-AGB stars (thick black line). Next to each double Gaussian are two sets of numbers in each magnitude bin; on the left is the fraction of stars from the “RHeB” Gaussian that are expected in the TP-AGB region and on the right is the fraction of stars in the “TP-AGB” Gaussian expected in the RHeB region. The TRGB magnitude is marked by a dashed horizontal line. Top right: optical–NIR TRGB color vs. M_{Br} . Galaxies brighter than $M_B = -17$ have a wider color selection box for the RGB than those that are fainter. Bottom: example results of the TP-AGB and RGB identification for NGC 3741 (left) and NGC 300 (right). RGB stars selected to scale the population synthesis models are within the red box and stars considered TP-AGB (above the TRGB and redward of the RHeB–TP-AGB line) are shown in blue, with the RHeB–TP-AGB line shown in black.

The top panel of Figure 2 shows the cumulative SFH of UGC 5139 reproduced from Figure 1 which used PARSEC models as input. Overplotted is the cumulative SFH from Weisz et al. (2011), who used Padua models as the stellar evolution models in MATCH. In the bottom panel, the inferred mass distributions calculated using TRILEGAL and the R14 TP-AGB models (see Section 4) are shown for each SFH in $0.1 M_\odot$ bins. Despite the clear offsets in the ancient SFR and mass distribution of the inferred stellar population, the relative TP-AGB mass distribution hardly changes. Furthermore, a Kolmogorov–Smirnov test (Kolmogoroff 1933; Smirnov 1948)⁹ of the TP-AGB mass distribution resulted in D_{KS} statistic of 0.036 and a p value of 0.9995, therefore we can reject the null hypothesis, which indicates that the SFHs from previous studies using MATCH/Padua give statistically indistinguishable TP-AGB

populations compared to those produced with MATCH/PARSEC V1.2S.

2.4. Identifying TP-AGB Stars

The region of TP-AGB stars in an optical–NIR CMD are easy to identify by eye, as TP-AGB stars are the brightest and reddest when present (see Figure 3). However, the boundary in CMD space near the TRGB as well as the blue edge of the TP-AGB can contain RGB stars and red core helium-burning (RHeB) stars, respectively. Confusing RGB and RHeB stars with TP-AGB stars will erroneously increase the apparent number of TP-AGB stars in the observations which will be interpreted as longer TP-AGB lifetimes.

We exclude RGB and RHeB stars and thereby isolate TP-AGB stars in the following way. First, to exclude RGB counts in the N_{TP-AGB}/N_{RGB} ratio we continue the method of R14 and

⁹ Implemented with `ks_2samp` in `scipy` (Jones et al. 2001).

Table 2
 $N_{\text{TP-AGB}}/N_{\text{RGB}}$ Comparisons

Target	Data	R14	Data	R14	Fractional Difference
	$N_{\text{TP-AGB}}$	$N_{\text{TP-AGB}}$	$\frac{N_{\text{TP-AGB}}}{N_{\text{RGB}}}$	$\frac{N_{\text{TP-AGB}}}{N_{\text{RGB}}}$	
NGC 300-WIDE1	422 ± 21	462 ± 21	0.511 ± 0.043	0.560 ± 0.046	−0.097 ± 0.016
UGC 8508	246 ± 16	167 ± 13	0.368 ± 0.038	0.249 ± 0.029	0.325 ± 0.071
NGC 4163	575 ± 24	693 ± 26	0.363 ± 0.024	0.437 ± 0.028	−0.203 ± 0.026
NGC 2403-DEEP	181 ± 13	150 ± 12	0.614 ± 0.081	0.508 ± 0.071	0.171 ± 0.047
NGC 2403-HALO-6	136 ± 12	93 ± 10	0.604 ± 0.092	0.414 ± 0.071	0.315 ± 0.102
NGC 3741	241 ± 16	213 ± 15	0.429 ± 0.046	0.377 ± 0.042	0.120 ± 0.026
UGC 4305-1	549 ± 23	340 ± 18	0.558 ± 0.042	0.346 ± 0.030	0.380 ± 0.061
UGC 4305-2	612 ± 25	602 ± 25	0.463 ± 0.031	0.454 ± 0.031	0.019 ± 0.003
UGC 4459	247 ± 16	327 ± 18	0.351 ± 0.036	0.465 ± 0.043	−0.324 ± 0.063
UGC 5139	475 ± 22	210 ± 15	0.401 ± 0.030	0.178 ± 0.017	0.557 ± 0.096
KDG 73	49 ± 7	30 ± 5	0.282 ± 0.062	0.171 ± 0.044	0.392 ± 0.187
UGC A292	71 ± 8	115 ± 11	0.477 ± 0.096	0.779 ± 0.137	−0.635 ± 0.239
Mean	317 ± 62	283 ± 58	0.452 ± 0.197	0.412 ± 0.200	0.089 ± 0.082

Note. The data and R14 $N_{\text{TP-AGB}}/N_{\text{RGB}}$. Both are calculated from the NIR matched catalogs (see Section 2.4). The uncertainties reported on the mean are the quadrature of the listed uncertainties of each field and the mean fractional difference is calculated from the mean ratio values, not for example, the mean of the individual fractional differences.

exclude a ± 0.2 m_{F160W} band around the TRGB. To minimize the contribution of RHeB stars counted as TP-AGB we used the M12 matched catalogs to find an approximate separation line in CMD space between TP-AGB stars and the RHeB sequence.

We first estimated the RHeB TP-AGB separation line for all of the galaxies stacked together and then tested its adequacy on each individual galaxy to refine the slope and intercept.

To stack the galaxies, we shifted all the optical–NIR data to absolute magnitudes using the distance moduli and extinction values (A_V) calculated in Dalcanton et al. (2009, 2012) and corrected the A_V values to the extinction values of the filters in question using coefficients derived from Cardelli et al. (1989) with an extinction curve with $R_V = 3.1$ applied to a G2V star ($A_{\text{F160W}}/A_V = 0.6056$ and $A_{\text{F814W}}/A_V = 0.2044$, see Girardi et al. 2008).

We tested the initial separation line by fitting a double Gaussian distribution as a function of absolute $F160W$ magnitude for each field within the color limits $F814W - F160W = [1, 3]$ using Levenberg–Marquardt least-squares minimization.¹⁰ We then calculated the expected numbers of stars from each Gaussian on each side of the separating line for each magnitude bin within the color limits.

We set the magnitude bins with the binning algorithm from Knuth (2006) implemented in the astroML package (Vanderplas et al. 2012). The resulting slope works well in many of the most populated systems, however, it fails in the regime where a single Gaussian can provide a better fit than a double, a scenario usually due to an underpopulated RHeB or a large color dispersion of the TP-AGB (a sign of an intrinsic metallicity spread).

We also found that the slope calculated worked well for all galaxies, but we needed to further take into account the intrinsic offsets in the color of the RGB from galaxy to galaxy due to differences in metallicity. Therefore, we adjusted the separation line to include a TRGB color term. The resulting

line can be reproduced following the equation:

$$M_{\text{F160W}} = -7.0 - 9.03[(M_{\text{F814W}} - M_{\text{F160W}}) - (M_{\text{F814W}} - M_{\text{F160W}})_{\text{TRGB}}].$$

An example of this procedure is shown in Figure 3. With this method, we find for each field that RHeB stars do not contribute to the TP-AGB by more than 5% and conversely, fewer than 6% of TP-AGB stars are found in the RHeB region (mean of 3%). We verified this result using TRILEGAL in stellar population synthesis mode and added photometric uncertainties from the ASTs (the same method as discussed below in Section 3). Since we are able to track the individual stellar phases in the TRILEGAL simulated stellar catalog, we find the mean contamination by RHeB stars to be 3% and the fraction of AGB stars (both TP-AGB and early-AGB) that is “lost” to the RHeB region is never more than 4%.

Stars redward of the separation line and brighter than the TRGB are considered TP-AGB stars in the following analysis. The number of TP-AGB stars identified in each galaxy field is listed in the first two columns of Table 2.

3. TP-AGB EVOLUTION MODELS

The TP-AGB models tested here are essentially the same as in R14; the full set spans initial TP-AGB masses from $0.50 < M < 6.0 M_{\odot}$ and a metallicity range of $0.0005 < Z < 0.06$ ($-1.48 < [\text{Fe}/\text{H}] < 0.60$). Obviously, we are limited by the metallicities in the data set and only test the lower end of the calculated metallicity range. Briefly, the TP-AGB models incorporate the following main characteristics:

1. The conditions at the first thermal pulse on the AGB are extracted from PARSEC V1.2S stellar evolution models.
2. The TP evolution is computed with the COLIBRI code which is extensively described in Marigo et al. (2013). The code includes several innovative features, for instance, detailed nucleosynthesis calculations both during the thermal pulses and the interpulse phase, and on-the-fly computation of atomic and molecular abundances and their Rosseland mean opacities (Marigo & Aringer 2009) as the envelope composition changes along the evolution due to mixing events and/or HBB. A few key

¹⁰ The python implementation is based on MINPACK-1, <http://cars9.uchicago.edu/software/python/mpfit.html>.

Figure 4. Evolution of $4 M_{\odot}$ TP-AGB stars at different metallicities, $Z = 0.001$ (left) and $Z = 0.008$ (right). From top to bottom the panels show effective temperature, temperature at the base of the convective envelope, luminosity, C/O ratio (with $C/O = 1$ as a dotted horizontal line), star mass, and mass-loss rate. HBB is stronger at lower metallicities due to higher temperatures at the base of the convective zone. HBB does not occur in the higher metallicity TP-AGB star, which exhibits a constant increase in the C/O ratio and does not show the bell-shaped increase in stellar luminosity exhibited by the lower metallicity star.

processes are parameterized, such as the efficiency of third dredge-up episodes and the efficiency of mass loss.

3. A complete nuclear network that includes the proton-proton chains, the CNO, NeNa, MgAl cycles, and several (α, γ) and (α, n) reactions is coupled to a time-dependent diffusive description of convection to follow the nucleosynthesis in the HBB phases of massive TP-AGB stars. The concentrations of the most abundant elements in the pulse-driven convection zone (e.g., ${}^4\text{He}$, ${}^{12}\text{C}$, ${}^{16}\text{O}$, ${}^{22}\text{Ne}$, ${}^{24,25}\text{Mg}$) are also predicted solving a proper nuclear network.
4. For the mass loss, we adopt a two-phase formalism that includes the pre-dust phase ($\dot{M}_{\text{pd}}$) using a modified version of the Schröder & Cuntz (2005) formula, inspired

by the theoretical work of Cranmer & Saar (2011) on Alfvén waves and atmospheric turbulence, and the pulsation-assisted dust-driven phase ($\dot{M}_{\text{dd}}$) following a formula similar to Bedijn (1988) but with parameters calibrated with a sample of Galactic Miras. The super wind mass loss ($\dot{M}_{\text{sw}}$) is assumed not to exceed a maximum limit determined according to Vassiliadis & Wood (1993). The largest mass loss between these two prescriptions is adopted as the reported mass loss.

Low-mass TP-AGB stars with the above mass loss prescriptions are fully discussed in R14, so we remark on the higher-mass TP-AGB stars. The evolution of two TP-AGB stars with initial masses $M = 4.0 M_{\odot}$ and metallicities

Figure 5. Duration of (left) and mass lost in (right) each mass-loss phase for TP-AGB stars (pre-dust, $\dot{M}_{pd}$, in blue; pulsation-assisted dust-driven, $\dot{M}_{dd}$, in yellow; and super wind, $\dot{M}_{sw}$, in red). Shown are initial masses between 3 and 5 $M_{\odot}$ at different metallicities, $Z = 0.001, 0.008$ (top, bottom). Pre-dust mass loss is less efficient than the other phases, however, it lasts for the largest proportion of intermediate-mass TP-AGB lifetimes. This trend makes the mass lost to pre-dust winds comparable to the mass lost during the dust-driven winds for lower masses and higher metallicities. For higher masses and lower metallicities, the mass lost in pre-dust winds is comparable to that lost in the super wind phase.

$Z = 0.001$ and $Z = 0.008$ ($[\text{Fe}/\text{H}] = -1.18, -0.28$) are shown in Figure 4. Each vertical panel shows a different feature as a function of TP-AGB age (motivated by Vassiliadis & Wood 1993) from top to bottom: the effective temperature, temperature at the base of the convective envelope, luminosity, surface carbon-to-oxygen ratio (C/O; with a horizontal dashed line at C/O = 1), stellar mass, and mass-loss rate. Vertical dotted lines separate the mass-loss regimes discussed above. The over-luminous structure of the lower metallicity track (second from the top panel, left side of Figure 4) is an effect of strong HBB which increases with decreasing metallicity (for a given stellar mass) due to higher temperatures at the base of the convective envelope.

Pre-dust mass loss is important across the full span of TP-AGB mass and metallicities, though decreasingly so with increasing metallicity. Figure 5 shows TP-AGB stars with initial masses between 3 and 5 $M_{\odot}$ (calculated at mass steps of $\Delta M = 0.2 M_{\odot}$) at two low metallicities ($Z = 0.001, 0.008$ or $[\text{Fe}/\text{H}] = -1.18, -0.28$). The left panels show the fraction of time a TP-AGB star of a given initial mass spends in the pre-dust, pulsation-assisted dust-driven, and super wind phases. The right panels show the amount of mass lost in each of those phases as a function of initial TP-AGB mass. Generally, the amount of time spent in the pre-dust phase increases with decreasing initial mass while the amount of mass lost from the pre-dust phase is roughly constant with mass (with the exception of the high mass $Z = 0.008$ TP-AGB stars, which shows a slight decrease in mass lost with increasing initial TP-AGB mass). With increasing metallicity, we can expect an overall decrease in the duration of the pre-dust phase as it becomes easier to make dust at lower effective temperatures, so the pulsation-assisted dust-driven wind phase will take over sooner. It is interesting to note the amount of mass lost in the low metallicity pre-dust phase is comparable to the mass lost in the super wind phase and at increasing metallicity, the pre-dust winds are responsible for a comparable amount of mass lost to the pulsation-assisted dust-driven winds. However, the trends

described above are derived with a simplified prescription for AGB mass loss. We postpone to future works the application of state-of-the-art dynamical models of AGB atmospheres (Höfner 2008; Bladh et al. 2013; Eriksson et al. 2014) for more accurate and physically grounded predictions of the relative effects of the pre-dust wind and the dust-driven wind as a function of mass and metallicity.

4. MODELING THE DATA

To model the data with TP-AGB stars, we use the population synthesis code TRILEGAL. Initially designed to simulate foreground stars in the Milky Way (Girardi et al. 2005), it is also able to simulate galaxies of known distance, mean extinction, and SFH. TRILEGAL uses as input the PARSEC V1.2S models. Although such models are computed at constant mass, mass loss during the RGB is taken into account internally by TRILEGAL at the phase of isochrone construction. We adopt RGB mass loss described by Reimers (1975) with a multiplicative factor $\eta = 0.2$ (Miglio et al. 2012), which typically produces just a modest (of the order of $\sim 0.2 M_{\odot}$ at most) change in the initial mass of the RGB and TP-AGB stars considered in this work.¹¹

Since Girardi & Marigo (2007b), TRILEGAL incorporates TP-AGB stars taking into account a series of their most basic evolutionary and spectral properties. For the present paper, suffice it to mention that the photometric properties are simulated taking fully into account the presence of L and T_{eff} variations due to the thermal pulse cycles and the changes both in the photospheric spectra and in dust-reprocessing properties taking place as the stars transit between O-rich and C-rich configurations (with either $\text{C}/\text{O} < 1$ or > 1 , where C/O is the surface carbon-to-oxygen ratio). A detailed description of how these processes are implemented can be found in Marigo et al. (2008) and Aringer et al. (2009). In particular, here we assume

¹¹ As mentioned in passing in R14, the R14 pre-dust mass-loss prescription may be too efficient to extend to the lower luminosities of RGB stars.

Figure 6. Comparisons between the LF predicted by our models using the **R14** TP-AGB tracks and the **M12** data. The data are marked by the dark gray line with Poisson uncertainties, with a separate red curve for the TP-AGB stars only. The models are derived from at least 100 **TRILEGAL** simulations sampling the random uncertainties in the SFH: they are marked by the gray shaded region denoting the 67% confidence level interval. The contribution from TP-AGB stars is shown in blue with the same confidence interval shaded. The data have been completeness corrected with AGB-SNAP ASTs (the faint fall off of the optical LFs is due to the constraint of both NIR and optical detections in the **M12** data). Over-plotted are the TRGB with a 0.2 mag dispersion as well as the 1 mag limit below the TRGB and the magnitude limits of the RGB selection box (vertical lines). The top of each panel shows individual entries in Table 2 to compare the data and model. Optical and near-infrared LFs are presented in the left and right panels, respectively.

the dust properties of the 60% aluminum oxide + 40% silicate mixture for O-rich stars, and 85% amorphous carbon + 15% silicon carbide mixture for C-rich stars (Groenewegen 2006). In addition, the present code presents some relatively minor updates that are discussed in P. Marigo et al. (2016, in preparation). We ignore the photometric changes due to the long-period variability since its effect is relatively small in the NIR and since at low metallicities a major fraction of the TP-AGB stars pulsate only in low-amplitude overtone modes. Therefore, pulsation is not expected to change the bulk numbers or the LFs of the TP-AGB stars we predict for the present sample of galaxies.

For each galaxy, we follow the same method as **R14**: briefly, we sample the random uncertainties in the SFH to produce

~100 **TRILEGAL** stellar populations of an arbitrarily large total mass and apply photometric uncertainty and completeness information from the data by applying ASTs to each simulation. To account for the mass difference between our simulation and the observation, we scale the simulation to match the number of a subset of RGB stars in **M12**'s four-filter matched catalogs.

To select the RGB stars to scale the simulation, unlike **R14**, we cannot uniformly apply the same $F110W - F160W$ color cut to exclude possible main sequence contaminants as the galaxies in our program also have a large number of RHeB stars. Instead, the RGB subset is defined by a box in optical-NIR CMD space with magnitude limits set at the $F160W$ TRGB and 1 mag below. The blue $F814W - F160W$ limit of

Figure 6. (Continued.)

the RGB box is set at 0.2 mag bluer than the $F814W - F160W$ TRGB color, which nicely maps to the integrated B-band magnitude, M_{B_T} (Karachentsev et al. 2004) as the slope of the RGB is primarily set by the metallicity (e.g., see Section 6 of Dalcanton et al. 2012). Finally, the width of the RGB box is set to 0.4 mag for galaxies with $M_{B_T} > -17$ and 0.6 mag for brighter galaxies (see Figure 3); as it is reasonable to expect at constant T type, higher mass galaxies had more metallicity evolution and therefore show a larger spread in RGB color.

5. ANALYSIS

Using many galaxies to constrain TP-AGB evolution has added benefits beyond statistically large samples of TP-AGB stars. The more galaxies available, the less the resultant constraints are subject to a specific TP-AGB population and a specific SFH. However, these benefits come with their own limitations. We only derive model constraints from the mean properties of the observed TP-AGB populations and therefore

do not tune models of one specific combination of age, mass, and metallicity. Our models are deemed satisfactory when they reach broad agreement with the largest span of age, metallicity, and mass accessible from our galaxy sample. Below we present three comparisons of the R14 models to the galaxy sample. First we calculate the relative number of TP-AGB stars to RGB stars, the N_{TP-AGB}/N_{RGB} ratio; second, we compare the shape of the observed and predicted LFs, and finally we compare the flux contribution of TP-AGB stars compared the integrated $F160W$ flux of the image.

5.1. Ratio of TP-AGB to RGB Stars

The comparison between the observed and predicted N_{TP-AGB}/N_{RGB} ratio is a rough estimate how well the predicted mean lifetimes of TP-AGB stars agree with those observed. The general assumption is that the number of stars in a given portion of a CMD is related to the amount of time they spend in that portion of the CMD. If by experimental design, the CMD

Figure 6. (Continued.)

filters are set to optimize the separation of stellar evolution phases, this allows researchers to measure relative lifetimes of different evolutionary phases. Table 2 lists the numbers of TP-AGB stars found in the data (column 1) and predicted in the model (column 2) and the $N_{\text{TP-AGB}}/N_{\text{RGB}}$ for each field (column 3). The model $N_{\text{TP-AGB}}/N_{\text{RGB}}$ listed for each field is the median of the ~ 100 TRILEGAL models (column 4). The fractional differences between model and data (column 5) serve as a comparison. Each column also contains Poisson uncertainties. A fractional difference of 0 would mean a perfect agreement and negative values indicate a relative overproduction of model TP-AGB stars compared to the number observed.

From field to field, the fractional difference varies largely, from nearly 65% overproduction of TP-AGB stars (UGCA 292) to nearly 40% underproduction (KDG 73). Due to uncertainties in SFH (see Section 2.3) and variations in TP-AGB populations from galaxy to galaxy, it is more useful to combine the measurements for a rough estimate on the

adequacy of the model. Taking the mean $N_{\text{TP-AGB}}/N_{\text{RGB}}$ of the data and model shows a fractional difference of 9%, meaning the R14 models show good agreement, but slightly underpredict the mean TP-AGB lifetimes based on these observations.

5.2. Luminosity Functions in the Optical and NIR

Model and observed LFs are shown in Figure 6. The M12 catalog LFs (dark gray) are completeness corrected using the ANGST and AGB-SNAP ASTs and are shown with Poisson uncertainties calculated for each magnitude bin. The ~ 100 model LFs are combined and shown as 67% confidence level intervals shaded in gray. The TP-AGB stars, identified by the method discussed in Section 2.4, are shown for the data (red) and for the model (blue) with the same confidence level intervals as the full LF (shaded blue). Vertical lines in each plot show the TRGB and 0.2 mag region around the TRGB excluded from the $N_{\text{TP-AGB}}/N_{\text{RGB}}$ calculations (see

Figure 6. (Continued.)

Section 2.4). The fainter vertical line outlines the RGB box magnitude (but not color) limits used to scale the TRILEGAL simulations. Finally, the top of each panel shows the listing from Table 2 of the number of TP-AGB and RGB stars observed and modeled as well as the corresponding $N_{\text{TP-AGB}}/N_{\text{RGB}}$ ratios.

The LFs all follow the same trends, nearly a power law in the $F160W$ at magnitudes faintward of the TRGB and some degree of agreement and spread in the TP-AGB region (brightward of the TRGB). The departure from a power law of the observations at faint magnitudes, especially in the optical (despite being completeness corrected), is due to stars not detected in both $F814W$ and $F160W$ in the M12 four-filter matched catalogs. In nearly every galaxy in both filters, the spread in the model TP-AGB LF from random SFH uncertainties is consistent with the data TP-AGB LF. The exceptions are the optical LFs of NGC 300-WIDE1, UGC 4305-1, and UGC 5139 where the model underpredicts

the number of TP-AGB stars ~ 0.5 mag brighter than the TRGB.

5.3. Flux Contribution of TP-AGB Stars

One issue quantified in the M12 study was the overproduction of TP-AGB stars using the Padua models outside of the MCs. With the R14 models, it is worth revisiting the issue of the flux contribution of TP-AGB stars. Figure 7 is a reproduction of Figure 7 of M12 (gray) with our results overplotted (blue). Perfect agreement between data and model would be a ratio equaling 1, which is shown as a black solid horizontal line. Mean values for each model are shown as a dashed line. The Marigo et al. (2008) model (gray diamonds) emphasizes the initial problem of extrapolating models calibrated in the MCs alone. On average, they overpredict both the number counts of TP-AGB stars (left panel) and the NIR ($F160W$) TP-AGB flux (right panel). The G10 models (gray squares), with the introduction of pre-dust mass loss,

Figure 7. R14 TP-AGB models improve the prediction the number counts of TP-AGB stars (left panel) as well as $F160W$ TP-AGB flux (right panel). Shown in comparison is a modified version of Figure 7 of M12: Padova 2008 (gray diamonds; Marigo et al. 2008) and Padova 2010 (gray squares; G10). The data and model are recalculated from M12 tables with their weighted mean values shown as dashed gray horizontal lines. On average, galaxies in the present study with R14 TP-AGB models (blue) show much better agreement with the data.

Figure 8. Predicted TP-AGB age (top) and initial mass (bottom) distributions of each galaxy from the median values of the ~ 100 TRILEGAL simulations that sampled the random uncertainties of the derived SFH.

Figure 9. Predicted TP-AGB mass distributions of each galaxy from the ~ 100 TRILEGAL simulations that sampled the random uncertainties of the derived SFH. The black histogram is the median LF plus the gray dashed area for the 67% confidence level interval of the combined simulations. Note masses $M \gtrsim 3 M_{\odot}$ are TP-AGB stars that are expected to experience HBB.

drastically reduce the offsets apparent in the Marigo et al. (2008) model. In contrast, the R14 models show the best agreement on average for both the observed number counts of TP-AGB stars (0.80) as well as the TP-AGB flux (0.90).

Finally, we present the relative age and mass distributions of TP-AGB stars in Figure 8 and a more detailed picture of the simulated TP-AGB mass distributions in Figure 9.

6. CONCLUSIONS

The potential of using nearby galaxies with resolved SFHs to constrain TP-AGB evolution was demonstrated in our previous works (G10, R14), in which evolutionary parameters have been changed in order to improve the agreement between the model predictions and the data. The present work is intended to continue along this line with the addition of galaxies containing younger populations, hence allowing us to calibrate the parameters of TP-AGB stars with larger initial masses even into the regime of TP-AGB stars that experience HBB. We find that the models produced by R14 show good agreement in most cases with a few exceptions. First, the $N_{\text{TP-AGB}}/N_{\text{RGB}}$ ratio, a proxy for the mean TP-AGB lifetimes, showed a mean fractional difference of data to model of 9% across the entire galaxy sample. Second, we compared predicted and observed LFs, a more complex estimate of the TP-AGB lifetimes than

the $N_{\text{TP-AGB}}/N_{\text{RGB}}$ ratio. All but a few cases show consistent agreement to within random SFH uncertainties. Finally, we compared the predicted TP-AGB flux to the observed and found significant improvement compared to older TP-AGB models. The few problematic cases we find (KDG 73, UGCA 292) probably correspond to cases in which the SFH was not as well measured due to the depth of the photometry and/or the large distance to the galaxies.

We consider the present R14 models as satisfactory and that the TP-AGB calibration process needs new and independent data to proceed, aiming to an even more stringent level of agreement between data and models.

The fact that the present models satisfactorily reproduce the numbers and LFs of the present *HST* optical and NIR data does not necessarily mean that they will equally well reproduce data in other passbands, even for the same galaxies. As discussed in R14, the present calibration involved changes mainly in the mass-loss formula, as this was the largest parameter affecting the lifetime along the TP-AGB. There may be other ways of obtaining similar results via changes in other processes as well, including dredge-up efficiency and the mixing length parameter along with different mass-loss formalisms. These alternative ways, however, would probably imply different results for the mid-infrared photometry, which we cannot test with the present observational data for this sample.

A TP-AGB calibration following a multi-band and multi-parametric approach, applied simultaneously to many galaxies, is the clear path forward in understanding the physical properties of large numbers of TP-AGB stars. This process will be pursued in forthcoming works by our team, targeting more complete databases in the Magellanic Clouds and in M31, in addition to ANGST galaxies.

Present TP-AGB models are available at the STARKEY web server <http://starkey.astro.unipd.it>. The derived isochrones are provided via the <http://stev.oapd.inaf.it/cmd> and <http://starkey.astro.unipd.it/cmd> web interfaces. The regular grid of a subset of PARSEC V1.2S tracks used in MATCH is maintained at https://github.com/philrosenfield/padova_tracks. The full grid of PARSEC V1.2S tracks is available at <http://people.sissa.it/~sbressan/parsec.html>.

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