




Flat F-Manifolds, F-CohFTs, and Integrable Hierarchies

Alessandro Arsie¹, Alexandr Buryak^{2,3,4}, Paolo Lorenzoni^{5,6}, Paolo Rossi⁷ 

¹ Department of Mathematics and Statistics, The University of Toledo, 2801W, Bancroft St., Toledo, OH 43606, USA. E-mail: alessandro.arsie@utoledo.edu

² Faculty of Mathematics, National Research University Higher School of Economics, 6 Usacheva str., Moscow, Russian Federation 119048. E-mail: aburyak@hse.ru

³ Center for Advanced Studies, Skolkovo Institute of Science and Technology, 1 Nobel str., Moscow, Russian Federation 143026

⁴ P.G. Demidov Yaroslavl State University, 14 Sovetskaya str., Yaroslavl, Russian Federation 150003

⁵ Dipartimento di Matematica e Applicazioni, Università di Milano-Bicocca, Via Roberto Cozzi 53, 20125 Milan, Italy

⁶ INFN sezione di Milano-Bicocca, Milan, Italy. E-mail: paolo.lorenzoni@unimib.it

⁷ Dipartimento di Matematica “Tullio Levi-Civita”, Università degli Studi di Padova, Via Trieste 63, 35121 Padua, Italy. E-mail: paolo.rossi@math.unipd.it

Received: 23 December 2020 / Accepted: 28 April 2021

© The Author(s) 2021

Abstract: We define the double ramification hierarchy associated to an F-cohomological field theory and use this construction to prove that the principal hierarchy of any semisimple (homogeneous) flat F-manifold possesses a (homogeneous) integrable dispersive deformation at all orders in the dispersion parameter. The proof is based on the reconstruction of an F-CohFT starting from a semisimple flat F-manifold and additional data in genus 1, obtained in our previous work. Our construction of these dispersive deformations is quite explicit and we compute several examples. In particular, we provide a complete classification of rank 1 hierarchies of DR type at the order 9 approximation in the dispersion parameter and of homogeneous DR hierarchies associated with all 2-dimensional homogeneous flat F-manifolds at genus 1 approximation.

Contents

1.	Double Ramification Hierarchy of an F-CohFT	
1.1	F-cohomological field theories	
1.2	Vector fields on the formal loop space	
1.3	Densities of local vector fields for the DR hierarchy	
1.4	Densities of integrals of motion for the DR hierarchy	
1.5	Homogeneous DR hierarchies	
2.	Principal Hierarchy of a Flat F-Manifold and Dispersive Deformations	
2.1	Flat F-manifolds	
2.2	Principal hierarchy of a flat F-manifold	
2.3	Dispersive deformations of a principal hierarchy: descendant DR hierarchies	
2.4	Nonequivalence of dispersive deformations	
2.5	Homogeneous dispersive deformations	
3.	Towards a Classification of Dispersive Deformations	
3.1	Dispersive deformations of DR type and the rank 1 case	

3.1.1	Integrable systems of DR type	
3.1.2	Classification of rank 1 hierarchies of DR type	
3.2	Homogeneous dispersive deformations and the rank 2 case	
3.2.1	Homogeneous deformations with string and dilaton property	
3.2.2	Classification of semisimple homogeneous flat F-manifolds in dimension 2	
3.2.3	Integrable deformations of rank 2 homogeneous principal hierarchies	
3.3	General integrable deformations and open problems	

Introduction

Since Witten’s conjecture [Wit91] and its proof by Kontsevich [Kon92], there have been growing and fruitful interactions between the area of integrable hierarchies of PDEs and algebraic geometry of the moduli spaces of algebraic curves. In this context, and in connection with topological field theory, Dubrovin introduced in the 90s the notion of *Frobenius manifold* [Dub96], a differential-geometric structure that encodes genus-zero information of a cohomological field theory (CohFT) on the moduli space of stable curves, besides having far reaching connections with other areas of mathematics.

From the point of view of integrable systems, given a Dubrovin–Frobenius manifold, there exists an associated integrable hierarchy of Hamiltonian quasilinear PDEs called *Dubrovin’s principal hierarchy*, or simply *principal hierarchy*. An important problem in the theory of integrable systems consists in constructing a full dispersive hierarchy starting from its dispersionless limit.

In the framework of moduli spaces, the principal hierarchy associated to a Dubrovin–Frobenius manifold and its dispersive deformation should satisfy additional constraints coming from the intersection theory of the CohFT. In the semisimple case, there exist two different (but conjecturally Miura-equivalent [Bur15, BDGR18, BGR19]) constructions defining such dispersive deformations:

- (1) The Dubrovin–Zhang construction [DZ01] is based on the idea that the partition function of the corresponding CohFT in all genera is the logarithm of the tau-function of a special solution (called the *topological solution*) to a full dispersive hierarchy (the *DZ hierarchy*). One can construct the hierarchy itself starting from this tau-function, and it turns out that the principal hierarchy is the dispersionless limit of DZ hierarchy. Moreover the full DZ hierarchy and the principal hierarchy are related by a special change of dependent variables, called a *quasi-Miura transformation*, which can be uniquely determined in the semisimple case from genus zero information.
- (2) The double ramification construction, introduced by one of the authors in [Bur15], is based on the definition of an infinite set of commuting Hamiltonian densities [BR16a] in terms of intersection numbers of the CohFT, the double ramification cycles and other natural tautological classes on the moduli space of curves.

For both constructions and in the (homogeneous) semisimple case, the reconstruction of the full dispersive hierarchy from its dispersionless limit (the principal hierarchy of the Dubrovin–Frobenius manifold encoding the genus 0 part of the CohFT) is possible thanks to the Givental–Teleman reconstruction theorem for the CohFT itself from its genus 0 part [Tel12, Giv01].

Notice that, by construction, the dispersionless limits of both the DZ and DR hierarchies coincide with the principal hierarchy of the Dubrovin–Frobenius manifold underlying the CohFT.

In the last 20 years, it has been observed that many constructions related to Dubrovin–Frobenius manifolds can be extended to a more general setting [Sab98, Get04, Man05, LPR09, SZ11, AL13a, Lor14, KMS15, AL17, DH17, BR18, KMS18, AL19, BB19, ABLR20]. For instance, it was observed in [LPR09] that the notion of principal hierarchy does not require the existence of an invariant flat metric. This leads naturally to the consideration of the generalization of Dubrovin–Frobenius manifolds, called F-manifolds with compatible flat structure [Man05] or simply *flat F-manifolds* [LPR09], obtained by replacing a flat metric with a flat torsionless connection and keeping all the axioms of Dubrovin–Frobenius manifolds apart from those involving explicitly the metric and not just the associated Levi–Civita connection. In flat coordinates for the flat connection, the flows of the principal hierarchy are systems of conservation laws. In the case of Dubrovin–Frobenius manifolds, the presence of an invariant flat metric has to deal with the presence of a local Hamiltonian structure. In this paper we construct (homogeneous) double ramification hierarchies starting from a (homogeneous) CohFT. In particular, in the semisimple case, leveraging on the results of [ABLR20], this provides dispersive deformations of the principal hierarchy associated to a semisimple (homogeneous) flat F-manifold. The existence of these dispersive integrable deformations relies on:

- (1) a generalization of the notion of cohomological field theory, called *F-cohomological field theory* (or F-CohFT for short) introduced in [BR18, ABLR20];
- (2) a reconstruction theorem for a semisimple (homogeneous) F-CohFT starting from a flat F-manifold and additional data in genus 1 [ABLR20];
- (3) the definition of an infinite set of commuting flows (the *DR hierarchy*) in terms of intersection numbers of the F-CohFT, the double ramification cycles, the top Hodge class, and psi classes on the moduli space of stable curves.

The paper is organized as follows.

Section 1 is devoted to the construction of the DR hierarchy of an F-CohFT (see also [BR18]). The main properties of this hierarchy are given in terms of densities of local vector fields on the formal loop space and a special basis for their integrals of motion. We also consider the additional properties of the hierarchy in the case of a homogeneous F-CohFT.

In Sect. 2, after recalling the definition of a flat F-manifold and the construction of its associated principal hierarchy, we present our main result: given an arbitrary semisimple flat F-manifold and an associated principal hierarchy, we construct a family of dispersive integrable deformations of the principal hierarchy. These deformations, called the *descendant DR hierarchies*, come from the family of DR hierarchies associated to a family of F-CohFTs parameterized by a semisimple point of our flat F-manifold. The descendant DR hierarchy depends on a choice of a certain vector field on the flat F-manifold, which we call a *framing*. We prove that the descendant DR hierarchies corresponding to different framings are not related to each other by a Miura transformation that is close to identity.

In Sect. 3, we discuss the role of (descendant) DR hierarchies in the problem of classification of integrable deformations of integrable dispersionless systems of conservation laws. One can impose various constraints for such integrable deformations, and we discuss the corresponding results (mostly at the approximation up to some finite power of ε) for flat F-manifolds of dimension 1 and 2 in Sects. 3.2 and 3.1.1. In Sect. 3.3, we briefly mention the problem of computing general integrable deformations of principal hierarchies of flat F-manifolds. It was conjectured in [AL18] that the equivalence classes of such deformations are labeled by certain functional parameters called *Miura invariants*. In the case of Dubrovin–Frobenius manifolds and bihamiltonian deformations, these

invariants are equivalent to central invariants, which are known to classify deformations of semisimple local bihamiltonian structures of hydrodynamic type [DLZ06, CPS18].

1. Double Ramification Hierarchy of an F-CohFT

In this section, we associate to any F-CohFT with a vector space V an infinite sequence of commuting vector fields on the formal loop space of V , i.e., an infinite sequence of compatible systems of evolutionary PDEs of rank $N := \dim V$ (in particular, in the form of conservation laws). This construction is a generalization of the double ramification hierarchy of [Bur15, BR16a] to the context of F-CohFTs and enjoys most of its properties (for instance, recursion formulas for the higher symmetries), but loses in general the Hamiltonian nature.

1.1. F-cohomological field theories. We recall from [BR18, ABLR20] the definition of an F-cohomological field theory on the moduli space $\overline{\mathcal{M}}_{g,n}$ of stable curves of genus g with n marked points. We will denote by $H^*(X)$ the cohomology ring with coefficients in \mathbb{C} of a topological space X . When considering the moduli space of stable curves, $X = \overline{\mathcal{M}}_{g,n}$, the even part $H^{\text{even}}(\overline{\mathcal{M}}_{g,n})$ in the cohomology ring $H^*(\overline{\mathcal{M}}_{g,n})$ can optionally be replaced by the Chow ring $A^*(\overline{\mathcal{M}}_{g,n})$. Here and in what follows we perform summation over repeated Greek indices.

Definition 1. An *F-cohomological field theory* (or F-CohFT) is a system of linear maps

$$c_{g,n+1}: V^* \otimes V^{\otimes n} \rightarrow H^{\text{even}}(\overline{\mathcal{M}}_{g,n+1}), \quad 2g - 1 + n > 0,$$

where V is an arbitrary finite dimensional vector space, together with a special element $e \in V$, called the *unit*, such that, chosen any basis $e_1, \dots, e_{\dim V}$ of V and the dual basis $e^1, \dots, e^{\dim V}$ of V^* , the following axioms are satisfied:

- (i) The maps $c_{g,n+1}$ are equivariant with respect to the S_n -action permuting the n copies of V in $V^* \otimes V^{\otimes n}$ and the last n marked points in $\overline{\mathcal{M}}_{g,n+1}$, respectively.
- (ii) $\pi^* c_{g,n+1}(e^{\alpha_0} \otimes \otimes_{i=1}^n e_{\alpha_i}) = c_{g,n+2}(e^{\alpha_0} \otimes \otimes_{i=1}^n e_{\alpha_i} \otimes e)$ for $1 \leq \alpha_0, \alpha_1, \dots, \alpha_n \leq \dim V$, where $\pi: \overline{\mathcal{M}}_{g,n+2} \rightarrow \overline{\mathcal{M}}_{g,n+1}$ is the map that forgets the last marked point. Moreover, $c_{0,3}(e^\alpha \otimes e_\beta \otimes e) = \delta_{\beta}^\alpha$ for $1 \leq \alpha, \beta \leq \dim V$.
- (iii) $\text{gl}^* c_{g_1+g_2,n_1+n_2+1}(e^{\alpha_0} \otimes \otimes_{i=1}^{n_1+n_2} e_{\alpha_i}) = c_{g_1,n_1+2}(e^{\alpha_0} \otimes \otimes_{i \in I} e_{\alpha_i} \otimes e_\mu) \otimes c_{g_2,n_2+1}(e^\mu \otimes \otimes_{j \in J} e_{\alpha_j})$ for $1 \leq \alpha_0, \alpha_1, \dots, \alpha_{n_1+n_2} \leq \dim V$, where $I \sqcup J = \{2, \dots, n_1+n_2+1\}$, $|I| = n_1$, $|J| = n_2$, and $\text{gl}: \overline{\mathcal{M}}_{g_1,n_1+2} \times \overline{\mathcal{M}}_{g_2,n_2+1} \rightarrow \overline{\mathcal{M}}_{g_1+g_2,n_1+n_2+1}$ is the corresponding gluing map.

An F-CohFT taking value in $H^0(\overline{\mathcal{M}}_{g,n+1}) = \mathbb{C}$ only is called an *F-topological field theory* (or F-TFT). Moreover, there is an obvious generalization of the notion of an F-CohFT where the maps $c_{g,n+1}$ take value in $H^{\text{even}}(\overline{\mathcal{M}}_{g,n+1}) \otimes K$, where K is a \mathbb{C} -algebra. We will call such objects *F-cohomological field theories with coefficients in K* .

Definition 2. An F-CohFT $c_{g,n+1}: V^* \otimes V^{\otimes n} \rightarrow H^{\text{even}}(\overline{\mathcal{M}}_{g,n+1})$ is called *homogeneous* if there exists an operator $Q \in \text{End}(V)$, a vector $\bar{r} \in V$, and a complex constant γ such that $Qe = 0$ and the following condition is satisfied:

$$\text{Deg} \circ c_{g,n+1} + \pi_* \circ c_{g,n+2} \circ (\otimes \bar{r})$$

$$= c_{g,n+1} \circ \left(-Q^t \otimes \text{Id}^{\otimes n} + \sum_{i+j=n-1} \text{Id} \otimes \text{Id}^{\otimes i} \otimes Q \otimes \text{Id}^{\otimes j} \right) + \gamma g c_{g,n+1}, \quad (1.1)$$

where $\text{Deg} \in \text{End}(H^*(\overline{\mathcal{M}}_{g,n}))$ is the operator acting on $H^i(\overline{\mathcal{M}}_{g,n})$ by the multiplication by $\frac{i}{2}$, $\pi: \overline{\mathcal{M}}_{g,n+2} \rightarrow \overline{\mathcal{M}}_{g,n+1}$ is the map that forgets the last marked point, $\otimes \bar{r}: V^* \otimes V^{\otimes n} \rightarrow V^* \otimes V^{\otimes(n+1)}$ is the operator of tensor multiplication from the right by \bar{r} , and $Q^t \in \text{End}(V^*)$ is the transposed operator. The constant γ is called the *conformal dimension* of our F-CohFT.

Remark 1. Our definition of a homogeneous F-CohFT is slightly more general, than the one from the paper [ABLR20] where the operator Q was required to be diagonalizable. However, it is easy to see that all the results from [ABLR20] about homogeneous F-CohFTs are true with the new definition (see also Sect. 2.1 with a new definition of a homogeneous flat F-manifold). An example of a homogeneous F-CohFT with a nondiagonalizable operator Q will appear in Sect. 3.

1.2. Vector fields on the formal loop space. Let $\widehat{\mathcal{A}}$ and $\widehat{\Lambda}$ be the spaces of differential polynomials and local functionals in formal (even) variables u_k^α , $1 \leq \alpha \leq N$, $k \geq 0$, and ε , with the differential grading $\text{deg}_{\partial_x} u_k^\alpha = k$, $\text{deg}_{\partial_x} \varepsilon = -1$, where the definitions and the notations are taken from [Ros17, Section 2.1].

The *space of densities of local multivector fields* (on the formal loop space of V) is the supercommutative associative algebra

$$\widehat{\mathcal{A}}^\bullet := \mathbb{C}[[u^*, \theta_*]][[u_{>0}^*, \theta_{*,>0}]][[\varepsilon]],$$

where the new formal variables $\theta_{\alpha,k}$, $1 \leq \alpha \leq N$, $k \geq 0$, are odd (anti-commuting among themselves and commuting with ε and u_k^α) with $\text{deg}_{\partial_x} \theta_{\alpha,k} = k$, $u^\alpha := u_0^\alpha$, and $\theta_\alpha := \theta_{\alpha,0}$, and the symbol $*$, as an index, denotes any of the allowed values for that index. The algebra $\widehat{\mathcal{A}}^\bullet$ is endowed with the *super grading*, denoted by deg_θ , which is defined by $\text{deg}_\theta \theta_{\alpha,k} := 1$ and $\text{deg}_\theta u_k^\alpha = \text{deg}_\theta \varepsilon := 0$. The sub-vector space of $\widehat{\mathcal{A}}^\bullet$ homogeneous of super degree $i \geq 0$ is denoted by $\widehat{\mathcal{A}}^i$ and called the *space of densities of local i -vector fields*. We have $\widehat{\mathcal{A}} = \widehat{\mathcal{A}}^0$, while $\widehat{\mathcal{A}}^1$ is called the *space of densities of local vector fields*. The homogeneous component of the space $\widehat{\mathcal{A}}^i$ of differential degree k will be denoted by $(\widehat{\mathcal{A}}^i)^{[k]}$.

The operator ∂_x is extended from $\widehat{\mathcal{A}}$ to $\widehat{\mathcal{A}}^\bullet$ as the super-derivation

$$\partial_x := \sum_{k \geq 0} \left(u_{k+1}^\alpha \frac{\partial}{\partial u_k^\alpha} + \theta_{\alpha,k+1} \frac{\partial}{\partial \theta_{\alpha,k}} \right).$$

The *space of local multivector fields* is defined as

$$\widehat{\Lambda}^\bullet := \widehat{\mathcal{A}}^\bullet / (\text{Im } \partial_x \oplus \mathbb{C}[[\varepsilon]])$$

and, for $i \geq 0$, the *space of local i -vector fields* $\widehat{\Lambda}^i$ is the image of $\widehat{\mathcal{A}}^i$ in the quotient. If $f \in \widehat{\mathcal{A}}^\bullet$, its image in $\widehat{\Lambda}^\bullet$ is denoted by $\widehat{f} = \int f dx$. As before, $\widehat{\Lambda} = \widehat{\Lambda}^0$, and $\widehat{\Lambda}^1$ is called the *space of local vector fields*. Naturally, the spaces $\widehat{\Lambda}^i$ inherit the differential grading deg_{∂_x} .

For any $1 \leq \alpha \leq N$, we define the (*super*) *variational derivatives*

$$\frac{\delta}{\delta u^\alpha} := \sum_{k \geq 0} (-\partial_x)^k \frac{\partial}{\partial u_k^\alpha}, \quad \frac{\delta}{\delta \theta_\alpha} := \sum_{k \geq 0} (-\partial_x)^k \frac{\partial}{\partial \theta_{\alpha,k}},$$

which are well defined on $\widehat{\Lambda}^\bullet$ since they vanish on $\text{Im } \partial_x \oplus \mathbb{C}[[\varepsilon]]$.

The *Schouten–Nijenhuis bracket* $[\cdot, \cdot]: \widehat{\Lambda}^i \times \widehat{\Lambda}^j \rightarrow \widehat{\Lambda}^{i+j-1}$ is defined by

$$[\bar{f}, \bar{g}] := \sum_{k \geq 0} \left(\partial_x^k \left(\frac{\delta \bar{f}}{\delta \theta_\alpha} \right) \frac{\partial \bar{g}}{\partial u_k^\alpha} + (-1)^i \partial_x^k \left(\frac{\delta \bar{f}}{\delta u^\alpha} \right) \frac{\partial \bar{g}}{\partial \theta_{\alpha,k}} \right). \quad (1.2)$$

This Schouten–Nijenhuis bracket is a lift of the Schouten–Nijenhuis bracket $[\cdot, \cdot]: \widehat{\Lambda}^i \times \widehat{\Lambda}^j \rightarrow \widehat{\Lambda}^{i+j-1}$ defined by

$$[\bar{f}, \bar{g}] := \int \left(\frac{\delta \bar{f}}{\delta \theta_\alpha} \frac{\delta \bar{g}}{\delta u^\alpha} + (-1)^i \frac{\delta \bar{f}}{\delta u^\alpha} \frac{\delta \bar{g}}{\delta \theta_\alpha} \right) dx. \quad (1.3)$$

A further lift of the Schouten–Nijenhuis bracket to $\widehat{\mathcal{A}}^i \times \widehat{\mathcal{A}}^j$ can be defined employing formal Dirac delta functions, similarly to what was done in [BR16b] for the quantum commutator of two differential polynomials,

$$[f(x), g(y)] := \sum_{k,l \geq 0} \left(\frac{\partial f}{\partial \theta_{\alpha,k}}(x) \frac{\partial g}{\partial u_l^\alpha}(y) \partial_x^k \partial_y^l \delta(x-y) + (-1)^i \frac{\partial f}{\partial u_k^\alpha}(x) \frac{\partial g}{\partial \theta_{\alpha,l}}(y) \partial_x^k \partial_y^l \delta(x-y) \right). \quad (1.4)$$

Taking the integral with respect to x of formula (1.4), using that $\int \delta(x-y)g(y)dx = g(y)$, reproduces indeed formula (1.2), and further integration with respect to y gives (1.3).

As usual, for $i = j = 1$, the above Schouten–Nijenhuis brackets are called the *Lie brackets*. For $i = 1$ and $j = 0$, the Schouten–Nijenhuis brackets reduce simply to the differentiation of (a density of) a local functional along (a density of) a vector field, from which we see that the symbol $\theta_{\alpha,k}$ can be interpreted as the operator $\partial_x^k \circ \frac{\delta}{\delta u^\alpha}: \widehat{\Lambda} \rightarrow \widehat{\mathcal{A}}$.

Given a local vector field $\bar{X} \in \widehat{\Lambda}^1$, there is a unique representative $X \in \widehat{\mathcal{A}}^1$ of \bar{X} such that $X = X^\alpha \theta_\alpha$ with $X^\alpha \in \widehat{\mathcal{A}}$. This representative is given by $X = \frac{\delta \bar{X}}{\delta \theta_\alpha} \theta_\alpha$. The system of evolutionary PDEs associated to \bar{X} is

$$\frac{\partial u^\alpha}{\partial t} = \frac{\delta \bar{X}}{\delta \theta_\alpha}(u_*^*, \varepsilon), \quad \alpha = 1, \dots, N. \quad (1.5)$$

Two systems of evolutionary PDEs

$$\begin{aligned} \frac{\partial u^\alpha}{\partial t} &= \frac{\delta \bar{X}}{\delta \theta_\alpha}(u_*^*, \varepsilon), & \alpha = 1, \dots, N, \\ \frac{\partial u^\alpha}{\partial s} &= \frac{\delta \bar{Y}}{\delta \theta_\alpha}(u_*^*, \varepsilon), & \alpha = 1, \dots, N, \end{aligned}$$

are compatible, in the sense that, for any $1 \leq \alpha \leq N$, $\frac{\partial}{\partial t} \frac{\partial u^\alpha}{\partial s} = \frac{\partial}{\partial s} \frac{\partial u^\alpha}{\partial t}$, if and only if the associated local vector fields $\bar{X}, \bar{Y} \in \widehat{\Lambda}^1$ satisfy $[\bar{X}, \bar{Y}] = 0$.

Under a Miura transformation (see [Ros17, Section 2.1] for more details) of the form

$$\tilde{u}^\alpha = \tilde{u}^\alpha(u_*^*; \varepsilon) \in \widehat{\mathcal{A}}^{[0]} = (\widehat{\mathcal{A}}^0)^{[0]}, \quad 1 \leq \alpha \leq N, \quad (1.6)$$

$$\tilde{u}^*|_{u_*^*=0} = 0, \quad \det \left(\frac{\partial \tilde{u}^*}{\partial u^*} \right) \Big|_{u_*^*=0} \neq 0, \quad (1.7)$$

the generators u_*^* and $\theta_{*,*}$ of $\widehat{\mathcal{A}}^\bullet$ transform according to the formulae

$$u_k^\alpha = \partial_x^k u^\alpha(\tilde{u}_*^*; \varepsilon), \quad \theta_{\alpha,k} = \partial_x^k \left(\sum_{s \geq 0} (-\partial_x)^s \left(\frac{\partial \tilde{u}^\alpha}{\partial u_s^\alpha} \Big|_{u_*^*=u_*^*(\tilde{u}_*^*; \varepsilon)} \tilde{\theta}_\mu \right) \right),$$

$$1 \leq \alpha \leq N, \quad k \geq 0,$$

where $u^\alpha(\tilde{u}_*^*; \varepsilon)$ is obtained by inverting $\tilde{u}^\alpha = \tilde{u}^\alpha(u_*^*; \varepsilon)$ order by order in ε . For a local vector field, these formulae give

$$\bar{X} = \int (X^\alpha \theta_\alpha) dx = \int \left(\left(\sum_{s \geq 0} \frac{\partial \tilde{u}^\alpha}{\partial u_s^\alpha} \partial_x^s X^\mu \right) \Big|_{u_*^*=u_*^*(\tilde{u}_*^*; \varepsilon)} \tilde{\theta}_\alpha \right) dx,$$

from which we obtain that a system of evolutionary PDEs (1.5) transforms into

$$\frac{\partial \tilde{u}^\alpha}{\partial t} = \tilde{X}^\alpha(\tilde{u}_*^*; \varepsilon) = \left(\sum_{s \geq 0} \frac{\partial \tilde{u}^\alpha}{\partial u_s^\alpha} \partial_x^s X^\mu \right) \Big|_{u_*^*=u_*^*(\tilde{u}_*^*; \varepsilon)}, \quad \alpha = 1, \dots, N.$$

Performing the change of formal variables

$$u_k^\alpha = \partial_x^k \left(\sum_{a \in \mathbb{Z}} p_a^\alpha e^{iax} \right), \quad \theta_{\alpha,k} = \partial_x^k \left(\sum_{a \in \mathbb{Z}} q_{\alpha,a} e^{iax} \right), \quad 1 \leq \alpha \leq N, \quad k \geq 0, \quad (1.8)$$

one can rewrite a density of a local multivector field $f(u_*^*, \theta_{*,*}; \varepsilon) \in (\widehat{\mathcal{A}}^m)^{[d]}$ as a formal Fourier series

$$f = \sum_{\substack{n,s \geq 0 \\ a_1, \dots, a_n \in \mathbb{Z} \\ b_1, \dots, b_m \in \mathbb{Z}}} f_{\alpha_1, \dots, \alpha_n, \beta_1, \dots, \beta_m; s}^{a_1, \dots, a_n, b_1, \dots, b_m} \varepsilon^s p_{a_1}^{\alpha_1} \dots p_{a_n}^{\alpha_n} q_{\beta_1, b_1} \dots q_{\beta_m, b_m} e^{i(\sum_{j=1}^n a_j + \sum_{j=1}^m b_j)x},$$

where the coefficient $f_{\alpha_1, \dots, \alpha_n, \beta_1, \dots, \beta_m; s}^{a_1, \dots, a_n, b_1, \dots, b_m}$, as a function of the indices $a_1, \dots, a_n, b_1, \dots, b_m$, is a homogeneous polynomial of degree $s+d$. Formal Fourier series of this type form a supercommutative associative algebra where the formal variables $q_{*,*}$ are odd. Moreover, the local multivector field \bar{f} corresponds to the constant term of the Fourier series. Similarly to the variables $\theta_{*,*}$, one should interpret the variable $q_{\alpha,a}$ to represent the vector $\frac{\partial}{\partial p_{-a}^\alpha}$. This is coherent with the following formulae for the variational derivatives in the variables p_*^* and $q_{*,*}$:

$$\frac{\delta}{\delta u^\alpha} = \sum_{a \in \mathbb{Z}} e^{iax} \frac{\partial}{\partial p_{-a}^\alpha}, \quad \frac{\delta}{\delta \theta_\alpha} = \sum_{a \in \mathbb{Z}} e^{iax} \frac{\partial}{\partial q_{\alpha, -a}},$$

acting on local multivector fields to give densities of local multivector fields. Accordingly, using the formal Fourier expansion $\delta(x) = \sum_{a \in \mathbb{Z}} e^{iax}$ for the formal Dirac delta function, it is easy to obtain the formula for the Schouten–Nijenhuis bracket (1.4) on $\widehat{\mathcal{A}}^i \times \widehat{\mathcal{A}}^j$ in the new variables:

$$[f(x), g(y)] = \sum_{a \in \mathbb{Z}} \left(\frac{\partial f}{\partial q_{\alpha, a}}(x) \frac{\partial g}{\partial p_{-a}^{\alpha}}(y) + (-1)^i \frac{\partial f}{\partial p_{-a}^{\alpha}}(x) \frac{\partial g}{\partial q_{\alpha, a}}(y) \right), \quad (1.9)$$

from which analogues of (1.2) and (1.3) are easily obtained by integration in x and then y .

1.3. Densities of local vector fields for the DR hierarchy. Denote by $\psi_i \in H^2(\overline{\mathcal{M}}_{g,n})$ the i -th *psi class*, which is the first Chern class of the line bundle over $\overline{\mathcal{M}}_{g,n}$ formed by the cotangent lines at the i -th marked point. Denote by \mathbb{E} the rank g Hodge vector bundle over $\overline{\mathcal{M}}_{g,n}$ whose fibers are the spaces of holomorphic one-forms on stable curves. Let $\lambda_j := c_j(\mathbb{E}) \in H^{2j}(\overline{\mathcal{M}}_{g,n})$, these classes are called the *Hodge classes*.

For any $a_1, \dots, a_n \in \mathbb{Z}$, $\sum_{i=1}^n a_i = 0$, denote by $\text{DR}_g(a_1, \dots, a_n) \in H^{2g}(\overline{\mathcal{M}}_{g,n})$ the *double ramification (DR) cycle*. We refer the reader, for example, to [BSSZ15] for the definition of the DR cycle on $\overline{\mathcal{M}}_{g,n}$, which is based on the notion of a stable map to $\mathbb{C}\mathbb{P}^1$ relative to 0 and ∞ . If not all the multiplicities a_i are equal to zero, then one can think of the class $\text{DR}_g(a_1, \dots, a_n)$ as the Poincaré dual to a compactification in $\overline{\mathcal{M}}_{g,n}$ of the locus of pointed smooth curves $(C; p_1, \dots, p_n)$ satisfying $\mathcal{O}_C(\sum_{i=1}^n a_i p_i) \cong \mathcal{O}_C$. Consider the Poincaré dual to the double ramification cycle $\text{DR}_g(a_1, \dots, a_n)$ in the space $\overline{\mathcal{M}}_{g,n}$. It is an element of $H_{2(2g-3+n)}(\overline{\mathcal{M}}_{g,n})$, and abusing notation it is also denoted by $\text{DR}_g(a_1, \dots, a_n)$.

The restriction $\text{DR}_g(a_1, \dots, a_n)|_{\mathcal{M}_{g,n}^{\text{ct}}}$, where $\mathcal{M}_{g,n}^{\text{ct}}$ is the moduli space of stable curves of compact type, is a homogeneous polynomial in a_1, \dots, a_n of degree $2g$ with the coefficients in $H^{2g}(\mathcal{M}_{g,n}^{\text{ct}})$. This follows from Hain’s formula [Hai13] for the version of the DR cycle defined using the universal Jacobian over $\mathcal{M}_{g,n}^{\text{ct}}$ and the result of the paper [MW13], where it is proved that the two versions of the DR cycle coincide on $\mathcal{M}_{g,n}^{\text{ct}}$ (the polynomiality of the DR cycle on $\overline{\mathcal{M}}_{g,n}$ is proved in [JPPZ17]). The polynomiality of the DR cycle on $\mathcal{M}_{g,n}^{\text{ct}}$ together with the fact that λ_g vanishes on $\overline{\mathcal{M}}_{g,n} \setminus \mathcal{M}_{g,n}^{\text{ct}}$ (see, e.g., [FP00, Section 0.4]) imply that the cohomology class $\lambda_g \text{DR}_g(-\sum_{j=1}^n a_j, a_1, \dots, a_n) \in H^{4g}(\overline{\mathcal{M}}_{g,n+1})$ is a degree $2g$ homogeneous polynomial in the coefficients a_1, \dots, a_n .

Given a vector space V with $\dim V = N$ and a basis $e_1, \dots, e_N \in V$, let $c_{g,n+1}: V^* \otimes V^{\otimes n} \rightarrow H^{\text{even}}(\overline{\mathcal{M}}_{g,n+1})$ be an F-CohFT with unit $e = A^\mu e_\mu$. For $1 \leq \beta \leq N$ and $d \geq 0$, we define the following system of formal Fourier series:

$$Y_{\beta,d} := - \sum_{\substack{g,n \geq 0, 2g+n > 0 \\ a_1, \dots, a_n \in \mathbb{Z}}} \frac{i a (-\varepsilon^2)^g}{n!} \left(\int_{\text{DR}_g(a, -a - \sum_{j=1}^n a_j, a_1, \dots, a_n)} \lambda_g \psi_2^d c_{g,n+2}(e^\alpha \otimes e_\beta \otimes \otimes_{j=1}^n e_{\alpha_j}) \right) q_{\alpha,a} \left(\prod_{j=1}^n p_{a_j}^{\alpha_j} \right) e^{i(a + \sum_{j=1}^n a_j)x}, \quad (1.10)$$

where the integral above is intended to vanish whenever the dimension of the DR cycle does not match the degree of the integrand. Thanks to the polynomiality property of the DR cycle, (1.10) can be rewritten as a system of densities of local vector fields $Y_{\beta,d} \in (\widehat{\mathcal{A}}^1)^{[1]}$ as

$$Y_{\beta,d} = - \sum_{\substack{g,n \geq 0, 2g+n > 0 \\ k_1, \dots, k_n \geq 0 \\ k + \sum_{j=1}^n k_j = 2g}} \frac{\varepsilon^{2g}}{n!} \text{Coef}_{a^k (a_1)^{k_1} \dots (a_n)^{k_n}} \left(\int_{\text{DR}_g(a, -a - \sum_{j=1}^n a_j, a_1, \dots, a_n)} \lambda_g \psi_2^d c_{g,n+2}(e^\alpha \otimes e_\beta \otimes \otimes_{j=1}^n e_{\alpha_j}) \right) \theta_{\alpha, k+1} \prod_{j=1}^n u_{k_j}^{\alpha_j}. \quad (1.11)$$

To this definition, we add the extra densities $Y_{\beta,-1} := -\theta_{\beta,1}$, $1 \leq \beta \leq N$.

The *double ramification hierarchy* associated to the given F-CohFT is the infinite system of local vector fields $\bar{Y}_{\beta,d}$, $1 \leq \beta \leq N$, $d \geq -1$, associated with the above densities or, in terms of evolutionary PDEs, the system

$$\frac{\partial u^\alpha}{\partial t_d^\beta} = \partial_x P_{\beta,d}^\alpha, \quad 1 \leq \alpha, \beta \leq N, \quad d \geq 0, \quad (1.12)$$

where

$$P_{\beta,d}^\alpha := \sum_{\substack{g,n \geq 0, 2g+n > 0 \\ k_1, \dots, k_n \geq 0 \\ \sum_{j=1}^n k_j = 2g}} \frac{\varepsilon^{2g}}{n!} \text{Coef}_{(a_1)^{k_1} \dots (a_n)^{k_n}} \left(\int_{\text{DR}_g(-\sum_{j=1}^n a_j, 0, a_1, \dots, a_n)} \lambda_g \psi_2^d c_{g,n+2}(e^\alpha \otimes e_\beta \otimes \otimes_{j=1}^n e_{\alpha_j}) \right) \prod_{j=1}^n u_{k_j}^{\alpha_j}. \quad (1.13)$$

Let us adopt the convention $P_{\beta,-1}^\alpha := \delta_\beta^\alpha$. Notice that the system of evolutionary PDEs (1.12) carries strictly less information than the corresponding densities (1.11). We have the following result.

Theorem 1 ([BR18]). *All the equations of the DR hierarchy (1.12) are compatible with each other, namely,*

$$\frac{\partial}{\partial t_{d_2}^{\beta_2}} \left(\frac{\partial u^\alpha}{\partial t_{d_1}^{\beta_1}} \right) = \frac{\partial}{\partial t_{d_1}^{\beta_1}} \left(\frac{\partial u^\alpha}{\partial t_{d_2}^{\beta_2}} \right), \quad 1 \leq \alpha, \beta_1, \beta_2 \leq N, \quad d_1, d_2 \geq 0.$$

This theorem is proved in [BR18], but we give another proof in Theorem 2 (see part (ii)). For $1 \leq \beta_1, \beta_2 \leq N$ and $d_1, d_2 \geq 0$, let us define the generating series

$$Y_{\beta_1, d_1; \beta_2, d_2}(x, y) := - \sum_{\substack{g, n \geq 0 \\ a, b_1, b_2, a_1, \dots, a_n \in \mathbb{Z}}} \frac{ia(-\varepsilon^2)^g}{n!} \left(\int_{\text{DR}_g(a, b_1, b_2, a_1, \dots, a_n)} \lambda_g \psi_2^{d_1} \psi_3^{d_2} c_{g, n+3} (e^\alpha \otimes e_{\beta_1} \otimes e_{\beta_2} \otimes \otimes_{j=1}^n e_{\alpha_j}) \right) \cdot q_{\alpha, a} \left(\prod_{j=1}^n p_{a_j}^{\alpha_j} \right) e^{-ib_1 x} e^{-ib_2 y}, \quad (1.14)$$

where we adopt the convention that $\text{DR}_g(a, b_1, b_2, a_1, \dots, a_n) := 0$ when $a + b_1 + b_2 + \sum_{j=1}^n a_j \neq 0$. To this definition, for future convenience, we add $Y_{\beta_1, -1; \beta_2, d}(x, y) = Y_{\beta_1, d; \beta_2, -1}(x, y) := 0$, $1 \leq \beta_1, \beta_2 \leq N$, $d \geq 0$.

We will use the symbol $\mathbb{1}$, as an index, to denote the sum over the values $1 \leq \alpha \leq N$ for that index with the coefficients A^α . For example, $Y_{\mathbb{1}, d} := A^\mu Y_{\mu, d}$, $\theta_{\mathbb{1}, k} := A^\mu \theta_{\mu, k}$, and $\frac{\partial}{\partial t_d^\mathbb{1}} := A^\mu \frac{\partial}{\partial t_d^\mu}$.

Theorem 2. For all $1 \leq \beta_1, \beta_2 \leq N$ and $d_1, d_2 \geq -1$ such that $d_1 + d_2 \geq -1$, we have

- (i) $[Y_{\beta_2, d_2}(y), Y_{\beta_1, d_1}(x)] = \partial_x Y_{\beta_1, d_1+1; \beta_2, d_2}(x, y) - \partial_y Y_{\beta_1, d_1; \beta_2, d_2+1}(x, y)$;
- (ii) $[\bar{Y}_{\beta_2, d_2}, \bar{Y}_{\beta_1, d_1}] = 0$;
- (iii) $[\bar{Y}_{\mathbb{1}, 1}, Y_{\beta_1, d_1}] = \partial_x (D - 1) Y_{\beta_1, d_1+1}$, where $D := \sum_{k \geq 0} \left(u_k^\alpha \frac{\partial}{\partial u_k^\alpha} + \theta_{\alpha, k} \frac{\partial}{\partial \theta_{\alpha, k}} \right) + \varepsilon \frac{\partial}{\partial \varepsilon}$;
- (iv) $[\bar{Y}_{\beta_2, 0}, Y_{\beta_1, d_1}] = \partial_x \frac{\partial}{\partial u^{\beta_2}} Y_{\beta_1, d_1+1}$;
- (v) $Y_{\mathbb{1}, 0} = -u^\alpha \theta_{\alpha, 1} + \partial_x^2 S$, $S \in (\widehat{\mathcal{A}}^1)^{[-1]}$, which implies $\frac{\partial u^\alpha}{\partial t_0^\mathbb{1}} = \partial_x u^\alpha$ for $1 \leq \alpha \leq N$;
- (vi) $\frac{\partial}{\partial u^\mathbb{1}} Y_{\beta_1, d_1+1} = Y_{\beta_1, d_1}$, $\frac{\partial}{\partial u^\mathbb{1}} P_{\beta_1, d_1+1}^{\beta_2} = P_{\beta_1, d_1}^{\beta_2}$;
- (vii) $\frac{\partial}{\partial u^{\beta_2}} P_{\mathbb{1}, 1}^{\beta_1} = D P_{\beta_2, 0}^{\beta_1}$.

Proof. For $n \geq 0$, let us use the notation $[n]$ for the set $\{1, \dots, n\}$.

Let us prove part (i). If $d_1 = -1$ or $d_2 = -1$, then the statement easily follows from the definitions. For $d_1, d_2 \geq 0$, the statement is analogous to [BR16b, Lemma 3.3], and we use [BSSZ15, Corollary 2.2], describing the intersection of the psi classes with the DR cycle, together with the fact that λ_g vanishes on $\overline{\mathcal{M}}_{g, n} \setminus \mathcal{M}_{g, n}^{\text{ct}}$. Let $n \geq 0$ and consider integers a_1, \dots, a_{n+3} with the vanishing sum. For a subset $I = \{i_1, \dots, i_{|I|}\} \subset [n+3]$, $i_1 < i_2 < \dots < i_{|I|}$, denote by A_I the string $a_{i_1}, a_{i_2}, \dots, a_{i_{|I|}}$. For $I, J \subset [n+3] \setminus \{2, 3\}$ with $I \sqcup J = [n+3] \setminus \{2, 3\}$, and for $g_1, g_2 > 0$ with $2g_1 + |I| > 0$, $2g_2 + |J| > 0$, let us denote by $\text{DR}_{g_1}(a_2, A_I, -k) \boxtimes \text{DR}_{g_2}(a_3, A_J, k)$ the cycle in $\overline{\mathcal{M}}_{g_1+g_2, n+3}$ obtained by gluing the two DR cycles at the marked points labeled by the integers $-k$ and k , respectively. Here, the coefficient a_j , $1 \leq j \leq n+3$, is attached to the marked point j . Then we have

$$(a_2\psi_2 - a_3\psi_3)\lambda_g \text{DR}_g(A_{[n+3]}) = \sum_{\substack{I \sqcup J = [n+3] \setminus \{2,3\} \\ k \in \mathbb{Z}, g_1 \geq 0, g_2 \geq 0 \\ g_1 + g_2 = g \\ 2g_1 + |I|, 2g_2 + |J| > 0}} \lambda_g \cdot k \cdot \text{DR}_{g_1}(a_2, A_I, -k) \boxtimes \text{DR}_{g_2}(a_3, A_J, k). \quad (1.15)$$

One then needs to intersect this relation with the class $-a_1 e^{-ia_2 x} e^{-ia_3 y} \psi_2^{d_1} \psi_3^{d_2} c_{g,n+3}$ ($e^{\alpha_1} \otimes \otimes_{i=2}^{n+3} e_{\alpha_i}$), where, as usual, the covector e^{α_1} is attached to the marked point 1 and each vector e_{α_i} is attached to the marked point i . Thanks to the gluing axiom of the F-CohFT, by the definitions (1.10) and (1.14), and after setting $\alpha_2 = \beta_1$ and $\alpha_3 = \beta_2$, the left-hand side of Eq. (1.15) produces the right-hand side of the equation in part (i) and depending on whether, in the above sum, the marked point 1 belongs to the subset I or J , we obtain either of the two terms in the Lie bracket on the left-hand side of the equation in part (i).

Part (ii) is immediately obtained from (i) upon integration in both x and y .

Part (iii) is obtained from (i) after setting $\beta_2 = \mathbb{1}$, $d_2 = 1$ and integrating in y . The generating series $\int Y_{\beta_1, d_1+1; \mathbb{1}, 1}(x, y) dy$ reduces to $(D - 1)Y_{\beta_1, d_1+1}$ thanks to the following simple equality:

$$\begin{aligned} \int_{\text{DR}_g(a, b_1, 0, a_1, \dots, a_n)} \lambda_g \psi_2^{d_1+1} \psi_3 c_{g,n+3}(e^\alpha \otimes e_{\beta_1} \otimes e \otimes \otimes_{j=1}^n e_{\alpha_j}) &= (2g + n) \\ \int_{\text{DR}_g(a, b_1, a_1, \dots, a_n)} \lambda_g \psi_2^{d_1+1} c_{g,n+2}(e^\alpha \otimes e_{\beta_1} \otimes \otimes_{j=1}^n e_{\alpha_j}), & \end{aligned}$$

which is in turn a consequence of the following behavior of the involved cohomology classes with respect to the morphism $\pi : \overline{\mathcal{M}}_{g,n+3} \rightarrow \overline{\mathcal{M}}_{g,n+2}$ forgetting the third marked point:

$$\text{DR}_g(a, b_1, 0, a_1, \dots, a_n) = \pi^* \text{DR}_g(a, b_1, a_1, \dots, a_n), \quad (1.16)$$

$$c_{g,n+3}(e^\alpha \otimes e_{\beta_1} \otimes e \otimes \otimes_{j=1}^n e_{\alpha_j}) = \pi^* c_{g,n+2}(e^\alpha \otimes e_{\beta_1} \otimes \otimes_{j=1}^n e_{\alpha_j}), \quad (1.17)$$

$$\lambda_g = \pi^* \lambda_g, \quad \pi_*(\psi_2^{d_1+1} \psi_3) = (2g + n) \psi_2^{d_1+1}. \quad (1.18)$$

Indeed, the operator D multiplies each term of Y_{β_1, d_1+1} by the number of variables ε , $u_{*,*}$, and $\theta_{*,*}$ appearing in that term, i.e., by $2g + n + 1$.

Part (iv) is similarly obtained from (i) by setting $d_2 = 0$ and integrating in y , as $\int Y_{\beta_1, d_1+1; \beta_2, 0}(x, y) dy$ reduces by definition to $\frac{\partial}{\partial u^{\beta_2}} Y_{\beta_1, d_1+1}$.

To deduce (v), we consider formula (1.10) and notice that, for $(g, n) \neq (0, 1)$,

$$\int_{\text{DR}_g(a, -a - \sum_{j=1}^n a_j, a_1, \dots, a_n)} \lambda_g c_{g,n+2}(e^\alpha \otimes e \otimes \otimes_{j=1}^n e_{\alpha_j}) = \int_{\pi_*(\lambda_g \text{DR}_g(a, -a - \sum_{j=1}^n a_j, a_1, \dots, a_n))} c_{g,n+1}(e^\alpha \otimes \otimes_{j=1}^n e_{\alpha_j}),$$

and $\pi_*(\lambda_g \text{DR}_g(a, -a - \sum_{j=1}^n a_j, a_1, \dots, a_n))$ is divisible by $(a + \sum_{j=1}^n a_j)^2$ as proved in [BDGR18, Lemma 5.1], where $\pi : \overline{\mathcal{M}}_{g,n+2} \rightarrow \overline{\mathcal{M}}_{g,n+1}$ is the map forgetting the second marked point. When $g = 0$ and $n = 1$, we have instead $\text{DR}_0(a, -a - a_1, a_1) = 1$, $\lambda_0 = 1$, and $c_{0,3}(e^\alpha \otimes e \otimes e_{\alpha_1}) = \delta_{\alpha_1}^\alpha$, which gives the desired result.

Part (vi) immediately follows from parts (iv), (v), the properties $\text{Ker}(\partial_x |_{\widehat{\mathcal{A}}_1}) = 0$, $\text{Ker}(\partial_x |_{\widehat{\mathcal{A}}}) = \mathbb{C}[[\varepsilon]]$, and the fact $\partial_x P_{\beta_1, d_1}^{\beta_2} = \frac{\delta}{\delta \theta_{\beta_2}} \bar{Y}_{\beta_1, d_1}$.

For part (vii), we compute $\partial_x \frac{\partial P_{\parallel,1}^{\beta_1}}{\partial u^{\beta_2}} = \frac{\delta}{\delta \theta_{\beta_1}} \int Y_{\parallel,1;\beta_2,0} dx dy = \frac{\delta}{\delta \theta_{\beta_1}} (D-1) \bar{Y}_{\beta_2,0} = D \frac{\delta}{\delta \theta_{\beta_1}} \bar{Y}_{\beta_2,0} = D \partial_x P_{\beta_2,0}^{\beta_1} = \partial_x D P_{\beta_2,0}^{\beta_1}$. \square

1.4. Densities of integrals of motion for the DR hierarchy. The DR hierarchy of a CohFT is a Hamiltonian integrable system [Bur15, BR16a], so the Hamiltonians both generate the commuting vector fields and provide integrals of motion for the hierarchy. In the non-Hamiltonian F-CohFT case, integrals of motion have a separate geometric definition in terms of intersection numbers on the moduli space of curves. For $1 \leq \beta \leq N$ and $d \geq 0$, we define the following system of formal Fourier series:

$$g^{\beta,d} := \sum_{\substack{g,n \geq 0 \\ 2g+n-1 > 0 \\ a_1, \dots, a_n \in \mathbb{Z}}} \frac{(-\varepsilon^2)^g}{n!} \left(\int_{\text{DR}_g(-\sum_{j=1}^n a_j, a_1, \dots, a_n)} \lambda_g \psi_1^d c_{g,n+1}(e^\beta \otimes \otimes_{j=1}^n e_{\alpha_j}) \right) \left(\prod_{j=1}^n p_{a_j}^{\alpha_j} \right) e^{i(\sum_{j=1}^n a_j)x}, \quad (1.19)$$

which, thanks to the polynomiality property of the DR cycle, can be rewritten as differential polynomials $g^{\beta,d} \in \widehat{\mathcal{A}}^{[0]}$ as

$$g^{\beta,d} = \sum_{\substack{g,n \geq 0, 2g+n-1 > 0 \\ k_1, \dots, k_n \geq 0 \\ \sum_{j=1}^n k_j = 2g}} \frac{\varepsilon^{2g}}{n!} \text{Coef}_{(a_1)^{k_1} \dots (a_n)^{k_n}} \left(\int_{\text{DR}_g(-\sum_{j=1}^n a_j, a_1, \dots, a_n)} \lambda_g \psi_1^d c_{g,n+2}(e^\beta \otimes \otimes_{j=1}^n e_{\alpha_j}) \right) \prod_{j=1}^n u_{k_j}^{\alpha_j}. \quad (1.20)$$

To this definition, we add the extra densities of conserved quantities $g^{\beta,-1} := u^\beta$, $1 \leq \beta \leq N$, and the ‘‘primary’’ local vector field $\bar{Y} := -\int g^{\beta,0} \theta_{\beta,1} dx$ or, in other words, $\partial_x g^{\beta,0} = \frac{\delta \bar{Y}}{\delta \theta_\beta}$, $1 \leq \beta \leq N$.

Finally, for $1 \leq \beta_1, \beta_2 \leq N$ and $d_1, d_2 \geq 0$, let us define the generating series

$$g^{\beta_1, d_1}(x, y) := \sum_{\substack{g,n \geq 0, 2g+n > 0 \\ b_1, b_2, a_1, \dots, a_n \in \mathbb{Z}}} \frac{(-\varepsilon^2)^g}{n!} \left(\int_{\text{DR}_g(b_1, b_2, a_1, \dots, a_n)} \lambda_g \psi_1^{d_1} \psi_2^{d_2} c_{g,n+2}(e^{\beta_1} \otimes e_{\beta_2} \otimes \otimes_{j=1}^n e_{\alpha_j}) \right) \left(\prod_{j=1}^n p_{a_j}^{\alpha_j} \right) e^{-ib_1 x} e^{-ib_2 y}. \quad (1.21)$$

To this definition, for future convenience, we add $g^{\beta_1, -1}(x, y) = g^{\beta_1, d}_{\beta_2, -1}(x, y) := 0$, $1 \leq \beta_1, \beta_2 \leq N$, $d \geq 0$.

Theorem 3. *For all $1 \leq \beta_1, \beta_2 \leq N$ and $d_1, d_2 \geq -1$ such that $d_1 + d_2 \geq -1$, we have*

- (i) $[Y_{\beta_1, d_1}(y), g^{\beta_2, d_2}(x)] = \partial_x g^{\beta_2, d_2+1}(x, y) - \partial_y g_{\beta_1, d_1+1}^{\beta_2, d_2}(x, y);$
- (ii) $[\bar{Y}_{\beta_1, d_1}, \bar{g}^{\beta_2, d_2}] = 0;$
- (iii) $[\bar{Y}_{\parallel, 1}, g^{\beta_2, d_2}] = \partial_x (D-1) g^{\beta_2, d_2+1};$
- (iv) $[\bar{Y}_{\beta_1, 0}, g^{\beta_2, d_2}] = \partial_x \frac{\partial}{\partial u^{\beta_1}} g^{\beta_2, d_2+1};$
- (v) $\frac{\partial}{\partial u^\parallel} g^{\beta_1, d_1+1} = g^{\beta_1, d_1};$

$$(vi) \bar{Y}_{\parallel,1} = (D-2)\bar{Y};$$

$$(vii) \bar{Y}_{\beta,0} = \frac{\partial}{\partial u^\beta} \bar{Y}.$$

Proof. The proof of (i) is completely analogous to the proof of (i) in Theorem 2. For $d_1 = -1$ or $d_2 = -1$, the statement easily follows from the definitions. Suppose $d_1, d_2 \geq 0$. Let $n \geq 0$ and consider integers a_1, \dots, a_{n+2} with the vanishing sum. Let us write the same relation as (1.15), but with the psi classes taken at other marked points:

$$(a_1\psi_1 - a_2\psi_2)\lambda_g \text{DR}_g(A_{[n+2]}) = \sum_{\substack{I \sqcup J = [n+2] \setminus \{1,2\} \\ k \in \mathbb{Z}, g_1 \geq 0, g_2 \geq 0 \\ g_1 + g_2 = g \\ 2g_1 + |I|, 2g_2 + |J| > 0}} \lambda_g \cdot k \cdot \text{DR}_{g_1}(a_1, A_I, -k) \boxtimes \text{DR}_{g_2}(a_2, A_J, k).$$

Intersecting this relation with the class $(-i)e^{-a_1ix}e^{-a_2iy}\psi_1^{d_2}\psi_2^{d_1}c_{g,n+2}(e^{\alpha_1} \otimes \otimes_{j=2}^{n+2} e^{\alpha_j})$ and forming the corresponding generating series, we obtain part (i) (after setting $\alpha_2 = \beta_1$ and $\alpha_1 = \beta_2$).

The proof of (ii) to (iv) follows strictly the arguments in the proof of the corresponding parts in Theorem 2.

The proof of part (v) is the same as the proof of part (vi) in Theorem 2.

For the proof of (vi), consider the equation of part (iii) with $d_2 = -1$. Multiplying it by θ_{β_2} , summing over β_2 , and integrating over x we obtain, on the left-hand side,

$$\int [\bar{Y}_{\parallel,1}, u^{\beta_2}] \theta_{\beta_2} dx = \int \frac{\delta \bar{Y}_{\parallel,1}}{\delta \theta_{\beta_2}} \theta_{\beta_2} dx = \bar{Y}_{\parallel,1}$$

and, on the right-hand side,

$$\int \partial_x \left((D-1)g^{\beta_2,0} \right) \theta_{\beta_2} dx = -(D-2) \int g^{\beta_2,0} \theta_{\beta_2,1} dx = (D-2)\bar{Y}.$$

Part (vii) is proved in an analogous fashion starting from (iv). \square

1.5. Homogeneous DR hierarchies. Let $Y_{\beta,q} \in (\widehat{\mathcal{A}}^1)^{[1]}$ and $g^{\alpha,p} \in \widehat{\mathcal{A}}^{[0]}$, $1 \leq \beta, \alpha \leq N$, $q, p \geq -1$, be the densities of local vector fields and of integrals of motion of the DR hierarchy associated to a homogeneous rank N F-CohFT. Let

$$c_{\beta\gamma}^\alpha := c_{0,3}(e^\alpha \otimes e_\beta \otimes e_\gamma) \in \mathbb{C}$$

for $1 \leq \alpha, \beta, \gamma \leq N$.

Consider the following vector field on the space of densities of local multivector fields on the formal loop space:

$$\widehat{E}_\gamma := \sum_{k \geq 0} \left(\left((\delta_\beta^\alpha - q_\beta^\alpha) u_k^\beta + \delta_{k,0} r^\alpha \right) \frac{\partial}{\partial u_k^\alpha} - (\delta_\beta^\alpha - q_\beta^\alpha) \theta_{\alpha,k} \frac{\partial}{\partial \theta_{\beta,k}} \right) + \frac{1-\gamma}{2} \varepsilon \frac{\partial}{\partial \varepsilon},$$

where $q_\alpha^\beta e_\beta := Qe_\alpha$ and $r^\alpha e_\alpha := \bar{r}$. For convenience, let us define $Y_{\alpha,-2} = P_{\alpha,-2}^\beta := 0$ and $g^{\alpha,-2} := A^\alpha$ for all $1 \leq \alpha, \beta \leq N$.

Proposition 1. *For all $1 \leq \alpha \leq N$ and $d \geq -1$, we have*

$$(i) \widehat{E}_\gamma(Y_{\alpha,d}) = dY_{\alpha,d} + q_\alpha^\beta Y_{\beta,d} + r^\gamma c_{\gamma\alpha}^\mu Y_{\mu,d-1};$$

$$\begin{aligned}
(ii) \quad & \widehat{E}_\gamma(P_{\beta,d}^\alpha) = (d+1)P_{\beta,d}^\alpha + q_\beta^\gamma P_{\gamma,d}^\alpha - q_\gamma^\alpha P_{\beta,d}^\gamma + r^\gamma c_{\gamma\beta}^\mu P_{\mu,d-1}^\alpha; \\
(iii) \quad & \widehat{E}_\gamma(g^{\alpha,d}) = (d+2)g^{\alpha,d} - q_\beta^\alpha g_{\beta,d}^\alpha + r^\gamma c_{\gamma\mu}^\alpha g^{\mu,d-1}; \\
(iv) \quad & \widehat{E}_\gamma(\bar{Y}) = \bar{Y} - r^\gamma c_{\gamma\mu}^\beta \int u^\mu \theta_{\beta,1} dx.
\end{aligned}$$

Proof. The proof is a simple consequence of Eq. (1.1) together with dimension counting for the intersection numbers involved in the definitions of $g^{\alpha,d}$, $Y_{\alpha,d}$, and \bar{Y} and the fact that $\pi^* \psi_i^d = \psi_i^d - \delta_{i,n+1}^0 \pi^* \psi_i^{d-1}$, $1 \leq i \leq n$, $d \geq 1$, where $\pi: \overline{\mathcal{M}}_{g,n+1} \rightarrow \overline{\mathcal{M}}_{g,n}$ forgets the last marked point and $\delta_{i,n+1}^0$ is the closure in $\overline{\mathcal{M}}_{g,n+1}$ of the locus of stable curves whose dual graph is a tree with two vertices, one of which has genus 0 and exactly two legs marked by i and $n+1$. \square

In [BRS20], the authors presented an explicit conjectural formula for a bihamiltonian structure of the DR hierarchy corresponding to a homogeneous CohFT. This in particular gives a recursion of certain type, called a *bihamiltonian recursion*, expressing the flows $\frac{\partial}{\partial t_{d+1}^\alpha}$, $1 \leq \alpha \leq N$, of the hierarchy in terms of the flows $\frac{\partial}{\partial t_d^\alpha}$, $1 \leq \alpha \leq N$. For a general homogeneous F-CohFT, we don't expect the corresponding DR hierarchy to have a Hamiltonian structure. However, we will now present a conjectural generalization of the bihamiltonian recursion in this setting.

Following [BRS20], we associate with a differential polynomial $f \in \widehat{\mathcal{A}}$ a sequence of differential operators indexed by $\alpha = 1, \dots, N$ and $k \geq 0$:

$$L_\alpha^k(f) := \sum_{i \geq k} \binom{i}{k} \frac{\partial f}{\partial u_i^\alpha} \partial_x^{i-k}.$$

Consider an arbitrary homogeneous F-CohFT and the corresponding DR hierarchy. Define an operator $R = (R_\beta^\alpha)$ by

$$R_\beta^\alpha := \widehat{E}_\gamma \left(L_\beta^0(g^{\alpha,0}) \right) \circ \partial_x + \left(\frac{1-\gamma}{2} \delta_\beta^\mu + q_\beta^\mu \right) L_\mu^0(g^{\alpha,0})_x + \partial_x \circ L_\beta^1(g^{\alpha,0}) \circ \partial_x,$$

where the notation $\widehat{E}_\gamma \left(L_\beta^0(g^{\alpha,0}) \right)$ (respectively, $L_\beta^0(g^{\alpha,0})_x$) means that we apply the operator \widehat{E}_γ (respectively, ∂_x) to the coefficients of the operator $L_\beta^0(g^{\alpha,0})$.

Conjecture 1. *The following recursion relation is satisfied:*

$$R_\mu^\alpha P_{\beta,d}^\mu = \left(\left(d + \frac{3-\gamma}{2} \right) \delta_\beta^\mu + q_\beta^\mu \right) \partial_x P_{\mu,d+1}^\alpha + (\partial_x P_{\mu,d}^\alpha) c_{\beta\nu}^\mu r^\nu, \quad 1 \leq \alpha, \beta \leq N, \quad d \geq -1. \quad (1.22)$$

Proposition 2. (1) *If our homogeneous F-CohFT comes from a homogeneous CohFT, then the recursion (1.22) coincides with the bihamiltonian recursion from part (2) of [BRS20, Conjecture 1.13].*

(2) *Conjecture 1 is true in genus 0, i.e., if we set $\varepsilon = 0$.*

Proof. For part (1), using the notations from paper [BRS20] let us note that $P_{\beta,d}^\alpha = \eta^{\alpha\mu} \frac{\delta \bar{g}_{\beta,d}}{\delta u^\mu}$. Therefore, we have to check that $\eta_{\beta\mu} K_2^{\alpha\mu} = R_\beta^\alpha$. This follows from the properties $\eta_{\beta\mu} \Omega^k(\bar{g})^{\alpha\mu} = L_\beta^k(g^{\alpha,0})$ and $q_\alpha^\mu \eta_{\mu\beta} + \eta_{\alpha\mu} q_\beta^\mu = \gamma \eta_{\alpha\beta}$.

The proof of part (2) follows closely the proof of [BRS20, Proposition 2.1]. \square

2. Principal Hierarchy of a Flat F-Manifold and Dispersive Deformations

In this section, using the results from the previous section, we construct a family of dispersive integrable deformations of a principal hierarchy associated to an arbitrary semisimple flat F-manifold. Moreover, we prove that different hierarchies from this family are not equivalent to each other by a Miura transformation that is close to identity.

2.1. Flat F-manifolds. Here we recall the notion of a flat F-manifold ([Get04, Man05], see also [AL18, LPR09]) and its main properties.

Definition 3. A flat F-manifold (M, ∇, \circ, e) is the datum of an analytic manifold M , an analytic connection ∇ in the tangent bundle TM , an algebra structure $(T_p M, \circ)$ with unit e on each tangent space, analytically depending on the point $p \in M$, such that the one-parameter family of connections $\nabla_z = \nabla + z\circ$ is flat and torsionless for any $z \in \mathbb{C}$, and $\nabla e = 0$.

The algebras $(T_p M, \circ)$ are commutative and associative. Let t^α , $1 \leq \alpha \leq N$, $N = \dim M$, be flat coordinates for the connection ∇ . Locally, there exist analytic functions $F^\alpha(t^1, \dots, t^N)$, $1 \leq \alpha \leq N$, such that the second derivatives

$$C_{\beta\gamma}^\alpha := \frac{\partial^2 F^\alpha}{\partial t^\beta \partial t^\gamma} \quad (2.1)$$

are the structure constants of the algebras $(T_p M, \circ)$, $\frac{\partial}{\partial t^\beta} \circ \frac{\partial}{\partial t^\gamma} = C_{\beta\gamma}^\alpha \frac{\partial}{\partial t^\alpha}$. Also, in the coordinates t^α the unit e has the form $e = A^\alpha \frac{\partial}{\partial t^\alpha}$ for some constants $A^\alpha \in \mathbb{C}$. Moreover, the following equations are satisfied:

$$A^\mu \frac{\partial^2 F^\alpha}{\partial t^\mu \partial t^\beta} = \delta_{\beta}^\alpha, \quad 1 \leq \alpha, \beta \leq N, \quad (2.2)$$

$$\frac{\partial^2 F^\alpha}{\partial t^\beta \partial t^\mu} \frac{\partial^2 F^\mu}{\partial t^\gamma \partial t^\delta} = \frac{\partial^2 F^\alpha}{\partial t^\gamma \partial t^\mu} \frac{\partial^2 F^\mu}{\partial t^\beta \partial t^\delta}, \quad 1 \leq \alpha, \beta, \gamma, \delta \leq N, \quad (2.3)$$

which are often called the *oriented WDVV equations*. The N -tuple of functions $\bar{F} = (F^1, \dots, F^N)$ is called a *vector potential* of the flat F-manifold.

Conversely, given an open subset M of \mathbb{C}^N and analytic functions F^1, \dots, F^N on M satisfying Eqs. (2.2) and (2.3), these functions define a flat F-manifold $(M, \nabla, \circ, A^\alpha \frac{\partial}{\partial t^\alpha})$ with the connection ∇ given by $\nabla_{\frac{\partial}{\partial t^\alpha}} \frac{\partial}{\partial t^\beta} = 0$, and the multiplication \circ given by the structure constants (2.1).

A point $p \in M$ of an N -dimensional flat F-manifold (M, ∇, \circ, e) is called *semisimple* if $T_p M$ has a basis π_1, \dots, π_N satisfying $\pi_\alpha \circ \pi_\beta = \delta_{\alpha, \beta} \pi_\alpha$. Moreover, locally around such a point one can choose coordinates u^i such that $\frac{\partial}{\partial u^\alpha} \circ \frac{\partial}{\partial u^\beta} = \delta_{\alpha, \beta} \frac{\partial}{\partial u^\alpha}$. These coordinates are called *canonical coordinates*. In particular, this means that the set of semisimple points is open in M . In the canonical coordinates, we have $e = \sum_\alpha \frac{\partial}{\partial u^\alpha}$. Following [Dub96, page 196], we call a flat F-manifold (M, ∇, \circ, e) *semisimple* if the set of semisimple points is dense in M .

A flat F-manifold given by a vector potential \bar{F} is called *homogeneous* if there exists a vector field E of the form

$$E = \underbrace{((\delta_\beta^\alpha - q_\beta^\alpha)t^\beta + r^\alpha)}_{=: E^\alpha} \frac{\partial}{\partial t^\alpha}, \quad q_\beta^\alpha, r^\alpha \in \mathbb{C}, \quad (2.4)$$

satisfying $[e, E] = e$ and such that

$$E^\mu \frac{\partial F^\alpha}{\partial t^\mu} = (2\delta_\beta^\alpha - q_\beta^\alpha) F^\beta + A_\beta^\alpha t^\beta + B^\alpha$$

for some $A_\beta^\alpha, B^\alpha \in \mathbb{C}$. Note that this equation can be written more invariantly as $\text{Lie}_E(\circ) = \circ$, where Lie_E denotes the Lie derivative. The vector field E is called the *Euler vector field*. Around a semisimple point, the Euler vector field has the following form in canonical coordinates: $E = \sum_{i=1}^N (u^i + a^i) \frac{\partial}{\partial u^i}$ for some $a^i \in \mathbb{C}$.

Remark 2. As we already mentioned in Remark 1, our definition of a homogeneous flat F-manifold is slightly more general than the one from [ABLR20], but all the results from that paper remain valid.

Remark 3. In [AL13a], the authors introduced the closely related notion of a *bi-flat F-manifold* that is the datum of two different flat F-manifold structures (∇, \circ, e) and $(\nabla^*, *, E)$ on the same manifold M intertwined by the following conditions: (1) $[e, E] = e$, $\text{Lie}_E(\circ) = \circ$; (2) $X * Y := (E\circ)^{-1} X \circ Y$ (or $X \circ Y = (e*)^{-1} X * Y$) for all local vector fields X, Y on M , where $(E\circ)^{-1}$ is the inverse of the endomorphism of the tangent bundle given by $E\circ$; (3) $(d_\nabla - d_{\nabla^*})(X \circ) = 0$ for all local vector fields X on M , where d_∇ is the exterior covariant derivative. For a bi-flat F-manifold, the flat structure given by $(\nabla^*, *, E)$ is called the *dual structure*. In the semisimple case, the flatness of the dual structure is equivalent to the condition $\nabla \nabla E = 0$ [AL17] (see [KMS18] for the regular case). Thus, the structure of a semisimple homogeneous flat F-manifold is equivalent to the structure of a semisimple bi-flat F-manifold.

Given an F-CohFT $c_{g,n+1} : V^* \otimes V^{\otimes n} \rightarrow H^{\text{even}}(\overline{\mathcal{M}}_{g,n+1})$, $\dim V = N$, and a basis $e_1, \dots, e_N \in V$, with $e = A^\alpha e_\alpha$, an N -tuple of functions (F^1, \dots, F^N) satisfying Eqs. (2.2) and (2.3) can be constructed as the following generating functions:

$$F^\alpha(t^1, \dots, t^N) := \sum_{n \geq 2} \frac{1}{n!} \sum_{1 \leq \alpha_1, \dots, \alpha_n \leq N} \left(\int_{\overline{\mathcal{M}}_{0,n+1}} c_{0,n+1}(e^\alpha \otimes \otimes_{i=1}^n e_{\alpha_i}) \right) \prod_{i=1}^n t^{\alpha_i},$$

thus yielding an associated flat F-manifold structure on a formal neighbourhood of 0 in V (see, e.g., [ABLR20, Proposition 3.2]). The flat F-manifold associated to a homogeneous F-CohFT is homogeneous with the Euler vector field (2.4) where $q_\beta^\alpha e_\alpha := Qe_\beta$ and $r^\alpha e_\alpha := \overline{r}$.

2.2. Principal hierarchy of a flat F-manifold. Given a flat F-manifold (M, ∇, \circ, e) , one can construct an integrable dispersionless hierarchy called a *principal hierarchy* associated to (M, ∇, \circ, e) (see [LPR09]). This construction generalizes the notion of a principal hierarchy associated to a Dubrovin–Frobenius manifold. Before presenting the construction, let us introduce a small generalization of the space of densities of local multivector fields.

Let U be an open subset of \mathbb{C}^N , with coordinates u^1, \dots, u^N . Denote by $\mathcal{O}(U)$ the space of analytic functions on U . Consider the following space:

$$\widehat{\mathcal{A}}_U^\bullet := \mathcal{O}(U)[u_{>0}^*, \theta_{*,*}][[\varepsilon]].$$

Clearly, the space $\widehat{\mathcal{A}}^\bullet$ can be considered as the space $\widehat{\mathcal{A}}_U^\bullet$ where U is a formal neighborhood of 0. The space $\widehat{\mathcal{A}}_U^\bullet$ will also be called the space of densities of local multivector

fields. It is easy to see that all constructions from Sect. 1.2 work with the space $\widehat{\mathcal{A}}_U^\bullet$. The space of local multivector fields corresponding to $\widehat{\mathcal{A}}_U^\bullet$ will be denoted by $\widehat{\Lambda}_U^\bullet$.

Consider a flat F-manifold (M, ∇, \circ, e) . For any point of M , on its open neighbourhood U , one can consider a basis (over $\mathbb{C}[[z]]$) $X_\alpha(z) = \sum_{d=-1}^{\infty} X_{\alpha,d} z^{d+1}$, $1 \leq \alpha \leq N = \dim M$, in the space of flat sections of the deformed connection $\nabla_{-z} = \nabla - z\circ$:

$$0 = (\nabla - z\circ)X_\alpha(z) = (\nabla - z\circ) \sum_{d=-1}^{\infty} X_{\alpha,d} z^{d+1}. \quad (2.5)$$

It is immediate to see from (2.5) that $X_{\alpha,-1}$, $\alpha = 1, \dots, N$, are flat vector fields for ∇ , while the vector fields $X_{\alpha,d+1}$ are obtained via the recurrence relation $\nabla X_{\alpha,d+1} = X_{\alpha,d}\circ$. If U is connected, then the collection of flat sections $X_\alpha(z)$ is determined uniquely up to a transformation of the form $X_\alpha(z) \mapsto X_\beta(z)G_\alpha^\beta(z)$, where $G(z) = (G_\alpha^\beta(z)) \in \text{Mat}_{N,N}(\mathbb{C}[[z]])$ is invertible. If M is simply connected, then flat sections $X_\alpha(z)$ can be constructed on the whole M .

Definition 4. A *calibration* of a flat F-manifold (M, ∇, \circ, e) is a basis $X_\alpha = \sum_{d=-1}^{\infty} X_{\alpha,d} z^{d+1}$, $X_{\alpha,d} \in \mathcal{T}(M)$, $1 \leq \alpha \leq \dim M$, in the space of flat sections of the deformed connection $\nabla - z\circ$. A flat F-manifold with a fixed calibration is called a *calibrated flat F-manifold*.

Consider now a flat F-manifold structure on $M \subset \mathbb{C}^N$ given by a vector potential \overline{F} , together with a calibration $X_\alpha(z)$. The *principal hierarchy* associated to our calibrated flat F-manifold is the following system of PDEs:

$$\frac{\partial u^\alpha}{\partial t_d^\beta} = \partial_x \left(X_{\beta,d}^\alpha \Big|_{t^\nu = u^\nu} \right), \quad 1 \leq \alpha, \beta \leq N, \quad d \geq 0, \quad (2.6)$$

where $X_{\beta,d}^\alpha \frac{\partial}{\partial t_d^\beta} := X_{\beta,d}$. We see that the system (2.6) has the form of a system of conservation laws. Moreover, this is a system of quasilinear evolutionary PDEs, which is dispersionless and integrable, in the sense that all the flows pairwise commute (see [LPR09]).

Suppose that M is a formal neighbourhood of $0 \in \mathbb{C}^N$. There exist unique flat sections $X_\alpha(z)$ on M satisfying the condition $X_{\alpha,-1} = \frac{\partial}{\partial t^\alpha}$ and the condition that $X_{\alpha,d}$ vanish at 0 for $d \geq 0$. The corresponding principal hierarchy is called the *ancestor principal hierarchy*.

Proposition 3. Consider an F-CohFT and the associated flat F-manifold and the DR hierarchy. Then the dispersionless part of the DR hierarchy coincides with the ancestor principal hierarchy of the flat F-manifold.

Proof. This immediately follows from the construction of the DR hierarchy and [ABLR20, Proposition 3.2] (see also an analogous statement in [Bur15, Section 4.2.2]). \square

We see that this proposition can be immediately used for a construction of dispersive deformations of ancestor principal hierarchies. In order to construct dispersive deformations of arbitrary principal hierarchies, we need a generalization of the construction of the DR hierarchy, which we will introduce in the next section.

2.3. *Dispersive deformations of a principal hierarchy: descendant DR hierarchies.* In order to construct dispersive deformations of a principal hierarchy associated to an arbitrary semisimple flat F-manifold, we first need to study analytic families of F-CohFTs depending on a semisimple point of a flat F-manifold together with a vector in the tangent space to that point.

Consider a semisimple flat F-manifold structure on $M \subset \mathbb{C}^N$ defined by a vector potential \bar{F} . Recall that on a connected open neighborhood U of an arbitrary semisimple point $\bar{t}_0 \in M$, one has the following objects (we use the notations from [ABLR20, Section 1.2]):

- canonical coordinates u^i ;
- the matrix $\tilde{\Psi} := \left(\frac{\partial u^i}{\partial t^\alpha} \right)$;
- the matrices \tilde{D} and $\tilde{\Gamma}$ defined by $d\tilde{\Psi} \cdot \tilde{\Psi}^{-1} = \tilde{D} + [\tilde{\Gamma}, dU]$, where \tilde{D} is a diagonal matrix consisting of one-forms, $\tilde{\Gamma}$ is a matrix with vanishing diagonal entries, and $U := \text{diag}(u^1, \dots, u^N)$ (in the homogeneous case this is the operator of multiplication by the Euler vector field);
- a diagonal nondegenerate matrix $H = \text{diag}(H_1, \dots, H_N)$ defined by $dH \cdot H^{-1} = -\tilde{D}$ (the entries of this matrix can be interpreted as the Lamé coefficients of a diagonal metric associated with the flat F-manifold);
- the matrices Ψ and Γ defined by $\Psi := H\tilde{\Psi}$ and $\Gamma := H\tilde{\Gamma}H^{-1}$;
- a sequence of matrices $R_0 = \text{Id}, R_1, R_2, \dots$ defined by the relations $dR_{k-1} + R_{k-1}[\Gamma, dU] = [R_k, dU]$, $k \geq 1$.

Note that the matrix H is defined uniquely up to the transformation $H \mapsto AH$, where A is a constant nondegenerate diagonal matrix. After such a transformation, the matrices Ψ , Γ , and R_k transform as follows: $\Psi \mapsto A\Psi$, $\Gamma \mapsto A\Gamma A^{-1}$, $R_k \mapsto AR_k A^{-1}$. Recall also that if we fix H , then the matrices R_k are defined uniquely up to the transformation

$$\text{Id} + \sum_{i \geq 1} R_i z^i \mapsto \left(\text{Id} + \sum_{i \geq 1} D_i z^i \right) \left(\text{Id} + \sum_{i \geq 1} R_i z^i \right), \quad (2.7)$$

where D_i , $i \geq 1$, are arbitrary constant diagonal matrices.

Using the notations from [ABLR20, Section 4.4], for any $G_0 \in \mathbb{C}^N$, let us define an analytic family of F-CohFTs parameterized by a point $\bar{t} \in U$ by

$$c^{G_0, \bar{t}} := \tilde{\Psi}^{-1}(\bar{t}) H^{-1}(\bar{t}) R^{-1}(-z, \bar{t}) H(\bar{t}) \cdot c^{\text{triv}, H^{-2}(\bar{t}) G_0},$$

where $R(z) := \sum_{i \geq 0} R_i z^i$, the above matrix action on F-CohFTs was introduced in [ABLR20, Section 4.2], for any $w_1 = (w_1^1, \dots, w_1^N) \in (\mathbb{C}^*)^N$, and $w_2 = (w_2^1, \dots, w_2^N) \in \mathbb{C}^N$,

$$c_{g, n+1}^{w_1, w_2}(e^{i_0} \otimes \otimes_{j=1}^n e_{i_j}) := \begin{cases} \frac{(w_2^{i_0})^g}{(w_1^{i_0})^{g+n-1}}, & \text{if } i_0 = i_1 = \dots = i_n, \\ 0, & \text{otherwise,} \end{cases}$$

and $c^{\text{triv}, w_2} := c^{(1, \dots, 1), w_2}$.

Note that the family $c^{G_0, \bar{t}}$ depends, first, on the choice of H and, second, on the choice of $R(z)$. However, the dependance on H is simple: under the transformation $H \mapsto AH$, where A is a nondegenerate constant diagonal matrix, $R(z)$ transforms as $R(z) \mapsto AR(z)A^{-1}$ (and $\tilde{\Psi}$ doesn't change), and therefore $c^{G_0, \bar{t}} \mapsto c^{A^{-2}G_0, \bar{t}}$.

We see that the simultaneous transformation $H \mapsto AH, G_0 \mapsto A^2G_0$ doesn't change $c^{G_0, \bar{t}}$. This transformation doesn't change the matrix $H^{-2}G_0$, which has a geometrical meaning: if we denote by \mathcal{X}^α the degree zero part of $c_{1,1}^{G_0, \bar{t}}(e^\alpha) \in H^*(\overline{\mathcal{M}}_{1,1})$, then $\mathcal{X}^\alpha \frac{\partial}{\partial t^\alpha} = \sum_{i=1}^N H_i^{-2} G_0^i \frac{\partial}{\partial u^i}$. Let us then choose H such that $H_i(\bar{t}_0) = 1$.

Notice that if $G_0 = 0$, then the maps $c_{g,n+1}^{G_0, \bar{t}}$ are zero for $g \geq 1$.

Let τ^1, \dots, τ^N be formal variables. Recall from [ABLR20, Section 3.2] (note, however, that we prefer to use a different notation here) that for an F-CohFT $c_{g,n+1} : V^* \otimes V^{\otimes n} \rightarrow H^{\text{even}}(\overline{\mathcal{M}}_{g,n+1})$ its formal shift $\text{Sh}_{\bar{\tau}}(c)_{g,n+1} : V^* \otimes V^{\otimes n} \rightarrow H^{\text{even}}(\overline{\mathcal{M}}_{g,n+1})[[\tau^*]]$ is defined by

$$\text{Sh}_{\bar{\tau}}(c)_{g,n+1} := \sum_{m \geq 0} \frac{1}{m!} \pi_{m*} \circ c_{g,n+m+1} \circ (\otimes (\tau^\alpha e_\alpha)^{\otimes m}),$$

where $\bar{\tau} = (\tau^1, \dots, \tau^N)$ and $\pi_m : \overline{\mathcal{M}}_{g,n+m+1} \rightarrow \overline{\mathcal{M}}_{g,n+1}$ forgets the last m marked points. The maps $\text{Sh}_{\bar{\tau}}(c)_{g,n+1}$ form an F-CohFT with the coefficients in $\mathbb{C}[[\tau^*]]$.

Proposition 4.1. *A vector potential of the flat F-manifold corresponding to the F-CohFT c^{G_0, \bar{t}_0} is equal to $\overline{F}(t^* - t_0^*)$.*

2. For $\bar{t} = \bar{t}_0 + \bar{\tau} \in U$, the Taylor expansion of $c^{G_0, \bar{t}}$ at \bar{t}_0 coincides with the formal shift of c^{G_0, \bar{t}_0} , i.e., $c_{g,n+1}^{G_0, \bar{t}_0 + \bar{\tau}} = \text{Sh}_{\bar{\tau}}(c^{G_0, \bar{t}_0})_{g,n+1}$, as elements of $\text{Hom}(V^* \otimes V^{\otimes n}, H^{\text{even}}(\overline{\mathcal{M}}_{g,n+1})[[\tau^*]])$.

Proof. 1. Since $H_i(\bar{t}_0) = 1$, we have $c^{G_0, \bar{t}_0} = \Psi^{-1}(\bar{t}_0)R^{-1}(-z, \bar{t}_0).c^{\text{triv}, G_0}$. The fact that a vector potential of the associated flat F-manifold is equal to $\overline{F}(t^* - t_0^*)$ was proved in [ABLR20, Section 4.4] (see equation (4.3) there).

2. An elementary computation shows that $c^{G_0, \bar{t}} = \Psi^{-1}R^{-1}(-z).c^{\overline{H}, H^{-1}G_0}$, where the vector $\overline{H} := (H_1, \dots, H_N)$ and the matrices $\Psi^{-1}, R^{-1}(-z)$, and H^{-1} are computed at the point \bar{t} . The statement of part 2 of the proposition is equivalent to the property

$$\frac{\partial}{\partial t^\beta} \left(\Psi^{-1}R^{-1}(-z).c^{\overline{H}, H^{-1}G_0} \right)_{g,n+1} = \pi_{1*} \circ \left(\Psi^{-1}R^{-1}(-z).c^{\overline{H}, H^{-1}G_0} \right)_{g,n+2} \circ (\otimes e_\beta),$$

which was proved in [ABLR20, proof of Proposition 4.11]. \square

The degree zero parts of $c_{1,1}^{G_0, \bar{t}}(e^\alpha) \in H^*(\overline{\mathcal{M}}_{1,1})$ induce a vector field on U , $\mathcal{X} = \mathcal{X}^\alpha \frac{\partial}{\partial t^\alpha}$, and we already noticed that $\mathcal{X} = \sum_{i=1}^N H_i^{-2} G_0^i \frac{\partial}{\partial u^i}$. This motivates the following definition.

Definition 5. Consider a semisimple flat F-manifold (M, ∇, \circ, e) . A vector field \mathcal{X} on M is called a *framing* if around each semisimple point of M , in canonical coordinates u^i , the field \mathcal{X} has the form $\mathcal{X} = \sum_{i=1}^N \alpha_i H_i^{-2} \frac{\partial}{\partial u^i}$ for some complex constants α_i , $1 \leq i \leq N$.

Using this language, we can say that our family of F-CohFTs $c^{G_0, \bar{t}}$ induces a framing $\mathcal{X} = \sum_{i=1}^N \mathcal{X}^i \frac{\partial}{\partial u^i}$ on U , with $\mathcal{X}^i(\bar{t}_0) = G_0^i$.

Suppose vice versa that all the points of our flat F-manifold M are semisimple and \mathcal{X} is a framing on M . We see that for any point $\bar{t}_0 \in M$ the above construction gives a family of F-CohFTs around \bar{t}_0 such that the induced framing coincides with \mathcal{X} . This family is not

unique because the matrix $R(z)$ is defined uniquely only up to the transformation (2.7). Suppose that M is simply connected. Then it is easy to see that there is a consistent choice of matrix $R(z)$ in all the charts such that the local families glue in a global family of F-CohFTs parameterized by $\bar{t} \in M$. Let us denote this global family by $c^{\mathcal{X}, \bar{t}}$. This global family is not unique: in order to fix the ambiguity, one can, for example, fix a choice of matrix $R(z)$ at some fixed point of M . Note that if $\mathcal{X} = 0$, then the maps $c_{g, n+1}^{\mathcal{X}, \bar{t}}$ are zero for $g \geq 1$.

Let us now apply the construction of the DR hierarchy to the F-CohFTs $c^{\mathcal{X}, \bar{t}}$. We obtain a family of densities $Y_{\beta, d}^{\bar{t}} \in (\widehat{\mathcal{A}}^1)^{[1]}$, where the superscript \bar{t} signals that the densities $Y_{\beta, d}^{\bar{t}}$ analytically depend on $\bar{t} \in M$. It is convenient to consider the generating series of densities $Y_{\beta, d}^{\bar{t}}$:

$$Y_{\beta}^{\bar{t}}(z) := \sum_{d \geq -1} Y_{\beta, d}^{\bar{t}} z^{d+1}.$$

Lemma 1. We have $\frac{\partial Y_{\beta}^{\bar{t}}(z)}{\partial t^{\gamma}} = \frac{\partial Y_{\beta}^{\bar{t}}(z)}{\partial u^{\gamma}} - C_{\beta\gamma}^{\mu} z Y_{\mu}^{\bar{t}}(z)$, $1 \leq \beta, \gamma \leq N$.

Proof. This follows from the definition of the densities $Y_{\beta, d}^{\bar{t}}$, the property $\frac{\partial}{\partial t^{\beta}} (c^{\mathcal{X}, \bar{t}})_{g, n+1} = \pi_{1*} \circ (c^{\mathcal{X}, \bar{t}})_{g, n+2} \circ (\otimes e_{\beta})$ (which is equivalent to part 2 of Proposition 4), and the fact that $\pi_1^* \psi_i^d = \psi_i^d - \delta_{i, n+1}^0 \pi_1^* \psi_i^{d-1}$, $1 \leq i \leq n, d \geq 1$, where the class $\delta_{i, n+1}^0$ was defined in the proof of Proposition 1. \square

Consider now a calibration $X_{\alpha}(z)$ of our flat F-manifold M . Define densities $\widetilde{Y}_{\beta, d}^{\bar{t}} \in (\widehat{\mathcal{A}}^1)^{[1]}$, $1 \leq \beta \leq N, d \geq -1$, by $\sum_{d \geq -1} \widetilde{Y}_{\beta, d}^{\bar{t}} z^{d+1} := \widetilde{Y}_{\beta}^{\bar{t}}(z)$ where

$$\widetilde{Y}_{\beta}^{\bar{t}}(z) := Y_{\mu}^{\bar{t}}(z) X_{\beta}^{\mu}(z).$$

Lemma 2. We have $\frac{\partial \widetilde{Y}_{\beta, d}^{\bar{t}}}{\partial t^{\gamma}} = \frac{\partial \widetilde{Y}_{\beta, d}^{\bar{t}}}{\partial u^{\gamma}}$.

Proof. This immediately follows from Lemma 1 and the property $\frac{\partial X_{\beta}^{\mu}(z)}{\partial t^{\gamma}} = C_{\gamma\nu}^{\mu} X_{\beta}^{\nu}(z)$. \square

Define densities of vector fields $Y_{\beta, d}^{\text{desc}} \in (\widehat{\mathcal{A}}_M^1)^{[1]}$, $1 \leq \beta \leq N, d \geq -1$, by

$$Y_{\beta, d}^{\text{desc}} := \left(\widetilde{Y}_{\beta, d}^{\bar{t}} \Big|_{u^*=0} \right) \Big|_{t^{\gamma}=u^{\gamma}} \in (\widehat{\mathcal{A}}_M^1)^{[1]}.$$

The previous lemma implies that for a fixed $\bar{t} \in M$ the density $\widetilde{Y}_{\beta, d}^{\bar{t}}$ is the Taylor expansion of the density $Y_{\beta, d}^{\text{desc}}$ at $u^{\gamma} = t^{\gamma}$, i.e., $\widetilde{Y}_{\beta, d}^{\bar{t}} = Y_{\beta, d}^{\text{desc}} \Big|_{u^{\gamma} \mapsto t^{\gamma} + u^{\gamma}}$, as elements of $(\widehat{\mathcal{A}}^1)^{[1]}$. Therefore, since for any $\bar{t} \in M$ the densities $\widetilde{Y}_{\beta, d}^{\bar{t}}$ produce a hierarchy of pairwise commuting flows, the densities $Y_{\beta, d}^{\text{desc}}$ also produce a hierarchy of pairwise commuting flows. This hierarchy is called the *descendant DR hierarchy*.

In more details, the equations of the descendant DR hierarchy are given by

$$\frac{\partial u^{\alpha}}{\partial t_d^{\beta}} = \partial_x P_{\beta, d}^{\text{desc}; \alpha}, \quad 1 \leq \alpha, \beta \leq N, \quad d \geq 0,$$

where $P_{\beta,d}^{\text{desc};\alpha} = \left(\tilde{P}_{\beta,d}^{\bar{t};\alpha} \Big|_{u^*=0} \right) \Big|_{t^\gamma=u^\gamma}$, $\tilde{P}_{\beta,d}^{\bar{t};\alpha} = \sum_{i=0}^{d+1} P_{\mu,d-i}^{\bar{t};\alpha} X_{i-1,\beta}^\mu$, and $P_{\beta,d}^{\bar{t};\alpha}$ are the differential polynomials (1.13) corresponding to the F-CohFT $c^{\mathcal{X},\bar{t}}$. Also, we adopt the convention $P_{\beta,-1}^{\text{desc};\alpha} := X_{\beta,-1}^\alpha$. Note that we have $\frac{\partial \tilde{P}_{\beta,d}^{\bar{t};\alpha}}{\partial t^\gamma} = \frac{\partial P_{\beta,d}^{\bar{t};\alpha}}{\partial u^\gamma}$.

We immediately see that $P_{\beta,d}^{\text{desc};\alpha} \Big|_{\varepsilon=0} = X_{\beta,d}^\alpha \Big|_{t^\gamma=u^\gamma}$, and therefore the dispersionless part of the descendant DR hierarchy coincides with the principal hierarchy. For $\mathcal{X} = 0$, the descendant DR hierarchy coincides with the principal hierarchy.

Statements analogous to the ones from Theorem 2 are true for the descendant DR hierarchy. We present here the proof of a couple of them.

Note that if $X_{\beta,-1}^\alpha = \delta_\beta^\alpha$, then $X_{\parallel,0}^\alpha$ coincides with t^α up to a constant. We will say that a calibration is of *standard type* if $X_{\beta,-1}^\alpha = \delta_\beta^\alpha$ and $X_{\parallel,0}^\alpha = t^\alpha$.

Proposition 5.1. *We have $\frac{\partial}{\partial u^\parallel} P_{\beta,d+1}^{\text{desc};\alpha} = P_{\beta,d}^{\text{desc};\alpha}$, $1 \leq \alpha, \beta \leq N$, $d \geq -1$.*

2. If our calibration is of standard type, then $\frac{\partial}{\partial u^\beta} P_{\parallel,1}^{\text{desc};\alpha} = DP_{\beta,0}^{\text{desc};\alpha}$, $1 \leq \alpha, \beta \leq N$.

Proof. To prove part 1, we compute $\frac{\partial}{\partial u^\parallel} P_{\beta,d+1}^{\text{desc};\alpha} \Big|_{u^\gamma \mapsto t^\gamma + u^\gamma} = \frac{\partial}{\partial u^\parallel} \tilde{P}_{\beta,d+1}^{\bar{t};\alpha} = \frac{\partial}{\partial u^\parallel} \sum_{i=0}^{d+2} P_{\mu,d+1-i}^{\bar{t};\alpha} X_{i-1,\beta}^\mu = \sum_{i=0}^{d+1} P_{\mu,d-i}^{\bar{t};\alpha} X_{i-1,\beta}^\mu = \tilde{P}_{\beta,d}^{\bar{t};\alpha} = P_{\beta,d}^{\text{desc};\alpha} \Big|_{u^\gamma \mapsto t^\gamma + u^\gamma}$.

For part 2, we compute $\frac{\partial}{\partial u^\beta} P_{\parallel,1}^{\text{desc};\alpha} \Big|_{u^\gamma \mapsto t^\gamma + u^\gamma} = \frac{\partial}{\partial u^\beta} \tilde{P}_{\parallel,1}^{\bar{t};\alpha} = \frac{\partial}{\partial u^\beta} \left(P_{\parallel,1}^{\bar{t};\alpha} + P_{\mu,0}^{\bar{t};\alpha} t^\mu \right) = DP_{\beta,0}^{\bar{t};\alpha} + t^\mu \frac{\partial}{\partial u^\mu} P_{\beta,0}^{\bar{t};\alpha} = DP_{\beta,0}^{\text{desc};\alpha} \Big|_{u^\gamma \mapsto t^\gamma + u^\gamma}$. \square

To summarize the above constructions, given the following data:

- a flat F-manifold structure on $M \subset \mathbb{C}^N$ given by a vector potential \bar{F} such that M is simply connected and all the points of M are semisimple;
- its calibration;
- a framing on M ;

we have constructed a dispersive integrable deformation of the principal hierarchy. In the next section, we will prove that the dispersive deformations corresponding to different framings are not related to each other by a Miura transformation that is close to identity.

2.4. Nonequivalence of dispersive deformations. We say that a Miura transformation (1.6)–(1.7) is *close to identity* if $\tilde{u}^\alpha \Big|_{\varepsilon=0} = u^\alpha$.

Definition 6. Two dispersive deformations of the principal hierarchy of a calibrated flat F-manifold are called *equivalent* if they are related by a Miura transformation that is close to identity.

Theorem 4. *Let us fix a calibrated flat F-manifold structure on a simply connected open subset $M \subset \mathbb{C}^N$, with a vector potential \bar{F} and which is semisimple at each point of M . Then, for different framings \mathcal{X} and $\tilde{\mathcal{X}}$ on M , the corresponding descendant DR hierarchies are not equivalent.*

Proof of Theorem 4. Following [AL18], for a system of evolutionary PDEs of the form

$$\frac{\partial u^\alpha}{\partial t} = Q^\alpha, \quad Q^\alpha \in \tilde{\mathcal{A}}_M^{[1]}, \quad 1 \leq \alpha \leq N,$$

let us consider the associated *Miura matrix* $S(z) = (S_\beta^\alpha(z)) \in \text{Mat}_{N,N}(\mathcal{O}(M)[[z]])$ defined by

$$S_\beta^\alpha(z) := \sum_{d \geq 0} \frac{\partial Q^\alpha}{\partial u_{d+1}^\beta} \Big|_{\substack{u_c^\gamma = \delta_{c,0} t^\gamma \\ \varepsilon = z}}.$$

For a Miura transformation (1.6)–(1.7) that is close to identity, introduce its *symbol* $T(z) = (T_\beta^\alpha(z)) \in \text{Mat}_{N,N}(\mathcal{O}(M)[[z]])$ by

$$T_\beta^\alpha(z) := \sum_{d \geq 0} \frac{\partial \tilde{u}^\alpha}{\partial u_d^\beta} \Big|_{\substack{u_c^\gamma = \delta_{c,0} t^\gamma \\ \varepsilon = z}}.$$

It is easy to see that under the Miura transformation the Miura matrix of our system of PDEs transforms as follows:

$$S(z) \mapsto T(z)S(z)T(z)^{-1}.$$

Now consider the descendant DR hierarchies corresponding to different framings \mathcal{X} and $\widehat{\mathcal{X}}$. Let us denote the Miura matrices of a flow $\frac{\partial}{\partial t^\alpha}$ from these two hierarchies by $S^{(\alpha,d)}(z)$ and $\widehat{S}^{(\alpha,d)}(z)$, respectively. Clearly, $S^{(\alpha,d)}(0) = \widehat{S}^{(\alpha,d)}(0)$. Suppose that the hierarchies are related by a Miura transformation that is close to identity. Denote its symbol by $T(z)$. For the calibration of our flat F-manifold, without loss of generality, we can assume that $X_{\beta,-1}^\alpha = \delta_\beta^\alpha$. Consider the expansions $S^{(\alpha,d)}(z) = \sum_{i \geq 0} S_{2i}^{(\alpha,d)} z^{2i}$, $\widehat{S}^{(\alpha,d)}(z) = \sum_{i \geq 0} \widehat{S}_{2i}^{(\alpha,d)} z^{2i}$, $T(z) = \sum_{i \geq 0} T_i z^i$. Then we have

$$\begin{aligned} \widehat{S}^{(\alpha,d)}(z) = T(z)S^{(\alpha,d)}(z)T(z)^{-1} &\Rightarrow \begin{cases} [T_1, S_0^{(\alpha,d)}] = 0, \\ S_2^{(\alpha,d)} + [S_0^{(\alpha,d)}, T_1]T_1 + [T_2, S_0^{(\alpha,d)}] = \widehat{S}_2^{(\alpha,d)}, \\ \Rightarrow S_2^{(\alpha,d)} - \widehat{S}_2^{(\alpha,d)} = [S_0^{(\alpha,d)}, T_2]. \end{cases} \end{aligned} \quad (2.8)$$

For the descendant DR hierarchy corresponding to the framing \mathcal{X} , we have

$$\widetilde{P}_{\mathbb{1},1}^{\bar{i};\alpha} = P_{\mathbb{1},1}^{\bar{i};\alpha} + P_{\mu,0}^{\bar{i};\alpha} X_{\mathbb{1},0}^\mu + X_{\mathbb{1},1}^\alpha, \quad \widetilde{P}_{\mu,0}^{\bar{i};\alpha} = P_{\mu,0}^{\bar{i};\alpha} + X_{\mu,0}^\alpha,$$

which implies that the matrix $S = (S_\beta^\alpha) := S_2^{(\mathbb{1},1)} - \sum_{\mu=1}^N X_{\mathbb{1},0}^\mu S_2^{(\mu,0)}$ is given by

$$\begin{aligned} S_\beta^\alpha &= \text{Coef}_{\varepsilon^2} \frac{\partial P_{\mathbb{1},0}^{\bar{i};\alpha}}{\partial u_2^\beta} \Big|_{u^* = 0} = \text{Coef}_{a^2} \int_{\text{DR}_1(a,0,-a)} \lambda_1 \psi_2 c_{1,3}^{\mathcal{X},\bar{i}}(e^\alpha \otimes e \otimes e_\beta) \\ &= 2 \text{Coef}_{a^2} \int_{\text{DR}_1(a,-a)} \lambda_1 c_{1,2}^{\mathcal{X},\bar{i}}(e^\alpha \otimes e_\beta), \end{aligned}$$

which is equal to $2 \text{Coef}_{a^2} \int_{\text{DR}_1(a,-a)} \lambda_1 = \frac{1}{12}$ times the degree zero part of $c_{1,2}^{\mathcal{X},\bar{i}}(e^\alpha \otimes e_\beta)$. By the construction of the cohomological field theory $c^{\mathcal{X},\bar{i}}$, the degree zero part of $c_{1,2}^{\mathcal{X},\bar{i}}(e^\alpha \otimes e_\beta)$ is equal to $\sum_{i=1}^N \frac{\partial t^\alpha}{\partial \widehat{u}^i} \mathcal{X}^i \frac{\partial \widehat{u}^i}{\partial \widehat{u}^\beta}$, where $\sum_{i=1}^N \mathcal{X}^i \frac{\partial}{\partial \widehat{u}^i} := \mathcal{X}$ and \widehat{u}^i are local canonical coordinates on M . Since $\mathcal{X} \neq \widehat{\mathcal{X}}$, we conclude that if we denote $\widehat{S} := \widehat{S}_2^{(\mathbb{1},1)} -$

$\sum_{\mu=1}^N X_{\mathbb{1},0}^{\mu} \widehat{S}_2^{(\mu,0)}$, then $\widetilde{\Psi}(S - \widehat{S})\widetilde{\Psi}^{-1}$ is a nonzero diagonal matrix. On the other hand, since $\widetilde{\Psi}S_0^{(\alpha,d)}\widetilde{\Psi}^{-1}$ is a diagonal matrix for any $1 \leq \alpha \leq N$ and $d \geq 0$, the diagonal part of $\left[\widetilde{\Psi}\left(S_0^{(\mathbb{1},1)} - \sum_{\mu=1}^N X_{\mathbb{1},0}^{\mu}S_0^{(\mu,0)}\right)\widetilde{\Psi}^{-1}, \widetilde{\Psi}T_2\widetilde{\Psi}^{-1}\right]$ is equal to zero, which contradicts (2.8). \square

Note that during the proof of the theorem we have obtained the following explicit relation between a framing and the differential polynomials defining the flows $\frac{\partial}{\partial t_1^{\alpha}}$ and $\frac{\partial}{\partial t_0^{\mu}}$ of a corresponding descendant DR hierarchy.

Lemma 3. *Consider a flat F-manifold, a calibration satisfying $X_{\beta,-1}^{\alpha} = \delta_{\beta}^{\alpha}$, a framing $\mathcal{X} = \mathcal{X}^{\alpha} \frac{\partial}{\partial t^{\alpha}}$, and a corresponding descendant DR hierarchy. Then we have*

$$\mathcal{X}^{\alpha} = 12 \frac{\partial}{\partial u_{xx}^{\mathbb{1}}} \text{Coef}_{\varepsilon^2} \left(P_{\mathbb{1},1}^{\text{desc};\alpha} - P_{\beta,0}^{\text{desc};\alpha} X_{\mathbb{1},0}^{\beta} \right) \Big|_{u^*=t^*}.$$

2.5. Homogeneous dispersive deformations. As at the beginning of Sect. 2.3, consider a semisimple flat F-manifold structure on $M \subset \mathbb{C}^N$ defined by a vector potential \overline{F} , a semisimple point, canonical coordinates u^i on an open neighborhood U of this point, the diagonal matrix of one-forms \widetilde{D} , a diagonal nondegenerate matrix H , and matrices R_k . Suppose that our flat F-manifold is homogeneous with an Euler vector field E of the form (2.4). By [ABLR20, Proposition 1.14], the diagonal matrix $i_E \widetilde{D}$ is constant, $i_E \widetilde{D} = -\text{diag}(\delta_1, \dots, \delta_N) = -\Delta$, $\delta_i \in \mathbb{C}$. Moreover, we have $E^{\alpha} \frac{\partial}{\partial t^{\alpha}} H = \Delta H$, and by [ABLR20, Proposition 1.16] we can fix a choice of matrices R_k by the additional conditions $E^{\alpha} \frac{\partial}{\partial t^{\alpha}} R_k = -k R_k + [\Delta, R_k]$ for $k \geq 1$. By [ABLR20, proof of Theorem 4.10], for an arbitrary $1 \leq l \leq N$ and an eigenvector G_0 of the matrix Δ corresponding to the eigenvalue δ_l the family of F-CohFTs $c^{G_0, \bar{t}}$ satisfies the property

$$\text{Deg} \circ c_{g,n+1}^{G_0, \bar{t}} + E^{\alpha} \frac{\partial}{\partial t^{\alpha}} c_{g,n+1}^{G_0, \bar{t}} = c_{g,n+1}^{G_0, \bar{t}} \circ \left(-Q^t \otimes \text{Id}^{\otimes n} + \sum_{i+j=n-1} \text{Id} \otimes \text{Id}^{\otimes i} \otimes Q \otimes \text{Id}^{\otimes j} \right) - 2\delta_l g c_{g,n+1}^{G_0, \bar{t}}.$$

This implies that for any $\bar{t} \in U$ the F-CohFT $c^{G_0, \bar{t}}$ is homogeneous of conformal dimension $-2\delta_l$. Note that the corresponding framing \mathcal{X} on U satisfies the property $[E, \mathcal{X}] = (-2\delta_l - 1)\mathcal{X}$.

Suppose that M is connected, then it is clear that up to permutations of the components the vector $(\delta_1, \dots, \delta_N)$ doesn't depend on a semisimple point. We come to the following natural definition.

Definition 7. The vector $\overline{\gamma} := (-2\delta_1, \dots, -2\delta_N)$ is called the *vector of conformal dimensions* corresponding to our flat F-manifold.

Suppose that all the points of M are semisimple. As in the previous section, we can now glue the local families of F-CohFTs in a global family. Note that given a framing \mathcal{X} on M satisfying $[E, \mathcal{X}] = (-2\delta_l - 1)\mathcal{X}$ we can now construct a unique global family $c^{\mathcal{X}, \bar{t}}, \bar{t} \in M$, of F-CohFTs fixing the choice of matrices R_k using the Euler vector field.

Summarizing the considerations of this section, we obtain the following result.

Theorem 5. Consider a homogeneous flat F -manifold structure on a connected open subset $M \subset \mathbb{C}^N$ defined by a vector potential \bar{F} . Suppose that all the points of M are semisimple. Let $\bar{\gamma} = (\gamma_1, \dots, \gamma_N)$ be the vector of conformal dimensions. Let $1 \leq l \leq N$ and let \mathcal{X} be a framing on M such that $[E, \mathcal{X}] = (\gamma_l - 1)\mathcal{X}$. Then the family of F -CohFTs $c^{\mathcal{X}, \bar{\gamma}}$ satisfies the property

$$\text{Deg} \circ c_{g,n+1}^{\mathcal{X}, \bar{\gamma}} + E^\alpha \frac{\partial}{\partial t^\alpha} c_{g,n+1}^{\mathcal{X}, \bar{\gamma}} = c_{g,n+1}^{\mathcal{X}, \bar{\gamma}} \circ \left(-Q^t \otimes \text{Id}^{\otimes n} + \sum_{i+j=n-1} \text{Id} \otimes \text{Id}^{\otimes i} \otimes Q \otimes \text{Id}^{\otimes j} \right) + \gamma_l g c_{g,n+1}^{\mathcal{X}, \bar{\gamma}}.$$

In particular, for any $\bar{t} \in M$ the F -CohFT $c^{\mathcal{X}, \bar{\gamma}}$ is homogeneous of conformal dimension γ_l .

Let us now discuss properties of the descendant DR hierarchies in the homogeneous case. Under the assumptions of the theorem, suppose also that M is simply connected. By [BB19, Proposition 4.4], there exists a calibration $X_\alpha(z)$ and complex matrices \tilde{R}_i , $i \geq 1$, such that $X_{\beta,-1}^\alpha = \delta_\beta^\alpha$, $[Q, \tilde{R}_i] = i\tilde{R}_i$, and

$$E^\mu \frac{\partial}{\partial t^\mu} X(z) = z \frac{\partial}{\partial z} X(z) + [X(z), Q] + X(z)\tilde{R}(z),$$

where $X(z) := \sum_{d \geq -1} (X_{\beta,d}^\alpha) z^{d+1}$ and $\tilde{R}(z) := \sum_{i \geq 1} \tilde{R}_i z^i$. Such a calibration is called *homogeneous*. Consider now the associated descendant DR hierarchy.

Let us introduce a generating series $P^{\text{desc}}(z)$ by

$$P^{\text{desc}}(z) := \sum_{d \geq -1} (P_{\beta,d}^{\text{desc};\alpha}) z^{d+1}.$$

Proposition 6. We have

$$\widehat{E}_\gamma P^{\text{desc}}(z) = z \frac{\partial}{\partial z} P^{\text{desc}}(z) + [P^{\text{desc}}(z), Q] + P^{\text{desc}}(z)\tilde{R}(z).$$

Proof. Let us introduce generating series $\tilde{P}^{\bar{t}}(z)$ and $P^{\bar{t}}(z)$ by

$$\tilde{P}^{\bar{t}}(z) := \sum_{d \geq -1} (\tilde{P}_{\beta,d}^{\bar{t};\alpha}) z^{d+1}, \quad P^{\bar{t}}(z) := \sum_{d \geq -1} (P_{\beta,d}^{\bar{t};\alpha}) z^{d+1}.$$

We have to check that

$$\widehat{E}_\gamma P^{\text{desc}}(z) \Big|_{u^\nu \mapsto t^\nu + u^\nu} = z \frac{\partial}{\partial z} \tilde{P}^{\bar{t}}(z) + [\tilde{P}^{\bar{t}}(z), Q] + \tilde{P}^{\bar{t}}(z)\tilde{R}(z).$$

For this, we compute

$$\begin{aligned} \widehat{E}_\gamma P^{\text{desc}}(z) \Big|_{u^\nu \mapsto t^\nu + u^\nu} &= \left(\widehat{E}_\gamma + t^\alpha (\delta_\alpha^\beta - Q_\alpha^\beta) \frac{\partial}{\partial u^\beta} \right) \tilde{P}^{\bar{t}}(z) = \\ &= \left(\widehat{E}_\gamma + t^\alpha (\delta_\alpha^\beta - Q_\alpha^\beta) \frac{\partial}{\partial u^\beta} \right) P^{\bar{t}}(z) \cdot X(z) \stackrel{\text{by Prop. 1.7}}{=} \end{aligned}$$

$$= z \frac{\partial}{\partial z} P^{\bar{i}}(z) \cdot X(z) + [P^{\bar{i}}(z), Q]X(z) + z P^{\bar{i}}(z) (E^\gamma C_\gamma) X(z), \quad (2.9)$$

where $C_\gamma := (C_{\gamma\beta}^\alpha)$, and we recall that $C_{\beta\gamma}^\alpha = \frac{\partial^2 F^\alpha}{\partial t^\beta \partial t^\gamma}$. Since $z E^\gamma C_\gamma X(z) = E^\gamma \frac{\partial}{\partial t^\gamma} X(z) = z \frac{\partial}{\partial z} X(z) + [X(z), Q] + X(z) \tilde{R}(z)$, the expression in line (2.9) is equal to

$$\begin{aligned} & z \frac{\partial}{\partial z} P^{\bar{i}}(z) \cdot X(z) + [P^{\bar{i}}(z), Q]X(z) + P^{\bar{i}}(z) \left(z \frac{\partial}{\partial z} X(z) + [X(z), Q] + X(z) \tilde{R}(z) \right) \\ &= z \frac{\partial}{\partial z} \tilde{P}^{\bar{i}}(z) + [\tilde{P}^{\bar{i}}(z), Q] + \tilde{P}^{\bar{i}}(z) \tilde{R}(z), \end{aligned}$$

as required. \square

3. Towards a Classification of Dispersive Deformations

In this section we consider the problem of classification of dispersive integrable deformations of principal hierarchies for flat F-manifolds and observe the central role played in it by the DR hierarchies. We propose two a priori different classes of deformations and we classify them, up to some finite order in ε , for 1 and 2 dimensional flat F-manifolds, respectively. Up to that approximation, we observe that both classes contain essentially the DR hierarchies considered in Sect. 2.

3.1. Dispersive deformations of DR type and the rank 1 case.

3.1.1. Integrable systems of DR type Given a local vector field $\bar{X} \in (\widehat{\mathcal{A}}^1)^{[1]}$, consider the operator $\mathcal{D}_{\bar{X}}: \widehat{\mathcal{A}}^1[[z]] \rightarrow \widehat{\mathcal{A}}^1[[z]]$ defined by

$$\begin{aligned} \mathcal{D}_{\bar{X}} Y(z) &:= \partial_x (D - 1) Y(z) - z [\bar{X}, Y(z)], \\ Y(z) &= \sum_{k \geq 0} Y_{k-1} z^k, \quad Y_{k-1} \in \widehat{\mathcal{A}}^1. \end{aligned}$$

Suppose there exist N solutions $Y_\alpha(z) \in (\widehat{\mathcal{A}}^1)^{[1]}$, $1 \leq \alpha \leq N$, to the equation $\mathcal{D}_{\bar{X}} Y_\alpha(z) = 0$ with the initial conditions $Y_\alpha(z=0) = -\theta_{\alpha,1}$. Then a new vector of solutions with the same initial conditions can be found by the following transformation:

$$Y_\alpha(z) \mapsto a_\alpha^\mu(z) Y_\mu(z), \quad (3.1)$$

where $a_\alpha^\mu(z) = \delta_\alpha^\mu + \sum_{i>0} a_{\alpha,i}^\mu z^i \in \mathbb{C}[[z]]$.

Theorem 6. *Assume that $\bar{X} \in (\widehat{\mathcal{A}}^1)^{[1]}$ satisfies the following properties:*

(a) *there exist N solutions $Y_\alpha(z) = \sum_{d \geq 0} Y_{\alpha,d-1} z^d \in (\widehat{\mathcal{A}}^1)^{[1]}[[z]]$, $1 \leq \alpha \leq N$, to the equation*

$$\mathcal{D}_{\bar{X}} Y_\alpha(z) = 0 \quad (3.2)$$

with the initial conditions $Y_\alpha(z=0) = -\theta_{\alpha,1}$,

(b) $\frac{\delta}{\delta u^{\mathbb{I}}} \bar{X} = -u^\alpha \theta_{\alpha,1} + \partial_x^2 R$, $R \in (\widehat{\Lambda}^1)^{[-1]}$, where $\frac{\delta}{\delta u^{\mathbb{I}}} = A^\alpha \frac{\delta}{\delta u^\alpha}$ and A^α are some complex constants.

Then, up to a transformation of type (3.1), we have

- (i) $Y_{\mathbb{I},0} = -u^\alpha \theta_{\alpha,1} + \partial_x^2 (D-1)^{-1} R$,
- (ii) $\bar{Y}_{\mathbb{I},1} = \bar{X}$,
- (iii) $[\bar{Y}_{\alpha_1, d_1}, \bar{Y}_{\alpha_2, d_2}] = 0$, $1 \leq \alpha_1, \alpha_2 \leq N$, $d_1, d_2 \geq -1$,
- (iv) $[\bar{Y}_{\alpha_2, 0}, Y_{\alpha_1, d}] = \partial_x \frac{\partial}{\partial u^{\alpha_2}} Y_{\alpha_1, d+1}$, $1 \leq \alpha_1, \alpha_2 \leq N$, $d \geq -1$,
- (v) $\frac{\partial}{\partial u^{\mathbb{I}}} Y_{\alpha, d+1} = Y_{\alpha, d}$, $1 \leq \alpha \leq N$, $d \geq -1$.

Proof. The proof follows closely the proof of [BDGR20, Theorem 5.1–5.2] with Lie brackets of densities of local vector fields replacing Poisson brackets of differential polynomials. \square

Remark 4. When we restrict to $\varepsilon = 0$, a particular local vector field satisfying condition (a) of Theorem 6 is given by $\bar{X} = -(D-2) \int F^\alpha(u^1, \dots, u^N) \theta_{\alpha,1} dx$ where the functions $F^\alpha(t^1, \dots, t^N)$ are solutions to the oriented WDVV Eqs. (2.2), (2.3). It is easy to check that for such \bar{X} solutions $Y_\alpha(z)$ are given by $Y_\alpha(z) = -\sum_{d \geq -1} X_{\alpha,d}^\beta \theta_{\beta,1} z^{d+1}$ where the functions $X_{\alpha,d}^\beta$ form a calibration of the flat F-manifold satisfying $X_{\alpha,-1}^\beta = \delta_\alpha^\beta$ (see Sect. 2.2). Therefore, the functions $Y_\alpha(z)$ are the generating series of densities of local vector fields of the principal hierarchy of the flat F-manifold. Note that condition (b) for our \bar{X} is equivalent to $\frac{\partial F^\alpha}{\partial t^{\mathbb{I}}} = t^\alpha$, which can always be fulfilled by adding to F^α appropriate linear terms.

Definition 8. Let $\bar{X} \in (\widehat{\Lambda}^1)^{[1]}$ satisfy the hypotheses of Theorem 6. Then we say that $\bar{X} = \bar{Y}_{\mathbb{I},1}$ and the induced hierarchy of compatible densities of local vector fields $Y_{\alpha,d}$, $1 \leq \alpha \leq N$, $d \geq -1$, are of *double ramification (DR) type*.

Theorem 7. *The double ramification hierarchy (1.11) associated to an F-CohFT is a hierarchy of double ramification type.*

Proof. Hypotheses (a) and (b) of Theorem 6 follow immediately from claims (iii) and (v), respectively, of Theorem 2. \square

3.1.2. Classification of rank 1 hierarchies of DR type Thanks to Theorem 6 and Remark 4, it makes sense to use Eq. (3.2) to find all possible deformations of DR type of a principal hierarchy associated to a given flat F-manifold, at low order in the dispersion parameter ε . These deformations will, in particular, include the ones coming from all F-CohFTs with the given genus 0 part.

Consider the ancestor principal hierarchy associated to the genus 0 part of the trivial CohFT, i.e., the CohFT with $V = \mathbb{C}\langle e \rangle$ and $c_{g,n}(e^{\otimes n}) = 1 \in H^0(\overline{\mathcal{M}}_{g,n})$ for all (g, n) in the stable range. Let $e_1 = e$, $u_k := u_k^1$, $\theta_k := \theta_{1,k}$ for $k \geq 0$, with $u := u_0$, $\theta := \theta_1$ as usual, and $Y_d := Y_{1,d} = Y_{\mathbb{I},d}$. A direct computation (at the approximation up to ε^9) shows that its most general deformation of DR type is either of the form

$$\frac{\partial u}{\partial t_1} = \frac{\delta \bar{Y}_1}{\delta \theta} = uu_1 + \varepsilon^2 C_{1,1} u_3 + \varepsilon^4 (C_{2,1} u_5 + C_{2,2} u_2 u_3)$$

$$\begin{aligned}
 & + \varepsilon^6 \left(C_{3,1} u_2 u_5 + C_{3,2} u_3 (u_2)^2 + \left(\frac{10(C_{2,1})^2}{7C_{1,1}} + \frac{9}{35} C_{1,1} C_{2,2} \right) u_7 + \left(-\frac{8C_{2,1}C_{2,2}}{3C_{1,1}} + 3C_{3,1} \right) u_3 u_4 \right) \\
 & + \varepsilon^8 \left(C_{4,1} (u_2)^2 u_5 + C_{4,2} (u_2)^3 u_3 + \left(\frac{15(C_{2,1})^3}{7(C_{1,1})^2} + \frac{58}{105} C_{2,1} C_{2,2} + \frac{4}{35} C_{1,1} C_{3,1} \right) u_9 \right. \\
 & \quad + \left(-\frac{123(C_{2,1})^2 C_{2,2}}{28(C_{1,1})^2} + \frac{57}{100} (C_{2,2})^2 + \frac{9}{16} C_{1,1} C_{3,2} + \frac{33C_{2,1}C_{3,1}}{7C_{1,1}} \right) u_2 u_7 \\
 & \quad + \left(-\frac{177(C_{2,1})^2 C_{2,2}}{4(C_{1,1})^2} + \frac{201}{100} (C_{2,2})^2 + \frac{333}{80} C_{1,1} C_{3,2} + \frac{33C_{2,1}C_{3,1}}{C_{1,1}} \right) u_4 u_5 \\
 & \quad + \left(\frac{44C_{2,1}(C_{2,2})^2}{21(C_{1,1})^2} - \frac{55C_{2,1}C_{3,2}}{12C_{1,1}} - \frac{44C_{2,2}C_{3,1}}{21C_{1,1}} + 2C_{4,1} \right) (u_3)^3 \\
 & \quad + \left(-\frac{24(C_{2,1})^2 C_{2,2}}{(C_{1,1})^2} + \frac{249}{175} (C_{2,2})^2 + \frac{9}{4} C_{1,1} C_{3,2} + \frac{132C_{2,1}C_{3,1}}{7C_{1,1}} \right) u_3 u_6 \\
 & \quad + \left(\frac{88C_{2,1}(C_{2,2})^2}{21(C_{1,1})^2} - \frac{55C_{2,1}C_{3,2}}{6C_{1,1}} - \frac{88C_{2,2}C_{3,1}}{21C_{1,1}} + 6C_{4,1} \right) u_2 u_3 u_4 \Big) \\
 & + O(\varepsilon^{10}),
 \end{aligned} \tag{3.3}$$

with $C_{i,j} \in \mathbb{C}$ and $C_{1,1} \neq 0$, or of the form

$$\frac{\partial u}{\partial t_1} = \frac{\delta \bar{Y}_1}{\delta \theta} = uu_1 + \varepsilon C u_2, \quad C \in \mathbb{C}. \tag{3.4}$$

Notice that, imposing $C_{k,2} = 0$ for all $k \geq 1$ in Eq. (3.3), we recover the most general Hamiltonian deformation of DR type, obtained in [BDGR20], which is in turn in one to one correspondence with the most general rank 1 CohFT. This shows that the extra parameters $C_{k,2}$, $k \geq 1$ control the strictly non-Hamiltonian deformations (at least with respect to the Hamiltonian operator ∂_x). We expect these to correspond to F-CohFTs that are not CohFTs.

Remark 5. It is easy to check that the r.h.s of Eq. (3.3) is a total x -derivative. Comparing with the results of [ALM15] we see that a similar result can be obtained starting from generic scalar conservation laws of the form

$$\begin{aligned}
 \frac{\partial u}{\partial t_d} &= \partial_x P_d, \quad d \geq 0, \\
 P_d &= \sum_{l \geq 0} \varepsilon^{2l} P_{d,l}, \quad P_{d,l} \in \mathcal{A}_M^{[2l]},
 \end{aligned} \tag{3.5}$$

choosing

- $P_{d,0} = \frac{u^{d+1}}{(d+1)!}$,
- $\frac{\partial P_l}{\partial u_x} = 0$; the reduction to this form by means of a Miura transformation is always possible and it is unique,

and imposing the following conditions:

- *Commutativity of the flows.*
- *String property:* $\frac{\partial}{\partial u} P_{d+1} = P_d$ for $d \geq -1$, where $P_{-1} := 1$.

According to the conjecture formulated in [ALM15], it should be possible to write all the coefficients appearing in the deformation as functions of the coefficients of the quasilinear part. Moreover, the coefficients of the quasilinear part should be constant (due to the string property) and arbitrary. This is consistent with the formula (3.3) since the additional free parameters appearing at the order ε^8 are related to the coefficients of the quasilinear part by constraints obtained considering higher order conditions.

Even more intriguing is the isolated deformation (3.4), which, up to reabsorbing the constant C into the factor ε , is the celebrated Burgers equation, which is dissipative and hence non-Hamiltonian. The appearance of terms with odd powers of ε in a hierarchy of DR type rules out the possibility that it is the double ramification hierarchy of an F-CohFT. However, considering that flat F-manifolds are known to appear in genus 0 open Gromov–Witten and Saito theory [PST14, BCT18, BCT19, BB19] it is tempting to conjecture that Burgers equation (3.4) and its higher symmetries might control some version of F-CohFT on the space of Riemann surfaces with boundaries, where curves can indeed possess half-integer genus accounting for odd powers of the genus parameter ε .

The fact that Burgers equation (3.4) and its higher symmetries form a hierarchy of DR type can be proved rigorously at all orders in ε as follows.

Theorem 8. *The vector field $\bar{X} = \int (uu_x + \varepsilon u_{xx})\theta dx$ of the Burgers equation defines a hierarchy of DR type, i.e., it satisfies conditions (a) and (b) of Theorem 6.*

Proof. Let us first present a reformulation of the Schouten–Nijenhuis bracket $[\cdot, \cdot]: \widehat{\mathcal{A}}^1 \times \widehat{\mathcal{A}}^1 \rightarrow \widehat{\mathcal{A}}^1$ in terms of formal differential operators. Consider an arbitrary local vector field $\bar{X} = \int X\theta dx \in \widehat{\mathcal{A}}^1$ and a density $Y = \sum_{k \geq 0} Y_k \theta_k \in \widehat{\mathcal{A}}^1$. The local vector field \bar{X} defines a flow on the space of differential polynomials by

$$\frac{\partial u}{\partial t} = \frac{\delta \bar{X}}{\delta \theta} = X,$$

and we consider also formal differential operators $\widetilde{L}_{\bar{X}}$ and L_Y defined by

$$\widetilde{L}_{\bar{X}} := \sum_{k \geq 0} (-\partial_x)^k \circ \frac{\partial X}{\partial u_k}, \quad L_Y := \sum_{k \geq 0} Y_k \partial_x^k.$$

Directly from the definition (1.2), we obtain the following identity:

$$L_{[\bar{X}, Y]} = \frac{\partial}{\partial t} L_Y - L_Y \circ \widetilde{L}_{\bar{X}},$$

where we apply the differentiation $\frac{\partial}{\partial t}$ to the operator L_Y coefficient-wise.

Let us now take $\bar{X} = \int (uu_x + \varepsilon u_{xx})\theta dx$.

Let us prove condition (a) of Theorem 6 by showing that a required solution $Y(z) = \sum_{k \geq -1} Y_k z^{k+1}$ of Eq. (3.2) is given by

$$Y(z) = -e^{z\varepsilon\partial_x} e^{z(u-2\varepsilon\partial_x)} \theta_1 \Leftrightarrow L_{Y(z)} = \sum_{k \geq -1} z^{k+1} L_{Y_k} = -e^{z\varepsilon\partial_x} \circ e^{z(u-2\varepsilon\partial_x)} \circ \partial_x.$$

Since $\widetilde{L}_{\bar{X}} = -u\partial_x + \varepsilon\partial_x^2$, Eq. (3.2) is equivalent to

$$\partial_x \circ \widetilde{D} L_{Y(z)} = z \left(\frac{\partial L_{Y(z)}}{\partial t} - L_{Y(z)} \circ (-u\partial_x + \varepsilon\partial_x^2) \right), \quad (3.6)$$

where $\tilde{D} := \sum_{n \geq 0} u_n \frac{\partial}{\partial u_n} + \varepsilon \frac{\partial}{\partial \varepsilon}$, and we apply \tilde{D} to $L_{Y(z)}$ coefficient-wise. Note that $\tilde{D}L_{Y(z)} = z \frac{\partial}{\partial z} L_{Y(z)}$. Therefore, Eq. (3.6) is equivalent to

$$\begin{aligned} -\partial_x \circ \frac{\partial}{\partial z} \left(e^{z\varepsilon\partial_x} \circ e^{z(u-2\varepsilon\partial_x)} \circ \partial_x \right) &= -\frac{\partial}{\partial t} \left(e^{z\varepsilon\partial_x} \circ e^{z(u-2\varepsilon\partial_x)} \circ \partial_x \right) \\ &\quad + e^{z\varepsilon\partial_x} \circ e^{z(u-2\varepsilon\partial_x)} \circ \partial_x \circ (-u\partial_x + \varepsilon\partial_x^2) \Leftrightarrow \\ \Leftrightarrow e^{z\varepsilon\partial_x} \circ (-\partial_x \circ u + \varepsilon\partial_x^2) \circ e^{z(u-2\varepsilon\partial_x)} \circ \partial_x &= -e^{z\varepsilon\partial_x} \circ \frac{\partial}{\partial t} e^{z(u-2\varepsilon\partial_x)} \circ \partial_x \\ &\quad + e^{z\varepsilon\partial_x} \circ e^{z(u-2\varepsilon\partial_x)} \circ (-\partial_x \circ u + \varepsilon\partial_x^2) \circ \partial_x \Leftrightarrow \\ \Leftrightarrow \frac{\partial}{\partial t} e^{z(u-2\varepsilon\partial_x)} &= [e^{z(u-2\varepsilon\partial_x)}, -\partial_x \circ u + \varepsilon\partial_x^2]. \end{aligned}$$

Note that the last equation follows from the elementary identity $\frac{\partial}{\partial t}(u - 2\varepsilon\partial_x) = [u - 2\varepsilon\partial_x, -\partial_x \circ u + \varepsilon\partial_x^2]$.

Condition (b) of Theorem 6 immediately follows from the equation $\frac{\delta}{\delta u} \bar{X} = -u\theta_1 + \varepsilon\theta_2$. \square

3.2. Homogeneous dispersive deformations and the rank 2 case.

3.2.1. Homogeneous deformations with string and dilaton property Let us fix a homogeneous flat F-manifold structure on an open subset $M \subset \mathbb{C}^N$ given by a vector potential \bar{F} , together with a homogeneous calibration of standard type. We consider systems of evolutionary PDEs of the form

$$\begin{aligned} \frac{\partial u^\alpha}{\partial t_d^\beta} &= \partial_x P_{\beta,d}^\alpha, \quad 1 \leq \alpha, \beta \leq N, \quad d \geq 0, \\ P_{\beta,d}^\alpha &= \sum_{l \geq 0} \varepsilon^{2l} P_{\beta,d,l}^\alpha, \quad P_{\beta,d,l}^\alpha \in \mathcal{A}_M^{[2l]}, \end{aligned} \quad (3.7)$$

such that the following properties are satisfied:

- (1) *Commutativity of the flows*: the flows $\frac{\partial}{\partial t_d^\beta}$ pairwise commute,
- (2) The dispersionless limit of the system (3.7) coincides with the principal hierarchy of the given calibrated flat F-manifold,
- (3) *String property*: $\frac{\partial}{\partial u^\alpha} P_{\beta,d+1}^\alpha = P_{\beta,d}^\alpha$ for $d \geq -1$, where $P_{\beta,-1}^\alpha := \delta_{\beta,\alpha}^\alpha$,
- (4) *Dilaton property*: $\frac{\partial P_{\beta,1}^\alpha}{\partial u^\beta} = D P_{\beta,0}^\alpha$,
- (5) *Homogeneity condition*: $\widehat{E}_\gamma P(z) = z \frac{\partial}{\partial z} P(z) + [P(z), Q] + P(z) \widetilde{R}(z)$ for some $\gamma \in \mathbb{C}$, where $P(z) := \sum_{d \geq -1} (P_{\beta,d}^\alpha) z^{d+1}$.

In this section, working out the $N = 2$ case, we observe how descendant DR hierarchies appear in the problem of classification of dispersive integrable deformations of principal hierarchies of flat F-manifolds of the above form, which we refer to as a *homogeneous deformation with string and dilaton properties*. The role played by conditions (3), (4), and (5) is central in producing finite dimensional spaces of deformations even without having to quotient with respect to equivalence up to Miura transformations of the dependent variables.

Remark 6. Axioms (1), (3), and (4) above correspond closely to properties (iii), (iv), and (v) of Theorem 6 for hierarchies of local vector fields of DR type. Homogeneity (5) corresponds to property (iii) of Proposition 1 for homogeneous DR hierarchies. Finally, condition (2) above is satisfied by hierarchies of DR type, see Remark 4. This means that homogeneous dispersive deformations with string and dilaton properties contain homogeneous descendant hierarchies of DR type whose local vector fields have only even powers of ε . It's not a priori clear that the converse is true and it would be interesting to investigate this point.

3.2.2. Classification of semisimple homogeneous flat F-manifolds in dimension 2 In the semisimple case, using canonical coordinates u_1, \dots, u_N , the structure of a homogeneous flat F-manifold can be recovered from a solution of the following system [AL19]:

$$\begin{aligned} \frac{\partial \Gamma_{ij}^i}{\partial u_k} &= -\Gamma_{ij}^i \Gamma_{ik}^i + \Gamma_{ij}^i \Gamma_{jk}^j + \Gamma_{ik}^i \Gamma_{kj}^k, & i \neq k \neq j \neq i, \\ \sum_{k=1}^N \frac{\partial \Gamma_{ij}^i}{\partial u_k} &= 0, \quad \sum_{k=1}^N u_k \frac{\partial \Gamma_{ij}^i}{\partial u_k} = -\Gamma_{ij}^i, & i \neq j. \end{aligned} \quad (3.8)$$

For $N = 2$, the above system reduces to (3.8), and the general solution is

$$\Gamma_{ij}^i = \frac{\epsilon_j}{u_i - u_j},$$

where ϵ_1 and ϵ_2 are arbitrary constants. Note that the corresponding vector of conformal dimensions is equal to $(2\epsilon_2, 2\epsilon_1)$. In order to compute a vector potential, we need to introduce flat coordinates u, v (these correspond to t^1, t^2 in Sect. 2.1). We have to distinguish 3 cases:

I. $\epsilon_1 + \epsilon_2 \neq 0, 1$. In this case, flat coordinates are

$$u = \left(\frac{u_1 - u_2}{4} \right)^{\frac{1}{m}}, \quad v = \frac{2+c}{4} u_1 + \frac{2-c}{4} u_2,$$

where $c = 2 \frac{\epsilon_1 - \epsilon_2}{\epsilon_1 + \epsilon_2}$, $m = \frac{1}{1 - \epsilon_1 - \epsilon_2} \neq 0, 1$, and a vector potential is

$$(F^1, F^2) = \begin{cases} \left(uv - 2c \frac{u^{m+1}}{m+1}, \frac{v^2}{2} + \frac{4-c^2}{2} \frac{mu^{2m}}{2m-1} \right), & \text{if } m \neq -1, \frac{1}{2}, 0, 1, \\ \left(uv - 2c \log u, \frac{v^2}{2} + \frac{4-c^2}{6} u^{-2} \right), & \text{if } m = -1, \\ \left(uv - \frac{4}{3} cu^{3/2}, \frac{v^2}{2} + \frac{4-c^2}{4} u \log u \right), & \text{if } m = \frac{1}{2}. \end{cases}$$

The unit is $\frac{\partial}{\partial v}$, the Euler vector field is $E = \frac{1}{m} u \frac{\partial}{\partial u} + v \frac{\partial}{\partial v}$, and $\bar{\gamma} = \left(\frac{(2-c)(m-1)}{2m}, \frac{(2+c)(m-1)}{2m} \right)$.

If m is a half-integer, these are the vector potentials of the bi-flat F-manifold structures defined on the orbit space of the dihedral group $I_2(2m)$ [AL17]. If also $c = 0$, the above vector potential comes from the Dubrovin–Frobenius manifold structure defined on the orbit space of the dihedral group.

II. $\epsilon_1 = c$, $\epsilon_2 = 1 - c$, $c \neq 0$ (see the remark about the case $c = 0$ below). Using the flat coordinates

$$u = u_1 - u_2 + \frac{u_2}{c}, \quad v = -\ln(u_1 - u_2),$$

we obtain

$$F^1 = \frac{c}{2}u^2 + \frac{1-c}{4}e^{-2v}, \quad F^2 = cuv + (2c-1)e^{-v}. \quad (3.9)$$

The unit is $\frac{1}{c}\frac{\partial}{\partial u}$, the Euler vector field is $E = u\frac{\partial}{\partial u} - \frac{\partial}{\partial v}$, and $\bar{\gamma} = (2 - 2c, 2c)$. For $c = \frac{1}{2}$, the above vector potential comes from the genus 0 Gromov–Witten potential of the complex projective line.

In the case $c = 0$, choosing the flat coordinates $u = u_2$ and $v = -\ln(u_1 - u_2)$, we obtain $F^1 = \frac{u^2}{2}$ and $F^2 = uv - e^{-v}$. This flat F-manifold is isomorphic to the flat F-manifold (3.9) with $c = 1$ and shifted by $v \mapsto v + \pi i$.

III. $\epsilon_1 = c$, $\epsilon_2 = -c$. If $c \neq 0$, then using the flat coordinates

$$\bar{u} = u_1 - u_2, \quad v = (u_1 - u_2) \ln(u_1 - u_2) + \frac{u_2}{c},$$

we obtain

$$F^1 = cuv + u^2 \left(\frac{c+1}{2} - c \ln u \right),$$

$$F^2 = \frac{c}{2}v^2 + u^2 \left(-\frac{3c+1}{4} + \frac{c+1}{2} \ln u - \frac{c}{2}(\ln u)^2 \right).$$

The unit is $\frac{1}{c}\frac{\partial}{\partial v}$, the Euler vector field is $E = u\frac{\partial}{\partial u} + (u+v)\frac{\partial}{\partial v}$, and $\bar{\gamma} = (-2c, 2c)$. If $c = 0$, then choosing as flat coordinates the canonical coordinates $u = u_1$ and $v = u_2$ we obtain

$$F^1 = \frac{u^2}{2}, \quad F^2 = \frac{v^2}{2}.$$

The unit is $\frac{\partial}{\partial u} + \frac{\partial}{\partial v}$, the Euler vector field is $E = u\frac{\partial}{\partial u} + v\frac{\partial}{\partial v}$, and $\bar{\gamma} = (0, 0)$.

3.2.3. Integrable deformations of rank 2 homogeneous principal hierarchies We now want to classify all homogeneous deformations with string and dilaton properties of principal hierarchies associated to the homogeneous two-dimensional flat F-manifolds considered above. In our computations below, we have observed the following remarkable facts:

- If such a deformation exists and is nontrivial at the ε^2 approximation, then γ **must be equal** to γ_1 or γ_2 .
- For $\gamma = \gamma_i$, at the ε^2 approximation, any such deformation coincides with the descendant DR hierarchy constructed using an appropriate framing. In particular, any such deformation at the approximation up to ε^2 can be extended to a deformation at all orders of ε .

Let us consider all three cases from Sect. 3.2.2 in detail.

Case I. For simplicity, we consider the case $\frac{1}{m} \neq \mathbb{Z}$, which guarantees that there is a unique homogeneous calibration of standard type such that $\tilde{R}_i = 0$ for $i \geq 1$. Recall that the vector of conformal dimensions is $(\gamma_1, \gamma_2) = \left(\frac{(2-c)(m-1)}{2m}, \frac{(2+c)(m-1)}{2m} \right)$. We have three subcases.

Case II. If $\gamma_1 \neq \gamma_2$ and $\gamma = \gamma_1$, we obtain

$$\begin{aligned} P_{1,0}^1 &= v - 2cu^m \\ &\quad + Au^{-\frac{1}{2}c(m-1)}\varepsilon^2 \left(m(c-2)(cm-c-2m+4)u^{m-3}u_x^2 + m(c-2)^2u^{m-2}u_{xx} - cu^{-1}v_{xx} \right) + \mathcal{O}(\varepsilon^4), \\ P_{1,0}^2 &= \frac{m^2(4-c^2)}{2m-1}u^{2m-1} + Am(c^2-4)u^{-\frac{1}{2}c(m-1)}\varepsilon^2 \left(m(cm-c-4m+6)u^{2m-4}u_x^2 \right. \\ &\quad \left. + m(c-4)u^{2m-3}u_{xx} - u^{m-2}v_{xx} \right) + \mathcal{O}(\varepsilon^4), \end{aligned}$$

and

$$\begin{aligned} P_{2,1}^1 &= uv - \frac{2mc}{m+1}u^{m+1} + Au^{-\frac{1}{2}c(m-1)}\varepsilon^2 \left(m(c-2)(cm-c-2m-2)u^{m-2}u_x^2 \right. \\ &\quad \left. + \frac{m(c-2)(cm-c-2m-2)}{m-1}u^{m-1}u_{xx} - \frac{cm-c-4}{m-1}v_{xx} \right) + \mathcal{O}(\varepsilon^4), \\ P_{2,1}^2 &= \frac{v^2}{2} + \frac{m(4-c^2)}{2}u^{2m} + Am(c+2)u^{-\frac{1}{2}c(m-1)}\varepsilon^2 \left(m(cm-c-4m)(c-2)u^{2m-3}u_x^2 \right. \\ &\quad \left. + \frac{m(cm-c-4m)(c-2)}{m-1}u^{2m-2}u_{xx} - \frac{cm-c-2m-2}{m-1}u^{m-1}v_{xx} \right) + \mathcal{O}(\varepsilon^4). \end{aligned}$$

Here A is an arbitrary complex constant. This deformation is given by the descendant DR hierarchy corresponding to the framing $(\mathcal{X}^1, \mathcal{X}^2) = 12A \left(\frac{4}{m-1}u^{-\frac{1}{2}c(m-1)}, \frac{4m(c+2)}{m-1}u^{-\frac{1}{2}(c-2)(m-1)} \right)$.

Case I2. If $\gamma_1 \neq \gamma_2$ and $\gamma = \gamma_2$, we obtain

$$\begin{aligned} P_{1,0}^1 &= v - 2cu^m \\ &\quad + Bu^{\frac{1}{2}c(m-1)}\varepsilon^2 \left(m(c+2)(cm-c+2m-4)u^{m-3}u_x^2 + m(c+2)^2u^{m-2}u_{xx} - cu^{-1}v_{xx} \right) + \mathcal{O}(\varepsilon^4), \\ P_{1,0}^2 &= \frac{m^2(4-c^2)}{2m-1}u^{2m-1} + Bm(c^2-4)u^{\frac{1}{2}c(m-1)}\varepsilon^2 \left(m(cm-c+4m-6)u^{2m-4}u_x^2 \right. \\ &\quad \left. + m(c+4)u^{2m-3}u_{xx} - u^{m-2}v_{xx} \right) + \mathcal{O}(\varepsilon^4), \end{aligned}$$

and

$$\begin{aligned} P_{2,1}^1 &= uv - \frac{2mc}{m+1}u^{m+1} + Bu^{\frac{1}{2}c(m-1)}\varepsilon^2 \left(m(c+2)(cm-c+2m+2)u^{m-2}u_x^2 \right. \\ &\quad \left. + \frac{m(c+2)(cm-c+2m+2)}{m-1}u^{m-1}u_{xx} - \frac{cm-c+4}{m-1}v_{xx} \right) + \mathcal{O}(\varepsilon^4), \\ P_{2,1}^2 &= \frac{v^2}{2} + \frac{m(4-c^2)}{2}u^{2m} + Bm(c-2)u^{\frac{1}{2}c(m-1)}\varepsilon^2 \left(m(cm-c+4m)(c+2)u^{2m-3}u_x^2 \right. \\ &\quad \left. + \frac{m(cm-c+4m)(c+2)}{m-1}u^{2m-2}u_{xx} - \frac{cm-c+2m+2}{m-1}u^{m-1}v_{xx} \right) + \mathcal{O}(\varepsilon^4). \end{aligned}$$

Here B is an arbitrary complex constant. This deformation is given by the descendant DR hierarchy corresponding to the framing $(\mathcal{X}^1, \mathcal{X}^2) = 12B \left(-\frac{4}{m-1} u^{\frac{1}{2}c(m-1)}, -\frac{4m(c-2)}{m-1} u^{\frac{1}{2}(c+2)(m-1)} \right)$.

Case I3. If $\gamma_1 = \gamma_2$ (which is equivalent to $c = 0$) and γ coincides with them, we get a two-parameter family of deformations formed by linear combinations of the deformations from Cases I1 and I2.

Case II. There is a unique homogeneous calibration of standard type such that $X_{\beta,0}^\alpha = \frac{\partial F^\alpha}{\partial t^\beta}$, $\tilde{R}_1 = \begin{pmatrix} 0 & 0 \\ -c & 0 \end{pmatrix}$, and $\tilde{R}_i = 0$ for $i \geq 2$. Recall that the vector of conformal dimensions is $(\gamma_1, \gamma_2) = (2 - 2c, 2c)$. We have three subcases.

Case III. If $\gamma_1 \neq \gamma_2$ and $\gamma = \gamma_1$, we obtain

$$P_{2,0}^1 = \frac{c-1}{2} e^{-2v} + A(c-1)e^{2(c-1)v} \varepsilon^2 \left(u_{xx} - \frac{2c-3}{2c} e^{-v} v_x^2 + \frac{2c-3}{c} e^{-v} v_{xx} \right) + \mathcal{O}(\varepsilon^4),$$

$$P_{2,0}^2 = cu - (2c-1)e^{-v} + Ae^{2(c-1)v} \varepsilon^2 \left(-\frac{2c-1}{2} e^v u_{xx} + \frac{(c-1)^2}{c} v_x^2 - \frac{(c-1)^2}{c} v_{xx} \right) + \mathcal{O}(\varepsilon^4),$$

and

$$P_{1,1}^1 = \frac{c^2}{2} u^2 + \frac{c(c-1)}{4} (2v+1)e^{-2v} + Ae^{2(c-1)v} \varepsilon^2 \left(c((c-1)v+1)u_{xx} - \frac{c-1}{2} ((2c-3)v+3)e^{-v} v_x^2 + \frac{c-1}{2} ((2c-3)v+2)e^{-v} v_{xx} \right) + \mathcal{O}(\varepsilon^4),$$

$$P_{1,1}^2 = c^2 uv - c(2c-1)(v+1)e^{-v} + Ae^{2(c-1)v} \varepsilon^2 \left(-\frac{c}{2} ((2c-1)v+2)e^v u_{xx} + \frac{c-1}{2} ((2c-2)v+3)v_x^2 - (c-1)((c-1)v+1)v_{xx} \right) + \mathcal{O}(\varepsilon^4),$$

where A is an arbitrary complex constant. This deformation is given by the descendant DR hierarchy corresponding to the framing $(\mathcal{X}^1, \mathcal{X}^2) = \frac{12A}{c} (e^{2(c-1)v}, -e^{2(c-1)v})$.

Case II2. If $\gamma_1 \neq \gamma_2$ and $\gamma = \gamma_2$, we obtain

$$P_{2,0}^1 = \frac{c-1}{2} e^{-2v} + B(c-1)e^{-2cv} \varepsilon^2 \left(-u_{xx} + \frac{2c+1}{2c} e^{-v} v_x^2 - \frac{2c+1}{2c} e^{-v} v_{xx} \right) + \mathcal{O}(\varepsilon^4),$$

$$P_{2,0}^2 = cu - e^{-v}(2c-1) + Be^{-2cv} \varepsilon^2 \left(\frac{2c-1}{2} e^v u_{xx} - cv_x^2 + cv_{xx} \right) + \mathcal{O}(\varepsilon^4),$$

and

$$P_{1,1}^1 = \frac{c^2}{2} u^2 + \frac{c(c-1)}{4} (2v+1)e^{-2v} + B(c-1)e^{-2cv} \varepsilon^2 \left(-(cv-1)u_{xx} + \frac{1}{2} ((2c+1)v-3)e^{-v} v_x^2 - \frac{1}{2} ((2c+1)v-2)e^{-v} v_{xx} \right) + \mathcal{O}(\varepsilon^4),$$

$$P_{1,1}^2 = c^2 uv - c(2c-1)(v+1)e^{-v} + Bce^{-2cv} \varepsilon^2 \left(\frac{1}{2} ((2c-1)v-2)e^v u_{xx} - \frac{1}{2} (2cv-3)v_x^2 + (cv-1)v_{xx} \right) + \mathcal{O}(\varepsilon^4),$$

where B is an arbitrary complex constant. This deformation is given by the descendant DR hierarchy corresponding to the framing $(\mathcal{X}^1, \mathcal{X}^2) = \frac{12B}{c} \left(\frac{c-1}{c} e^{-2cv}, -e^{-(2c-1)v} \right)$.

Case II3. If $\gamma_1 = \gamma_2$ (which is equivalent to $c = \frac{1}{2}$) and γ coincides with them, we get a two-parameter family of deformations formed by linear combinations of the deformations from Cases II1 and II2.

Case III. There is a unique homogeneous calibration of standard type such that $\tilde{R}_i = 0$ for $i \geq 1$. Recall that the vector of conformal dimensions is $(\gamma_1, \gamma_2) = (-2c, 2c)$.

Case III1. If $\gamma_1 \neq \gamma_2$ (equivalently, $c \neq 0$) and $\gamma = \gamma_1 = -2c$, we obtain

$$\begin{aligned} P_{1,0}^1 &= u(1 - 2c \ln u) + cv + Au^{-2c-1} \varepsilon^2 \left(- \left(c + \frac{1}{2} \right) u^{-1} u_x^2 + cv_{xx} + \left(\frac{3}{2} - c(1 + \ln u) \right) u_{xx} \right) + \mathcal{O}(\varepsilon^4), \\ P_{1,0}^2 &= u(\ln u - c(1 + \ln^2 u)) + Au^{-2c-1} \varepsilon^2 \left(-u^{-1} \left(\left(c + \frac{1}{2} \right) \ln u + c \right) u_x^2 + \right. \\ &\quad \left. + \left(c(1 + \ln u) - \frac{1}{2} \right) v_{xx} + \left(2(1 + \ln u) - c(1 + \ln u)^2 - \frac{1}{2c} \right) u_{xx} \right) + \mathcal{O}(\varepsilon^4), \end{aligned}$$

and

$$\begin{aligned} P_{2,1}^1 &= -\frac{c}{2} u(cu(1 + 2 \ln u) - 2cv - u) + Au^{-2c-1} \varepsilon^2 \left(c(c-1)uv_{xx} + c(1-c)u_x^2 + \right. \\ &\quad \left. + u \left(c(1-c) \ln u - \left(c - \frac{1}{2} \right) (c-2) \right) u_{xx} \right) + \mathcal{O}(\varepsilon^4), \\ P_{2,1}^2 &= -\frac{c}{4} u^2(2(c \ln u + c - 1) + c - 1) \ln u + \frac{c^2}{2} v^2 + \\ &\quad + Au^{-2c-1} \varepsilon^2 \left(c \left((1-c) \ln u + \frac{3}{2} - c \right) u_x^2 + cu \left((c-1) \ln u + c - \frac{3}{2} \right) v_{xx} + \right. \\ &\quad \left. + u \left(c(1-c) \ln^2 u - (2c^2 - 4c + 1) \ln u - c^2 + 3c - \frac{3}{2} \right) u_{xx} \right) + \mathcal{O}(\varepsilon^4), \end{aligned}$$

where A is an arbitrary complex constant. This deformation is given by the descendant DR hierarchy corresponding to the framing $(\mathcal{X}^1, \mathcal{X}^2) = -\frac{12A}{c} (u^{-2c}, u^{-2c}(1 + \ln u))$.

Case III2. If $\gamma_1 \neq \gamma_2$ (equivalently, $c \neq 0$) and $\gamma = \gamma_2 = 2c$, we obtain

$$\begin{aligned} P_{1,0}^1 &= u(1 - 2c \ln u) + cv + Bu^{2c-1} \varepsilon^2 \left(\left(c - \frac{1}{2} \right) u^{-1} u_x^2 - cv_{xx} + \left(c \ln u + c + \frac{1}{2} \right) u_{xx} \right) + \mathcal{O}(\varepsilon^4), \\ P_{1,0}^2 &= u(\ln u - c(1 + \ln^2 u)) + Bu^{2c-1} \varepsilon^2 \left(\left(\frac{1}{2} - c(1 + \ln u) \right) v_{xx} + \right. \\ &\quad \left. + \left(c(\ln u + 1)^2 - \frac{1}{2c} \right) u_{xx} + u^{-1} \left(\left(c - \frac{1}{2} \right) \ln u + c - 1 + \frac{1}{2c} \right) u_x^2 \right) + \mathcal{O}(\varepsilon^4), \end{aligned}$$

and

$$\begin{aligned} P_{2,1}^1 &= -\frac{c}{2} u(cu(1 + 2 \ln u) - 2cv - u) + \frac{c}{2} Bu^{2c-1} \varepsilon^2 \left(2(c+1)u_x^2 - 2(c+1)uv_{xx} + \right. \\ &\quad \left. + u(2(c+1) \ln u + 2c + 3)u_{xx} \right) + \mathcal{O}(\varepsilon^4), \\ P_{2,1}^2 &= -\frac{c}{4} u^2(2(c \ln u + c - 1) + c - 1) \ln u + \frac{c^2}{2} v^2 + \\ &\quad + Bu^{2c-1} \varepsilon^2 \left(\left(c^2 + \frac{c}{2} - 1 + c(c+1) \ln u \right) u_x^2 + u \left(1 - \frac{c}{2} - c^2 - c(c+1) \ln u \right) v_{xx} + \right. \\ &\quad \left. + u \left(c^2 + c - \frac{3}{2} + c(c+1)(\ln u + 2) \ln u \right) u_{xx} \right) + \mathcal{O}(\varepsilon^4), \end{aligned}$$

Table 1. Functional parameters for the integrable deformations at the approximation up to ε^2

Case	Values of c	Integrable first order deformations	Integrable second order deformations
I	$c \neq \pm 2$	Miura trivial	Two functional parameters
I	$c = \pm 2$	One functional parameter	Two functional parameters
II	$c \neq 1$	Miura trivial	Two functional parameters
II	$c = 1$	One functional parameter	Two functional parameters
III	arbitrary c	Miura trivial	Two functional parameters

where B is an arbitrary complex constant. This deformation is given by the descendant DR hierarchy corresponding to the framing $(\mathcal{X}^1, \mathcal{X}^2) = \frac{12B}{c} (-u^{2c}, u^{2c} (\frac{1}{c} - 1 - \ln u))$.

Case III3. If $\gamma_1 = \gamma_2$ (equivalently, $c = 0$) and $\gamma = \gamma_i = 0$, we get the two-parameter family of deformations

$$P_{1,1}^1 = \frac{u^2}{2} + A\varepsilon^2 u_{xx} + \mathcal{O}(\varepsilon^4),$$

$$P_{1,1}^2 = \frac{v^2}{2} + B\varepsilon^2 v_{xx} + \mathcal{O}(\varepsilon^4),$$

which is given by the descendant DR hierarchy corresponding to the framing $(\mathcal{X}^1, \mathcal{X}^2) = 12(A, B)$. This hierarchy is just the DR hierarchy of the rank 2 F-topological field theory

$$c_{g,n+1}(e^{\alpha_0} \otimes \otimes_{i=1}^n e^{\alpha_i}) = \begin{cases} A^g, & \text{if } \alpha_0 = \dots = \alpha_n = 1, \\ B^g, & \text{if } \alpha_0 = \dots = \alpha_n = 2, \\ 0, & \text{otherwise,} \end{cases}$$

and it coincides with the system of two uncoupled KdV hierarchies.

3.3. General integrable deformations and open problems. In Sect. 3.2.1, we considered the problem of classification of dispersive deformations, containing only even powers of ε and satisfying properties (1)–(5), of principal hierarchies of two-dimensional homogeneous semisimple flat F-manifolds. We observed that at the approximation up to ε^2 all such deformations are given by the descendant DR hierarchies.

In this section, we consider more general dispersive deformations of the same rank 2 principal hierarchies: first, we allow odd powers of ε in the dispersive deformation (3.7), and, second, we require that only properties (1)–(2) are satisfied. In other words, we require only integrability, i.e., pairwise commutativity of the flows. In the table below, we summarize the results of computations of such deformations at the approximation up to ε^2 (the results for Case I were already obtained in [AL18]). When we refer to a functional parameter relative to an integrable deformation, we mean that at a specified order the equivalence classes of deformations depend on an arbitrary function. Recall (see Definition 6) that two deformations are said to be equivalent if they are related by a Miura transformation that is close to identity.

For special values of the functional parameters, we recover the genus one approximations of the descendant DR hierarchies from Sect. 3.2.3. Unfortunately, for generic choices of the functional parameters the existence of a full dispersive hierarchy is an open problem. Concerning this we point out that, in [AL18], it was conjectured that, up to equivalence, integrable deformations for systems of any rank are labelled by a simple

set of invariants called *Miura invariants*. Consider a system of evolutionary PDEs of the form

$$\begin{aligned} \frac{\partial u^\alpha}{\partial t} = & A_\beta^\alpha(u^*)u_x^\beta + \varepsilon \left(B_\beta^\alpha(u^*)u_{xx}^\beta + B_{\beta\gamma}^\alpha(u^*)u_x^\beta u_x^\gamma \right) \\ & + \varepsilon^2 \left(C_\beta^\alpha(u^*)u_{xxx}^\beta + C_{\beta\gamma}^\alpha(u^*)u_x^\beta u_{xx}^\gamma + C_{\beta\gamma\delta}^\alpha(u^*)u_x^\beta u_x^\gamma u_x^\delta \right) + \dots, \\ & \alpha = 1, \dots, N, \end{aligned} \tag{3.10}$$

and as in the proof of Theorem 4 consider the associated Miura matrix

$$M_\beta^\alpha(u^*, p) = A_\beta^\alpha(u^*) + B_\beta^\alpha(u^*)p + C_\beta^\alpha(u^*)p^2 + \dots$$

The *Miura invariants* of the system (3.10) are the eigenvalues $\lambda^i(u^*, p)$ of the Miura matrix. If the eigenvalues of the matrix (A_β^α) are pairwise distinct at some point $(u^1, \dots, u^N) = (u_{\text{orig}}^1, \dots, u_{\text{orig}}^N) \in \mathbb{C}^N$, then the Miura invariants are well defined as formal power series whose coefficients are functions on an open neighbourhood of $(u_{\text{orig}}^1, \dots, u_{\text{orig}}^N)$:

$$\lambda^i = v^i + \lambda_1^i p + \lambda_2^i p^2 + \dots, \quad i = 1, \dots, N. \tag{3.11}$$

The functional parameters of Table 1 can be identified with a part of the coefficients $\lambda_1^1, \lambda_1^2, \lambda_2^1, \lambda_2^2$ in formula (3.11). The presence of odd powers of p in the expansion (3.11) seems an exceptional phenomenon. In the case of special deformations satisfying all the properties (1)–(5) from Sect. 3.2.1, there are examples related to open Gromov–Witten theory [BR18] and we expect that this is not a coincidence. However, this point requires further investigation.

Acknowledgements The work of A. B. is supported by the Russian Science Foundation (Grant No. 20-71-10110). P. L. is supported by MIUR–FFABR funds 2017 and by funds of H2020-MSCA-RISE-2017 Project No. 778010 IPaDEGAN.

Funding Open access funding provided by Università degli Studi di Padova within the CRUI-CARE Agreement.

Open Access This article is licensed under a Creative Commons Attribution 4.0 International License, which permits use, sharing, adaptation, distribution and reproduction in any medium or format, as long as you give appropriate credit to the original author(s) and the source, provide a link to the Creative Commons licence, and indicate if changes were made. The images or other third party material in this article are included in the article’s Creative Commons licence, unless indicated otherwise in a credit line to the material. If material is not included in the article’s Creative Commons licence and your intended use is not permitted by statutory regulation or exceeds the permitted use, you will need to obtain permission directly from the copyright holder. To view a copy of this licence, visit <http://creativecommons.org/licenses/by/4.0/>.

Publisher’s Note Springer Nature remains neutral with regard to jurisdictional claims in published maps and institutional affiliations.

References

[ABLR20] Arsie, A., Buryak, A., Lorenzoni, P., Rossi, P.: Semisimple flat F-manifolds in higher genus. [arXiv:2001.05599](https://arxiv.org/abs/2001.05599)

[AL13a] Arsie, A., Lorenzoni, P.: From the Darboux–Egorov system to bi-flat F-manifolds. *J. Geom. Phys.* **70**, 98–116 (2013)

[ALM15] Arsie, A., Lorenzoni, P., Moro, A.: On integrable conservation laws. *Proceedings A* **471**(2173), 20140124 (2015)

- [AL17] Arsie, A., Lorenzoni, P.: Complex reflection groups, logarithmic connections and bi-flat F-manifolds. *Lett. Math. Phys.* **107**(10), 1919–1961 (2017)
- [AL18] Arsie, A., Lorenzoni, P.: Flat F-manifolds, Miura invariants and integrable systems of conservation laws. *J. Integrable Sys.* **3**(1), xyy004 (2018)
- [AL19] Arsie, A., Lorenzoni, P.: F-manifolds, multi-flat structures and Painlevé transcendents. *Asian J. Math.* **23**(5), 877–904 (2019)
- [BB19] Basalaev, A., Buryak, A.: Open WDVV equations and Virasoro constraints. *Arnold Math. J.* **5**(2–3), 145–186 (2019)
- [Bur15] Buryak, A.: Double ramification cycles and integrable hierarchies. *Commun. Math. Phys.* **336**, 1085–1107 (2015)
- [BCT18] Buryak, A., Clader, E., Tessler, R.J.: Open r -spin theory II: The analogue of Witten’s conjecture for r -spin disks. [arXiv:1809.02536v4](https://arxiv.org/abs/1809.02536v4)
- [BCT19] Buryak, A., Clader, E., Tessler, R.J.: Closed extended r -spin theory and the Gelfand-Dickey wave function. *J. Geom. Phys.* **137**, 132–153 (2019)
- [BDGR18] Buryak, A., Dubrovin, B., Guéré, J., Rossi, P.: Tau-structure for the double ramification hierarchies. *Commun. Math. Phys.* **363**(1), 191–260 (2018)
- [BDGR20] Buryak, A., Dubrovin, B., Guéré, J., Rossi, P.: Integrable systems of double ramification type. *Int. Math. Res. Not.* **2020**(24), 10381–10446 (2020)
- [BGR19] Buryak, A., Dubrovin, B., Guéré, J.: DR/DZ equivalence conjecture and tautological relations. *Geom. Topol.* **23**(7), 3537–3600 (2019)
- [BR16a] Buryak, A., Rossi, P.: Recursion relations for double ramification hierarchies. *Commun. Math. Phys.* **342**, 533–568 (2016)
- [BR16b] Buryak, A., Rossi, P.: Double ramification cycles and quantum integrable systems. *Lett. Math. Phys.* **106**, 289–317 (2016)
- [BR18] Buryak, A., Rossi, P.: Extended r -spin theory in all genera and the discrete KdV hierarchy. [arXiv:1806.09825](https://arxiv.org/abs/1806.09825)
- [BRS20] Buryak, A., Rossi, P., Shadrin, S.: Towards a bihamiltonian structure for the double ramification hierarchy. *Lett. Math. Phys.* **111**, 13 (2021)
- [BSSZ15] Buryak, A., Shadrin, S., Spitz, L., Zvonkine, D.: Integrals of ψ -classes over double ramification cycles. *Am. J. Math.* **137**(3), 699–737 (2015)
- [CPS18] Carlet, G., Posthuma, H., Shadrin, S.: Deformations of semisimple Poisson pencils of hydrodynamic type are unobstructed. *J. Differ. Geom.* **108**(1), 63–89 (2018)
- [DH17] David, L., Hertling, C.: Regular F-manifolds: initial conditions and Frobenius metrics. *Annali della Scuola Normale di Pisa, Classe di Scienze* **17**(3), 1121–1152 (2017)
- [Dub96] Dubrovin, B.: Geometry of 2D topological field theories. Integrable systems and quantum groups (Montecatini Terme, 1993), 120–348, *Lecture Notes in Math.*, 1620, Fond. CIME/CIME Found. Subser., Springer, Berlin (1996)
- [DLZ06] Dubrovin, B., Liu, S.-Q., Zhang, Y.: On Hamiltonian perturbations of hyperbolic systems of conservation laws I: quasi-triviality of bi-Hamiltonian perturbations. *Commun. Pure Appl. Math.* **59**(4), 559–615 (2006)
- [DZ01] Dubrovin, B., Zhang, Y.: Normal forms of hierarchies of integrable PDEs, Frobenius manifolds and Gromov–Witten invariants. [arXiv:math/0108160](https://arxiv.org/abs/math/0108160)
- [FP00] Faber, C., Pandharipande, R.: Logarithmic series and Hodge integrals in the tautological ring. With an appendix by Don Zagier. *Mich. Math. J.* **48**(1), 215–252 (2000)
- [Get04] Getzler, E.: The jet-space of a Frobenius manifold and higher-genus Gromov–Witten invariants. Frobenius manifolds, 45–89, *Aspects Math.*, E36, Friedr. Vieweg, Wiesbaden (2004)
- [Giv01] Givental, A.: Semisimple Frobenius structures at higher genus. *Int. Math. Res. Not.* **2001**(23), 1265–1286 (2001)
- [Hai13] Hain, R.: Normal functions and the geometry of moduli spaces of curves. *Handbook of moduli*. Vol. I, 527–578, *Adv. Lect. Math. (ALM)*, 24, Int. Press, Somerville, MA (2013)
- [JPPZ17] Janda, F., Pandharipande, R., Pixton, A., Zvonkine, D.: Double ramification cycles on the moduli spaces of curves. *Publications Mathématiques Institut de Hautes Études Scientifiques.* **125**, 221–266 (2017)
- [KMS15] Kato, M., Mano, T., Sekiguchi, J.: Flat structure on the space of isomonodromic deformations. [arXiv:1511.01608](https://arxiv.org/abs/1511.01608)
- [KMS18] Konishi, Y., Minabe, S., Shiraishi, Y.: Almost duality for Saito structure and complex reflection groups. *J. Integrable Syst.* **3**(1), xyy003 (2018)
- [Kon92] Kontsevich, M.: Intersection theory on the moduli space of curves and the matrix Airy function. *Commun. Math. Phys.* **147**, 1–23 (1992)
- [LPR09] Lorenzoni, P., Pedroni, M., Raimondo, A.: F-manifolds and integrable systems of hydrodynamic type. *Arch. Math.* **47**(3), 163–180 (2011)

- [Lor14] Lorenzoni, P.: Darboux-Egorov system, bi-flat F-manifolds and Painlevé VI. *Int. Math. Res. Not.* **2014**(12), 3279–3302 (2014)
- [Man05] Manin, Y.: F-manifolds with flat structure and Dubrovin’s duality. *Adv. Math.* **198**(1), 5–26 (2005)
- [MW13] Marcus, S., Wise, J.: Stable maps to rational curves and the relative Jacobian. [arXiv:1310.5981](https://arxiv.org/abs/1310.5981)
- [PST14] Pandharipande, R., Solomon, J.P., Tessler, R.J.: Intersection theory on moduli of disks, open KdV and Virasoro. [arXiv:1409.2191v2](https://arxiv.org/abs/1409.2191v2)
- [Ros17] Rossi, P.: Integrability, quantization and moduli spaces of curves. *SIGMA* **13**, 060 (2017)
- [Sab98] Sabbah, C.: Frobenius manifolds: isomonodromic deformations and infinitesimal period mappings. *Expos. Math.* **16**(1), 1–57 (1998)
- [SZ11] Shadrin, S., Zvonkine, D.: A group action on Losev–Manin cohomological field theories. *Ann. Inst. Fourier* **61**(7), 2719–2743 (2011)
- [Tel12] Teleman, C.: The structure of 2D semi-simple field theories. *Invent. Math.* **188**(3), 525–588 (2012)
- [Wit91] Witten, E.: Two-dimensional gravity and intersection theory on moduli space. *Surv. Differ. Geom.* **1**, 243–310 (1991)

Communicated by H.-T. Yau