## Research Article

Elio Marconi*

# Rectifiability of entropy defect measures in a micromagnetics model 

https://doi.org/10.1515/acv-2021-0012
Received February 2, 2021; revised April 21, 2021; accepted April 29, 2021


#### Abstract

We study the fine properties of a class of weak solutions $u$ of the eikonal equation arising as asymptotic domain of a family of energy functionals introduced in [T. Rivière and S. Serfaty, Limiting domain wall energy for a problem related to micromagnetics, Comm. Pure Appl. Math. 54 (2001), no. 3, 294-338]. In particular, we prove that the entropy defect measure associated to $u$ is concentrated on a 1-rectifiable set, which detects the jump-type discontinuities of $u$.


Keywords: Eikonal equation, rectifiability, Lagrangian representation, entropy production, kinetic formulation, micromagnetics

MSC 2010: 35F20, 35L67, 35L65

Communicated by: Tristan Riviére

## 1 Introduction

We consider a bounded simply connected domain $\Omega \subset \mathbb{R}^{2}$ and we investigate the fine properties of the following class of divergence free unit vector fields:

Definition 1.1. We denote by $\mathcal{M}_{\text {div }}(\Omega)$ the set of vector fields $u: \Omega \rightarrow \mathbb{C}$ for which the following conditions hold:
(i) $\operatorname{div} u=0$ in the sense of distributions.
(ii) There exists $\phi \in L^{\infty}(\Omega)$ such that $u=e^{i \phi}$ and

$$
\left\langle U_{\phi}, \psi(x, a)\right\rangle:=\int_{\Omega \times \mathbb{R}} e^{i(\phi(x) \wedge a)} \cdot \nabla_{x} \psi(x, a) d x d a \in \mathcal{M}(\Omega \times \mathbb{R}),
$$

where $\mathcal{M}(\Omega \times \mathbb{R})$ denotes the set of finite Radon measures on $\Omega \times \mathbb{R}$.
The space $\mathcal{M}_{\mathrm{div}}(\Omega)$ is the conjectured asymptotic domain as $\varepsilon \rightarrow 0$ of the following family of energy functionals introduced in [20] in the context of micromagnetics:

$$
E_{\varepsilon}(u):=\varepsilon \int_{\Omega}|\nabla u|^{2}+\frac{1}{\varepsilon} \int_{\mathbb{R}^{2}}\left|H_{u}\right|^{2},
$$

where $u \in W^{1,2}\left(\Omega, \mathbb{S}^{1}\right)$ and the so-called demagnetizing field $H_{u} \in L^{2}\left(\mathbb{R}^{2} ; \mathbb{R}^{2}\right)$ is such that curl $H_{u}=0$ and $\operatorname{div}\left(\tilde{u}+H_{u}\right)=0$ in $\mathcal{D}^{\prime}\left(\mathbb{R}^{2}\right)$, where

$$
\tilde{u}(x)= \begin{cases}u(x) & \text { if } x \in \Omega \\ 0 & \text { otherwise }\end{cases}
$$

[^0]The following compactness result was proven in [21]: let $\phi_{\varepsilon_{n}}$ be a bounded sequence in $L^{\infty}(\Omega)$ such that $E_{\varepsilon_{n}}\left(u_{\varepsilon_{n}}\right)$ is uniformly bounded, where $u_{\varepsilon_{n}}=e^{i \phi_{\varepsilon_{n}}}$ and $\varepsilon_{n} \rightarrow 0$; then $\phi_{\varepsilon_{n}}$ is relatively compact in $L^{p}(\Omega)$ for every $p \in[1, \infty)$ and for every limit point $\bar{\phi}$ it holds

$$
\begin{equation*}
e^{i \bar{\phi}} \in \mathcal{M}_{\operatorname{div}}(\Omega) \quad \text { and } \quad\left|U_{\bar{\phi}}\right|(\Omega \times \mathbb{R}) \leq \liminf _{n \rightarrow \infty} E_{\varepsilon_{n}}\left(u_{\varepsilon_{n}}\right) \tag{1.1}
\end{equation*}
$$

Although the $\Gamma$-lim inf inequality (1.1) was proved in full generality, the corresponding $\Gamma$-lim sup inequality was obtained only in special cases. In particular, the energy-minimizing configurations were characterized by the results in $[5,21]$. It is expected that the energy $E_{\varepsilon}$ is concentrated on lines at a scale $\varepsilon>0$ around the lines, allowing for sharper and sharper jumps as $\varepsilon \rightarrow 0$; the latter correspond in three dimensions to jumps across surfaces, called domain walls in the theory of micromagnetism (see [20]). These lines are detected by the measure $U_{\phi}$ : in particular, if we denote by $p_{x}: \Omega \times \mathbb{R} \rightarrow \Omega$ the standard projection on the first component and if $\phi \in \operatorname{BV}(\Omega)$, then the measure

$$
v:=\left(p_{x}\right)_{\sharp}\left|U_{\phi}\right|
$$

is concentrated on the 1 -rectifiable jump set of $\phi$.
However, vector fields in $\mathcal{M}_{\text {div }}(\Omega)$ do not have necessarily bounded variation and a study of their fine properties must therefore be independent of the theory of BV functions. This program was announced in [5] and carried on in [4] leading to the following result.

Theorem 1.2. Let $\phi$ be a lifting of $u \in \mathcal{M}_{\text {div }}(\Omega)$ as in Definition 1.1. Then the following assertions hold:
(1) The jump set $J$ of $\phi$ is countably $\mathscr{H}^{1}$-rectifiable and coincides, up to $\mathscr{H}^{1}$-negligible sets, with

$$
\begin{equation*}
\Sigma:=\left\{x \in \Omega: \limsup _{r \rightarrow 0} \frac{v\left(B_{r}(x)\right)}{r}>0\right\} . \tag{1.2}
\end{equation*}
$$

Moreover, for every $a \in \mathbb{R}$ it holds

$$
\operatorname{div}\left(e^{i \phi \wedge a}\right)_{\llcorner } J=\mathbf{1}_{\phi^{-}<a<\phi^{+}}\left(e^{i a}-e^{i \phi^{-}}\right) \cdot \mathbf{n}_{J} \mathscr{H}_{\llcorner }^{1} J
$$

where $\mathbf{n}_{J}$ denotes the normal to $J$.
(2) Every $x \in \Omega \backslash \Sigma$ is a vanishing mean oscillation point of $\phi$, namely

$$
\lim _{r \rightarrow 0} \frac{1}{r^{2}} \int_{B_{r}(x)}\left|\phi-\phi_{r}(x)\right|=0
$$

where $\phi_{r}(x)$ is the average of $\phi$ on $B_{r}(x)$.
(3) The measure $v_{\mathrm{L}}(\Omega \backslash J)$ is orthogonal to $\mathscr{H}^{1}$, namely

$$
B \subset(\Omega \backslash J) \text { Borel with } \mathscr{H}^{1}(B)<\infty \quad \text { implies } \quad v(B)=0
$$

We observe that for functions $\phi \in \operatorname{BV}_{\mathrm{loc}}(\Omega)$ the above properties (2) and (3) can be improved to the following: (2') $\mathscr{H}^{1}$-a.e. point in $\Omega \backslash J$ is a Lebesgue point of $\phi$.
(3') The measure $v_{\mathrm{L}}(\Omega \backslash J)$ is identically 0 .
In [4], it was conjectured that both (2') and (3') hold for every $u \in \mathcal{M}_{\operatorname{div}}(\Omega)$. The following weaker version of ( $2^{\prime}$ ) was recently obtained in [14] in the close setting of weak solutions $u$ with finite entropy production of the Burgers equation:
(2^) The set of non-Lebesgue points of $u$ has Hausdorff dimension at most 1.
This result was extended for general conservation laws in [18], implying in particular that property (2*) holds in the setting of this paper, namely for functions $\phi \in L^{\infty}$ corresponding to vector fields $u \in \mathcal{M}_{\text {div }}(\Omega)$.

The main result of this paper is the proof of property ( $3^{\prime}$ ) for general vector fields $u \in \mathcal{M}_{\text {div }}(\Omega)$.
Theorem 1.3. Let $\phi$ be a lifting of $u \in \mathcal{M}_{\operatorname{div}}(\Omega)$ as in Definition 1.1. Then the measure $v$ is concentrated on the countably $\mathscr{H}^{1}$-rectifiable set $\Sigma$ defined in (1.2). In particular, for every $a \in \mathbb{R}$ it holds

$$
\operatorname{div}\left(e^{i \phi \wedge a}\right)=\mathbf{1}_{\phi^{-}<a<\phi^{+}}\left(e^{i a}-e^{i \phi^{-}}\right) \cdot \mathbf{n}_{J} \mathscr{H}^{1}\llcorner J
$$

Theorem 1.3 establishes that the concentration property expected for the $\Gamma$-limit functional of $E_{\varepsilon}$ as $\varepsilon \rightarrow 0$ holds for the candidate $\Gamma$-limit; this property is also considered as a fundamental step to complete the $\Gamma$-lim sup analysis (see [16]).

### 1.1 Main tool and strategy of the proof

The strategy of the proof of Theorem 1.3 was introduced in [19] to prove the analogous result for weak solutions with finite entropy production of the Burgers equation (or, more generally, 1D scalar conservation laws with uniformly convex flux). Indeed, there is a strong analogy between weak solutions to conservation laws with finite entropy production and the solutions to the eikonal equation arising in this model or the related model introduced by Aviles and Giga in [6]. In particular, Theorem 1.2 has an analogous version for scalar conservation laws (see [11, 15]) and for the model by Aviles and Giga [10]. In order to compare the setting of this paper and the one of conservation laws, we observe that for $u=e^{i \phi} \in \mathcal{M}_{\text {div }}(\Omega)$ it holds

$$
\partial_{x_{1}} \cos \phi+\partial_{x_{2}} \sin \phi=0
$$

Let us assume that $\phi$ takes values in $(0, \pi)$ so that the cosine is invertible in the range of $\phi$ and $v=\cos \phi$ satisfies the equation

$$
\partial_{x_{1}} v+\partial_{x_{2}}\left(\sin \left(\cos ^{-1}(v)\right)\right)=0
$$

Since the map $\sin \circ \cos ^{-1}$ is convex on $(-1,1)$, it is possible to transfer the results obtained for conservation laws with convex fluxes to solutions of the eikonal equation taking values in $(0, \pi)$. When instead the oscillation of $\phi$ is larger than $\pi$, the approach above fails and more refined arguments are needed.

The main tool used to prove Theorem 1.3 is the so called Lagrangian representation, which was introduced in [7] for entropy solutions to general conservation laws and then extended in [18] to weak solutions with finite entropy production. This Lagrangian representation (see Definition 3.1) is an extension of the classical method of characteristics to this non-smooth setting and it is strongly inspired by Ambrosio's superposition principle in the context of positive measure valued solutions to the linear continuity equation (see, for example, [1]). Roughly speaking, the evolution of the solution is obtained as superposition of single trajectories traveling with characteristic speed. This tool is well suited for our purposes since also the kinetic measure $U_{\phi}$ can be decomposed along the characteristic trajectories detected by the Lagrangian representation. In Section 3, we prove the existence of a Lagrangian representation for vector fields in $\mathcal{M}_{\text {div }}(\Omega)$ building on the following kinetic formulation obtained in [21] (see also [13] in the study of the model by Aviles and Giga, and the fundamental paper [17] in the setting of entropy solutions to scalar conservation laws): by setting $\chi(x, a):=\mathbf{1}_{\phi(x) \geq a}$, it holds

$$
\begin{equation*}
i e^{i a} \cdot \nabla_{\chi} \chi=-\partial_{a} U_{\phi} \quad \text { in } \mathcal{D}^{\prime}(\Omega \times \mathbb{R}) \tag{1.3}
\end{equation*}
$$

The proof of the existence of a Lagrangian representation follows the strategy of [18], but additional work is required since we consider here solutions on bounded domains instead of the whole $\mathbb{R}^{2}$.

Once a Lagrangian representation is available for vector fields in $\mathcal{M}_{\operatorname{div}}(\Omega)$, we implement the strategy introduced in [19] to prove Theorem 1.3. Since the oscillation of $\phi$ is bigger than $\pi$, the argument does not apply straightforwardly. Still a partial result is obtained in Section 4.2 by covering the image of $\phi$ with finitely many intervals $\left(I_{l}\right)_{l=1}^{L}$ of length less than $\pi$ and appropriately localizing the argument of [19]. A new regularity estimate is proven in Section 4.3 and this allows to conclude the proof of Theorem 1.3, relying on Theorem 1.2.

## 2 Preliminaries

### 2.1 Duality for $L^{1}$-optimal transport

In this section, we recall a few facts about $L^{1}$-optimal transport. We state the results in the form that we will need in Section 3. Given a metric space $X$, we denote by $\mathcal{M}_{+}(X)$ the set of finite non-negative Borel measures on $X$.

Definition 2.1. Let $(X, d)$ be a complete and separable metric space and let $\mu_{1}, \mu_{2} \in \mathcal{M}_{+}(X)$ be such that $\mu_{1}(X)=\mu_{2}(X)$. The Wasserstein distance of order 1 between $\mu_{1}$ and $\mu_{2}$ is defined by

$$
\begin{equation*}
W_{1}\left(\mu_{1}, \mu_{2}\right):=\inf _{\pi \in \Pi\left(\mu_{1}, \mu_{2}\right)} \int_{X} d(x, y) d \pi(x, y) \tag{2.1}
\end{equation*}
$$

where $\Pi\left(\mu_{1}, \mu_{2}\right)$ is the set of transport plans from $\mu_{1}$ to $\mu_{2}$, i.e.

$$
\Pi\left(\mu_{1}, \mu_{2}\right):=\left\{\omega \in \mathcal{M}_{+}\left(X^{2}\right): \pi_{1 \sharp} \omega=\mu_{1}, \pi_{2 \sharp} \omega=\mu_{2}\right\},
$$

denoting by $\pi_{1}, \pi_{2}: X^{2} \rightarrow X$ the two natural projections.
Notice that $W_{1}$ can take the value $+\infty$.
In order to prove the existence of a Lagrangian representation for vector fields in $\mathcal{M}_{\text {div }}(\Omega)$, we will take advantage of the dual formulation of the $L^{1}$-optimal transport. The following duality formula can be found, for example, in [22].

Proposition 2.2. For any $\mu_{1}, \mu_{2} \in \mathcal{M}_{+}(X)$ with $\mu_{1}(X)=\mu_{2}(X)$, it holds

$$
W_{1}\left(\mu_{1}, \mu_{2}\right)=\sup _{\psi \in L^{1}\left(\mu_{1}\right),\|\psi\|_{\text {Lip }} \leq 1}\left(\int_{X} \psi d \mu_{1}-\int_{X} \psi d \mu_{2}\right)
$$

Since it will be convenient to allow that the two measures $\mu_{1}$, $\mu_{2}$ have different masses, we deduce from Proposition 2.2 the following result.

Corollary 2.3. Let $(X, d)$ be bounded and let $\mu_{1}, \mu_{2} \in \mathcal{N}_{+}(X)$. Assume that there exist $C_{1}, C_{2}>0$ such that for every $\psi \in \operatorname{Lip}(X)$ it holds

$$
\begin{equation*}
\left|\int_{X} \psi d \mu_{1}-\int_{X} \psi d \mu_{2}\right| \leq C_{1}|\psi|_{\text {Lip }}+C_{2}\|\psi\|_{L^{\infty}} \tag{2.2}
\end{equation*}
$$

Then there exist $\tilde{\mu}_{1} \leq \mu_{1}, \tilde{\mu}_{2} \leq \mu_{2}$ such that $\left\|\mu_{1}-\tilde{\mu}_{1}\right\| \leq C_{2},\left\|\mu_{2}-\tilde{\mu}_{2}\right\| \leq C_{2}$ and

$$
\begin{equation*}
W_{1}\left(\tilde{\mu}_{1}, \tilde{\mu}_{2}\right) \leq C_{1}+C_{2} \operatorname{diam}(X) \tag{2.3}
\end{equation*}
$$

Proof. We assume without loss of generality that $\alpha:=\left\|\mu_{1}\right\|-\left\|\mu_{2}\right\| \geq 0$. Let $\bar{\mu}_{2}=\mu_{2}+\alpha \delta_{\bar{x}}$ for some $\bar{x} \in X$. Then we have

$$
\begin{aligned}
\left|\int_{X} \psi d \mu_{1}-\int_{X} \psi d \bar{\mu}_{2}\right| & =\left|\int_{X}(\psi-\psi(\bar{x})) d \mu_{1}-\int_{X}(\psi-\psi(\bar{x})) d \bar{\mu}_{2}\right| \\
& =\left|\int_{X}(\psi-\psi(\bar{x})) d \mu_{1}-\int_{X}(\psi-\psi(\bar{x})) d \bar{\mu}_{2}\right| \\
& \leq C_{1}|\psi|_{\text {Lip }}+C_{2}|\psi|_{\text {Lip }} \operatorname{diam}(X) .
\end{aligned}
$$

By Proposition 2.2, it follows that $W_{1}\left(\mu_{1}, \bar{\mu}_{2}\right) \leq C_{1}+C_{2} \operatorname{diam}(X)$. Let $\pi \in \mathcal{N}\left(X^{2}\right)$ be an optimal plan with marginals $\mu_{1}$ and $\bar{\mu}_{2}$ and let $\tilde{\pi} \leq \pi$ be such that $\left(p_{2}\right)_{\sharp} \tilde{\pi}=\mu_{2}$. Then we check that the statement is true for $\tilde{\mu}_{1}=\left(p_{1}\right)_{\sharp} \tilde{\pi}$ and $\tilde{\mu}_{2}=\mu_{2}$ : indeed, $\tilde{\pi}$ is an admissible plan between $\tilde{\mu}_{1}$ and $\tilde{\mu}_{2}$ by construction. Since $\tilde{\pi} \leq \pi$ and $W_{1}\left(\mu_{1}, \bar{\mu}_{2}\right) \leq C_{1}+C_{2} \operatorname{diam}(X)$, inequality (2.3) holds true. Moreover,

$$
\left\|\mu_{2}-\tilde{\mu}_{2}\right\|=0 \quad \text { and } \quad\left\|\mu_{1}-\tilde{\mu}_{1}\right\|=\mu_{1}(X)-\tilde{\mu}_{1}(X)=\alpha
$$

since $\tilde{\mu}_{1} \leq \mu_{1}$. Finally, we observe that, by choosing $\psi \equiv 1$ in (2.2), we obtain $\alpha \leq C_{2}$. This shows that $\left\|\mu_{1}-\tilde{\mu}_{1}\right\| \leq C_{2}$ and completes the proof.

The next theorem from [8] provides the existence of an $L^{1}$-optimal map with respect to quite general distances on $\mathbb{R}^{N}$.

Theorem 2.4. Let $X=\mathbb{R}^{N}$ with $N \in \mathbb{N}$ be the Euclidean space equipped with the distance induced by a convex norm $|\cdot|_{D *}$. Let $\mu_{1}, \mu_{2} \in \mathcal{P}\left(\mathbb{R}^{N}\right)$ be two probability measures such that $\mu_{1} \ll \mathscr{L}^{N}$ and the infimum in (2.1) is finite. Then there exists an optimal plan $\pi$ in (2.1) induced by a map, i.e. there exists a measurable map $T: \mathbb{R}^{N} \rightarrow \mathbb{R}^{N}$ such that $T_{\sharp} \mu_{1}=\mu_{2}$ and

$$
W_{1}\left(\mu_{1}, \mu_{2}\right)=\int_{X}|T(x)-x|_{D *} d \mu_{1}(x)
$$

### 2.2 Weak convergence of measures

We will say that a sequence of measures $\left(\mu_{n}\right)_{n \in \mathbb{N}} \subset \mathcal{M}_{+}(X)$ is narrowly convergent to $\mu \in \mathcal{M}_{+}(X)$ if

$$
\lim _{n \rightarrow \infty} \int_{X} f d \mu_{n}=\int_{X} f d \mu \quad \text { for all } f \in C_{b}(X)
$$

where $C_{b}(X)$ denotes the set of continuous real-valued bounded functions on $X$. Moreover, we say that a bounded family $\mathscr{F} \subset \mathcal{M}_{+}(X)$ is tight if for every $\varepsilon>0$ there exists a compact set $K \subset X$ such that for every $\mu \in \mathscr{F}$ it holds

$$
\mu(X \backslash K)<\varepsilon
$$

The following classical theorem characterizes the relatively compact families in $\mathcal{M}_{+}(X)$ (see [9]).
Theorem 2.5 (Prokhorov). Let $X$ be a metric space. If a bounded family $\mathscr{F} \subset \mathcal{M}_{+}(X)$ is tight, then it is relatively compact with respect to the narrow convergence. If moreover $X$ is complete and separable, then also the converse implication holds.

## 3 Lagrangian representation for vector fields in $\mathcal{M}_{\text {div }}$

In this section, we introduce the notions of Lagrangian representations of the hypograph and of the epigraph for the liftings $\phi$ of vector fields in $\mathcal{M}_{\operatorname{div}}(\Omega)$. Moreover, we provide a suitable decomposition along characteristics of the kinetic measure $U_{\phi}$ introduced in (1.3).

### 3.1 Notation and main definition

We will consider the standard decomposition of the measure $D f \in \mathcal{M}(\mathbb{R})$, where $f \in \operatorname{BV}(\mathbb{R}, \mathbb{R})$ (see, for example, [2]). We will adopt the following notation:

$$
D f=D^{\mathrm{ac}} f+D^{c} f+D^{j} f=\tilde{D} f+D^{j} f
$$

where $D^{\text {ac }} f, D^{c} f$ and $D^{j} f$ denote the absolutely continuous part, the Cantor part and the atomic part of $D f$, respectively; we refer to $\tilde{D} f$ as the diffuse part of $D f$.

For every function $\phi: \Omega \rightarrow[0, M]$, we denote its hypograph and its epigraph by

$$
H_{\phi}:=\{(x, a) \in \Omega \times[0, M]: a \leq \phi(x)\} \quad \text { and } \quad E_{\phi}:=\{(x, a) \in \Omega \times[0, M]: a \geq \phi(x)\}
$$

respectively.
We denote by $B_{R}$ an open ball of radius $R$ such that $\overline{B_{R}} \subset \Omega$ and we set

$$
\Gamma:=\left\{\left(\gamma, t_{\gamma}^{-}, t_{y}^{+}\right): 0 \leq t_{\gamma}^{-} \leq t_{y}^{+} \leq 1, \gamma \in \mathrm{BV}\left(\left(t_{\gamma}^{-}, t_{y}^{+}\right) ; B_{R} \times[0, M]\right), \gamma_{x} \text { is Lipschitz }\right\}
$$

In order to fix a representative, we will always assume that $y$ is continuous from the right. For every $t \in(0,1)$, we consider the section

$$
\Gamma(t):=\left\{\left(\gamma, t_{\gamma}^{-}, t_{\gamma}^{+}\right) \in \Gamma: t \in\left(t_{\gamma}^{-}, t_{\gamma}^{+}\right)\right\}
$$

and we set

$$
\begin{aligned}
e_{t}: \Gamma(t) & \rightarrow B_{R} \times[0, M], \\
\left(\gamma, t_{\gamma}^{-}, t_{\gamma}^{+}\right) & \mapsto \gamma(t) .
\end{aligned}
$$

Sometimes we will identify the triple $\left(\gamma, t_{\gamma}^{-}, t_{\gamma}^{+}\right) \in \Gamma$ with the curve $\gamma$ itself to make the notation less heavy.
Definition 3.1. Suppose $u \in \mathcal{M}_{\operatorname{div}}(\Omega)$ and $\phi \in L^{\infty}(\Omega)$ as in Definition 1.1. We say that the Radon measure $\omega_{h} \in \mathcal{M}(\Gamma)$ is a Lagrangian representation of the hypograph of $\phi$ on $B_{R}$ if the following conditions hold:
(1) For every $t \in(0,1)$, it holds

$$
\left(e_{t}\right)_{\#}\left[\omega_{h\llcorner } \Gamma(t)\right]=\mathscr{L}^{3}{ }_{\llcorner } H_{\phi} .
$$

(2) The measure $\omega_{h}$ is concentrated on the set of curves $\gamma \in \Gamma$ such that for $\mathscr{L}^{1}$-a.e. $t \in\left(t_{\gamma}^{-}, t_{\gamma}^{+}\right)$the following characteristic equation holds:

$$
\begin{equation*}
\dot{\gamma}_{x}(t)=i e^{i y_{a}(t)} \tag{3.1}
\end{equation*}
$$

(3) It holds the integral bound

$$
\int_{\Gamma} \operatorname{Tot} \operatorname{Var}_{[0,1)} \gamma_{a} d \omega_{h}(\gamma)<\infty .
$$

Similarly, we say that $\omega_{e} \in \mathcal{M}(\Gamma)$ is a Lagrangian representation of the epigraph of $u$ on $B_{R}$ if conditions (2) and (3) hold and (1) is replaced by

$$
\left(e_{t}\right)_{\sharp}\left[\omega_{e\llcorner } \Gamma(t)\right]=\mathscr{L}_{\llcorner }^{3} E_{\phi} \quad \text { for every } t \in(0,1) .
$$

In the following, we will adopt the slight abuse of notation

$$
\left(e_{t}\right)_{\sharp} \omega_{h}:=\left(e_{t}\right)_{\sharp}\left(\omega_{h\llcorner } \Gamma(t)\right) .
$$

A fundamental property of the Lagrangian representations $\omega_{h}, \omega_{e}$ above is that it is possible to decompose the Radon measure $U_{\phi}$ along the characteristic curves.

Given $\gamma \in \Gamma$, we consider

$$
\mu_{y}=(\mathrm{I}, \gamma)_{\#} \tilde{D}_{t} \gamma_{a}+\mathscr{H}^{1}{ }_{L} E_{\gamma}^{+}-\mathscr{H}^{1}{ }_{L} E_{\gamma}^{-} \in \mathcal{M}\left((0,1) \times B_{R} \times[0, M]\right),
$$

where

$$
\begin{aligned}
& E_{\gamma}^{+}:=\left\{(t, x, a): \gamma_{x}(t)=x, \gamma_{a}(t-)<\gamma_{a}(t+), a \in\left(\gamma_{a}(t-), \gamma_{a}(t+)\right)\right\}, \\
& E_{\gamma}^{-}:=\left\{(t, x, a): \gamma_{x}(t)=x, \gamma_{a}(t+)<\gamma_{a}(t-), a \in\left(\gamma_{a}(t+), \gamma_{a}(t-)\right)\right\},
\end{aligned}
$$

I : $[0,1) \rightarrow[0,1)$ denotes the identity and $\tilde{D}_{t} y_{a}$ denotes the diffuse part of the measure $D_{t} y_{a}$.
The main result of this section is the following theorem.
Theorem 3.2. Let $u \in \mathcal{M}_{\operatorname{div}}(\Omega)$ and $\phi \in L^{\infty}(\Omega)$ as in Definition 1.1. Let $B_{R}$ be an open ball of radius $R$ such that $\overline{B_{R}} \subset \Omega$ and $\mathscr{H}^{1}$-a.e. $x \in \partial B_{R}$ is a Lebesgue point of $\phi$. Then there exist Lagrangian representations $\omega_{h}, \omega_{e}$ of the hypograph and of the epigraph of $u$, respectively, on $B_{R}$ enjoying the additional properties:

$$
\begin{align*}
& \int_{\Gamma} \mu_{\gamma} d \omega_{h}(\gamma)=\mathscr{L}^{1} \times U_{\phi}=-\int_{\Gamma} \mu_{\gamma} d \omega_{e}(\gamma),  \tag{3.2}\\
& \int_{\Gamma}\left|\mu_{\gamma}\right| d \omega_{h}(\gamma)=\mathscr{L}^{1} \times\left|U_{\phi}\right|=\int_{\Gamma}\left|\mu_{\gamma}\right| d \omega_{e}(\gamma) . \tag{3.3}
\end{align*}
$$

Equations (3.2) and (3.3) are equalities in the space $\mathcal{N}\left((0,1) \times B_{R} \times[0, M]\right)$; equation (3.2) asserts that the measure $\mathscr{L}^{1} \times U_{\phi}$ can be decomposed along characteristics and equation (3.3) says that it can be done minimizing

$$
\int_{\Gamma} \operatorname{Tot}^{\operatorname{Var}}(0,1) \gamma_{a} d \omega_{h}(\gamma) \quad \text { and } \quad \int_{\Gamma} \operatorname{Tot}^{\operatorname{Var}}(0,1) \gamma_{a} d \omega_{e}(\gamma) .
$$

Moreover, it follows from (3.2) and (3.3) that we can separately represent the negative and the positive parts of $\mathscr{L}^{1} \times U_{\phi}$ in terms of the negative and positive parts of the measures $\mu_{\gamma}$ :

$$
\begin{equation*}
\int_{\Gamma} \mu_{\gamma}^{-} d \omega_{h}(\gamma)=\mathscr{L}^{1} \times U_{\phi}^{-}=\int_{\Gamma} \mu_{\gamma}^{+} d \omega_{e}(\gamma) \quad \text { and } \quad \int_{\Gamma} \mu_{\gamma}^{+} d \omega_{h}(\gamma)=\mathscr{L}^{1} \times U_{\phi}^{+}=\int_{\Gamma} \mu_{\gamma}^{-} d \omega_{e}(\gamma) \tag{3.4}
\end{equation*}
$$

The proof of Theorem 3.2 follows the strategy used in [18] to deal with general conservation laws; some additional work is required to obtain representations of solutions defined on $B_{R}$ and not on the whole Euclidean space.

### 3.2 An $L^{1}$-transport estimate

In this section, we prove an $L^{1}$-transport estimate that will be used as building block in the construction of approximate characteristics. First, we need the following lemma.
Lemma 3.3. Let $\bar{x} \in \Omega$ and let $B_{R}$ be an open ball of radius $R$ centered at $\bar{x}$ such that $\overline{B_{R}} \subset \Omega$ and $\mathscr{H}^{1}$-a.e. $x \in \partial B_{R}$ is a Lebesgue point of $\phi$. Let $\bar{t}>0$ be such that $\bar{t}<\operatorname{dist}\left(B_{R}, \partial \Omega\right)$ and let $\chi, U_{\phi}$ be as in (1.3). We define $\chi^{1}, \chi^{2}:[0, \bar{t}] \times \Omega \times[0, M] \rightarrow\{0,1\}$ by

$$
\chi^{1}(t, x, a)=\chi(x, a) \mathbf{1}_{B_{R}}(x) \quad \text { and } \quad \chi^{2}(t, x, a)=\chi\left(x-i e^{i a} t, a\right) \mathbf{1}_{B_{R}}(x)
$$

Then there exist two Radon measure $\mu_{\bar{t}}^{1}, \mu_{t}^{2} \in \mathcal{M}([0, \bar{t}] \times \Omega \times[0, M])$ absolutely continuous with respect to $\mathscr{H}^{3}{ }_{L}\left([0, \bar{t}] \times \partial B_{R} \times[0, M]\right)$ such that

$$
\left\{\begin{array}{l}
\partial_{t} \chi^{1}+i e^{i a} \cdot \nabla_{x} \chi^{1}=-\partial_{a}\left(\mathbf{1}_{[0, T] \times B_{R} \times[0, M]} U_{\phi}\right)+\mu_{\bar{t}}^{1}  \tag{3.5}\\
\partial_{t} \chi^{2}+i e^{i a} \cdot \nabla_{x} \chi^{2}=\mu_{\bar{t}}^{2} \\
\varepsilon_{\bar{t}}:=\frac{\left\|\mu_{\bar{t}}^{1}-\mu_{\bar{t}}^{2}\right\|}{\bar{t}} \rightarrow 0 \quad \text { as } \bar{t} \rightarrow 0
\end{array}\right.
$$

Proof. Let $\delta \in(0, R)$ and let $\psi_{\delta} \in C^{1}([0,+\infty) ; \mathbb{R})$ be such that

$$
\psi_{\delta} \equiv 1 \text { in }\left(0, R_{\delta}\right), \quad \psi_{\delta} \equiv 0 \text { in }(R,+\infty), \quad\left\|\psi_{\delta}^{\prime}\right\|_{C^{0}} \leq \frac{2}{\delta}
$$

Let us consider $\chi_{\delta}^{1}(t, x, a):=\chi(x, a) \psi_{\delta}(|x-\bar{x}|)$. Since $\mathscr{H}^{1}$-a.e. $x \in \partial B_{R}$ is a Lebesgue point of $\phi$, we have that $\mathscr{H}^{2}$-a.e. $(x, a) \in \partial B_{R} \times[0, M]$ is a Lebesgue point of $\chi$. Therefore, testing (1.3) with $\psi_{\delta}(|\cdot-\bar{x}|)$ and letting $\delta \rightarrow 0^{+}$, we get

$$
i e^{i a} \cdot \nabla_{\chi}\left(\chi(x, a) \mathbf{1}_{B_{R}}(x)\right)=-\partial_{a}\left(\mathbf{1}_{B_{R} \times[0, M]} U_{\phi}\right)+g \mathscr{H}_{\llcorner }^{2}\left(\partial B_{R} \times[0, M]\right)
$$

where $g(x, a):=i e^{i a} \cdot n(x) \chi(x, a)$ and $n$ denotes the inner normal to $B_{R}$. Since $\chi^{1}(t, x, a)=\chi(x, a) \mathbf{1}_{B_{R}}(x)$ for every $t \in[0, T]$, system (3.5) holds for

$$
\mu_{\bar{t}}^{1}=\bar{g} \mathscr{H}_{\mathrm{L}}^{3}\left([0, \bar{t}] \times \partial B_{R} \times[0, M]\right) \quad \text { with } \bar{g}(t, x, a)=i e^{i a} \cdot n(x) \chi(x, a)
$$

From the definition of $\chi^{2}$, the second equation in (3.5) holds with

$$
\mu_{\bar{t}}^{2}=i e^{i a} \cdot n(x) \chi\left(x-i e^{i a} t, a\right) \mathscr{H}_{\mathrm{L}}^{3}\left([0, \bar{t}] \times \partial B_{R} \times[0, M]\right)
$$

Since $\mathscr{H}^{2}$-a.e. $(x, a) \in \partial B_{R} \times[0, M]$ is a Lebesgue point of $\chi$, for $\mathcal{L}^{1}$-a.e. $a \in[0, M]$ we have that $\mathscr{H}^{1}$-a.e. $x \in \partial B_{R}$ is a Lebesgue point of $\chi(\cdot, a)$. In particular, for $\mathcal{L}^{1}$-a.e. $a \in[0, M]$ it holds

$$
\begin{equation*}
\lim _{\bar{t} \rightarrow 0} \frac{1}{\bar{t}} \int_{0}^{\bar{t}} \int_{\partial B_{R}}\left|\chi(x, a)-\chi\left(x-i e^{i a} t, a\right)\right| d \mathscr{H}^{1}(x) d t=0 \tag{3.6}
\end{equation*}
$$

Since for every $t \in[0, \bar{t}]$ and every $a \in[0, M]$ it holds

$$
\int_{\partial B_{R}}\left|\chi(x, a)-\chi\left(x-i e^{i a} t, a\right)\right| d \mathscr{H}^{1}(x) \leq\left|\partial B_{R}\right|
$$

by integrating (3.6) with respect to $a$, it follows by the dominated convergence theorem that

$$
\left\|\mu_{\bar{t}}^{1}-\mu_{\bar{t}}^{2}\right\|=\int_{0}^{\bar{t}} \int_{0}^{M} \int_{\partial B_{R}}\left|\chi(x, a)-\chi\left(x-i e^{i a} t, a\right)\right| d \mathscr{H}^{1}(x) d a d t=o(\bar{t}) \quad \text { as } \bar{t} \rightarrow 0
$$

since $\mathscr{H}^{1}$-a.e. $x \in \partial B_{R}$ is a Lebesgue point of $\phi$, and therefore $\mathscr{H}^{2}$-a.e. $(x, a) \in \partial B_{R} \times[0, M]$ is a Lebesgue point of $\chi$.

Proposition 3.4. In the setting of Lemma 3.3, let $\psi \in C_{c}^{1}(\Omega \times \mathbb{R})$. Then

$$
\int_{\Omega \times \mathbb{R}} \psi(x, a)\left(\chi^{1}(\bar{t})-\chi^{2}(\bar{t})\right) d x d a \leq\left(\bar{t}\left\|\partial_{a} \psi\right\|_{L^{\infty}}+\frac{\bar{t}^{2}}{2}\left\|\nabla_{x} \psi\right\|_{L^{\infty}}\right) v\left(B_{R}\right)+\|\psi\|_{L^{\infty}} \varepsilon_{\bar{t}} \bar{t}
$$

Proof. We set $\tilde{\chi}:=\chi^{1}-\chi^{2}$ and $\tilde{\psi}(t, x, a):=\psi\left(\chi+i e^{i a}(\bar{t}-t), a\right)$. It is straightforward to check that

$$
\begin{equation*}
\partial_{t}(\tilde{\chi} \tilde{\psi})+i e^{i a} \cdot \nabla_{x}(\tilde{\chi} \tilde{\psi})=-\tilde{\psi} \partial_{a}\left(\mathscr{L}^{1} \times U_{\phi}\right)+\tilde{\psi}\left(\mu_{\bar{t}}^{1}-\mu_{t}^{2}\right) \quad \text { in } \mathcal{D}^{\prime}((0, \bar{t}) \times \Omega \times \mathbb{R}) \tag{3.7}
\end{equation*}
$$

Let $g:[0, \bar{t}] \rightarrow \mathbb{R}$ be defined by

$$
g(t)=\int_{\Omega \times \mathbb{R}} \tilde{\chi}(t) \tilde{\psi}(t) d x d a
$$

It follows from (3.7) that

$$
g^{\prime}(t)=-\int_{\Omega \times \mathbb{R}} \partial_{a} \tilde{\psi}(t) d U_{\phi}+\int_{\Omega \times \mathbb{R}} \tilde{\psi}(t) d\left(\mu_{\bar{t}}^{1}-\mu_{\bar{t}}^{2}\right)_{t}
$$

holds in the sense of distributions, where $\left(\mu_{\bar{t}}^{1}-\mu_{\bar{t}}^{2}\right)_{t}$ denotes the disintegration of the measure $\mu_{\bar{t}}^{1}-\mu_{\bar{t}}^{2}$ in $t \in(0, \bar{t})$ with respect to $\mathscr{L}^{1}\left\llcorner(0, \bar{t})\right.$. Therefore, $g \in C^{1}([0, \bar{t}])$. Since $g(0)=0$, it holds

$$
\begin{aligned}
\int_{\Omega \times \mathbb{R}} \psi\left(\chi^{1}(\bar{t})-\chi^{2}(\bar{t})\right) d x d a & =g(\bar{t})-g(0) \\
& =\int_{0}^{\bar{t}} g^{\prime}(t) d t \\
& =-\int_{0}^{\bar{t}} \int_{\Omega \times \mathbb{R}} \partial_{a} \tilde{\psi}(t) d U_{\phi} d t+\int_{(0, \bar{t}) \times \Omega \times \mathbb{R}} \tilde{\psi} d\left(\mu_{\bar{t}}^{1}-\mu_{\bar{t}}^{2}\right) \\
& =-\int_{0}^{\bar{t}} \int_{\Omega \times \mathbb{R}}\left(\partial_{v} \phi-(\bar{t}-t) e^{i a} \cdot \nabla_{x} \psi\right) d U_{\phi} d t+\int_{(0, \bar{t}) \times \Omega \times \mathbb{R}} \tilde{\psi} d\left(\mu_{\bar{t}}^{1}-\mu_{\bar{t}}^{2}\right) \\
& \leq\left(\bar{t}\left\|\partial_{a} \psi\right\|_{L^{\infty}}+\frac{\bar{t}^{2}}{2}\left\|\nabla_{x} \psi\right\|_{L^{\infty}}\right) v\left(B_{R}\right)+\|\psi\|_{L^{\infty}}\left\|\mu_{\bar{t}}^{1}-\mu_{\bar{t}}^{2}\right\|
\end{aligned}
$$

and this concludes the proof.
We set $L_{\bar{t}}=\left(\varepsilon_{\bar{t}} \vee \bar{t}\right)^{-\frac{1}{2}}$ and we consider the anisotropic distance

$$
\begin{aligned}
& d_{\bar{t}}:\left(B_{R} \times[0, M]\right)^{2} \rightarrow[0,+\infty) \\
& \left(\left(x_{1}, a_{1}\right),\left(x_{2}, a_{2}\right)\right) \mapsto L_{\bar{t}}\left|x_{1}-x_{2}\right|+\left|a_{1}-a_{2}\right|
\end{aligned}
$$

A test function $\psi: B_{R} \times[0, M] \rightarrow \mathbb{R}$ is 1-Lipschitz with respect to $d_{\bar{t}}$ if and only if

$$
\left\|\partial_{a} \psi\right\|_{L^{\infty}} \leq 1 \quad \text { and } \quad\left\|\nabla_{x} \psi\right\|_{L^{\infty}} \leq L_{\bar{t}}
$$

Applying Corollary 2.3 to $\mu^{1}=\chi^{1}(\bar{t}) \mathscr{L}^{3}$ and $\mu^{2}=\chi^{2}(\bar{t}) \mathscr{L}^{3}$ on the space ( $B_{R} \times[0, M], d_{\bar{t}}$ ), we obtain the following result as a consequence of Proposition 3.4 and Theorem 2.4.

Corollary 3.5. There exist $\rho_{\bar{t}}^{1} \leq \chi^{1}(\bar{t})$ and $\rho_{\bar{t}}^{2} \leq \chi^{2}(\bar{t})$ such that

$$
\int_{B_{R} \times[0, M]}\left(\chi^{1}(\bar{t})-\rho_{\bar{t}}^{1}\right) d x d a \leq \varepsilon_{\bar{t}} \bar{t}, \quad \int_{B_{R} \times[0, M]}\left(\chi^{2}(\bar{t})-\rho_{\bar{t}}^{2}\right) d x d a \leq \varepsilon_{\bar{t}} \bar{t}
$$

and

$$
W_{1}\left(\rho_{\bar{t}}^{1} \mathscr{L}^{3}, \rho_{\bar{t}}^{2} \mathscr{L}^{3}\right) \leq\left(\bar{t}+\bar{t}^{\frac{3}{2}}\right) v\left(B_{R}\right)+\varepsilon_{\bar{t}}^{\frac{1}{2}} \bar{t}\left(2 R+\varepsilon_{\bar{t}}^{\frac{1}{2}} M\right)
$$

In particular, there exists

$$
T=\left(T_{x}, T_{a}\right): B_{R} \times[0, M] \rightarrow B_{R} \times[0, M]
$$

such that $T_{\sharp}\left(\rho_{\bar{t}}^{2} \mathscr{L}^{3}\right)=\rho_{\bar{t}}^{1} \mathscr{L}^{3}$ and

$$
\begin{equation*}
\int_{B_{R} \times[0, M]}^{t}\left(L_{\bar{t}}\left|T_{x}(x, a)-x\right|+\left|T_{a}(x, a)-a\right|\right) \rho_{\bar{t}}^{2}(x, a) d x d a \leq\left(\bar{t}+\bar{t}^{\frac{3}{2}}\right) v\left(B_{R}\right)+\varepsilon_{\bar{t}}^{\frac{1}{2}} \bar{t}\left(2 R+\varepsilon_{\bar{t}}^{\frac{1}{2}} M\right) \tag{3.8}
\end{equation*}
$$

Remark 3.6. Now, we observe that in order to build a Lagrangian representation as in Theorem 3.2, the use of Theorem 2.4 can be replaced by a more elementary argument: indeed, the infimum in (2.1) can be equivalently taken only on the plans induced by transport maps (see, for example, [3]). In particular, the second part of the statement in Corollary 3.5 can be replaced by the following slightly weaker version: for every $\varepsilon^{\prime}>0$, there exists a map

$$
T=\left(T_{x}, T_{a}\right): B_{R} \times[0, M] \rightarrow B_{R} \times[0, M]
$$

such that $T_{\sharp}\left(\rho_{\bar{t}}^{2} \mathscr{L}^{3}\right)=\rho_{\bar{t}}^{1} \mathscr{L}^{3}$ and

$$
\int_{B_{R} \times[0, M]}^{*}\left(L_{\bar{t}}\left|T_{x}(x, a)-x\right|+\left|T_{a}(x, a)-a\right|\right) \rho_{\bar{t}}^{2}(x, a) d x d a \leq\left(\bar{t}+\bar{t}^{\frac{3}{2}}\right) v\left(B_{R}\right)+\varepsilon_{\bar{t}}^{\frac{1}{2}} \bar{t}\left(2 R+\varepsilon_{\bar{t}}^{\frac{1}{2}} M\right)+\varepsilon^{\prime} .
$$

The only property that we will use of (3.8) is that the right-hand side is of the form $\bar{t} v\left(B_{R}\right)+o(\bar{t})$ as $\bar{t} \rightarrow 0$. In particular, choosing $\varepsilon^{\prime}=o(\bar{t})$, we can avoid the use of Theorem 2.4.

### 3.3 Construction of approximate characteristics

### 3.3.1 Building block

For a fixed $\bar{t}>0$, we consider the following sets:

$$
\begin{aligned}
& E_{1}:=\left\{(x, a) \in B_{R} \times[0, M]: x+i e^{i a} \bar{t} \in B_{R}\right\}, \\
& E_{2}:=\left\{(x, a) \in B_{R} \times[0, M]: x+i e^{i a} \bar{t} \notin B_{R}\right\}, \\
& E_{3}:=\left\{(x, a) \in\left(\Omega \backslash B_{R}\right) \times[0, M]: x+i e^{i a} \bar{t} \in B_{R}\right\} .
\end{aligned}
$$

For every $(x, a) \in E_{1}$, we define $\gamma_{\bar{t}, x, a}:[0, \bar{t}] \rightarrow B_{R} \times[0, M]$ by

$$
y_{\bar{t}, x, a}(t)= \begin{cases}\left(x+i e^{i a} t, a\right) & \text { if } t \in[0, \bar{t}), \\ T\left(x+i e^{i a} \bar{t}, a\right) & \text { if } t=\bar{t}\end{cases}
$$

where the transport map $T$ is defined in Corollary 3.5. For every $(x, a) \in E_{2}$, we set

$$
t^{+}(x, a):=\sup \left\{t \in[0, \bar{t}]: x+i e^{i a} t \in B_{R}\right\}
$$

and we define $\gamma_{\bar{t}, x, a}:\left[0, t^{+}(x, a)\right) \rightarrow B_{R} \times[0, M]$ by

$$
\gamma_{\bar{t}, x, a}(t)=\left(x+i e^{i a} t, a\right)
$$

For every $(x, a) \in E_{3}$, we set

$$
t^{-}(x, a):=\inf \left\{t \in[0, \bar{t}]: x+i e^{i a} t \in B_{R}\right\}
$$

and we define $\gamma_{\bar{t}, x, a}:\left(t^{-}(x, a), \bar{t}\right] \rightarrow B_{R} \times[0, M]$ by

$$
y_{\bar{t}, x, a}(t)= \begin{cases}\left(x+i e^{i a} t, a\right) & \text { if } t \in\left(t^{-}(x, a), \bar{t}\right) \\ T\left(x+i e^{i a} \bar{t}, a\right) & \text { if } t=\bar{t}\end{cases}
$$

### 3.3.2 Approximate characteristics

Fix $n \in N$ and set $\bar{t}_{n}=2^{-n}$. For every $(x, a) \in E_{2}$, we consider the curve

$$
\gamma_{x, a}^{0, n}:\left(t_{\gamma_{x, a}^{0, n}}^{-}, t_{\gamma_{x, a}^{0, n}}^{+}\right) \rightarrow B_{R} \times[0, M]
$$

with

$$
t_{\gamma_{x, a}^{0, n}}^{-}=0, \quad t_{\gamma_{x, a}^{0, n}}^{+}=t^{+}(x, a), \quad y_{x, a}^{0, n}(t)=\gamma_{2^{-n}, x, a}(t) \quad \text { for all } t \in\left(t_{\gamma_{x, a}^{0, n}}^{-}, t_{\gamma_{x, a}^{0, n}}^{+}\right)
$$

For every $(x, a) \in E_{1}$, we define

$$
\gamma_{x, a}^{0, n}:\left(t_{\gamma_{x, a}^{0, n}}^{-}, t_{\gamma_{x, a}^{0, n}}^{+}\right) \rightarrow B_{R} \times[0, M]
$$

with

$$
t_{y_{x, a}^{0, n}}^{-}=0, \quad t_{y_{x, a}^{0, n}}^{+} \geq 2^{-n}
$$

to be determined in the construction and

$$
\gamma_{x, a}^{0, n}(t)=\gamma_{2^{-n}, x, a}(t) \quad \text { for all } t \in\left(t_{\gamma_{x, a}^{0, n}}^{-}, 2^{-n}\right]
$$

For every $k=1, \ldots, 2^{n}$ and for every $(x, a) \in E_{3}$, we introduce a curve

$$
\gamma_{x, a}^{k, n}:\left(t_{\gamma_{x, a}^{k, n}}^{-}, t_{\gamma_{x, a}^{k, n}}^{+}\right) \rightarrow B_{R} \times[0, M]
$$

with

$$
t_{\gamma_{x, a}^{k, n}}^{-}=(k-1) 2^{-n}+t^{-}(x, a), \quad t_{\gamma_{x, a}^{k, n}}^{+} \geq k 2^{-n}
$$

to be determined and

$$
\gamma_{x, a}^{k, n}(t)=\gamma_{2^{-n}, x, a}\left(t-(k-1) 2^{-n}\right) \quad \text { for all } t \in\left(t_{\gamma_{x, a}^{k, n}}^{-}, k 2^{-n}\right]
$$

It remains to define the evolution of the curves $\gamma_{x, a}^{0, n}$ for $(x, a) \in E_{1}$ and $t \geq 2^{-n}$ and of the curves $\gamma_{x, a}^{k, n}$ for $(x, a) \in E_{3}$ and $t \geq k 2^{-n}$. Let us fix $k=1, \ldots, 2^{n}$ and $(x, a) \in E_{3}$. We define the evolution of $\gamma_{x, a}^{k, n}$ by recursion: assume that $y_{x, a}^{k, n}$ is defined on

$$
\left(t_{y_{x, a}^{k, n}}^{-}, l 2^{-n}\right] \quad \text { for some } l \geq k .
$$

If $l=2^{n}$, then we set

$$
t_{\gamma_{x, a}^{k, n}}^{+}=1
$$

Otherwise, if $l<2^{n}$, then we distinguish two cases.
If $y_{x, a}^{k, n}\left(l 2^{-n}\right) \in E_{2}$, then we set

$$
t_{\gamma_{x, a}^{k, n}}^{+}=l 2^{-n}+t^{+}\left(y_{x, a}^{k, n}\left(l 2^{-n}\right)\right)
$$

and

$$
\gamma_{x, a}^{k, n}(t)=\gamma_{2^{-n}, \gamma_{x, a}^{k, n}\left(12^{-n)}\right.}\left(t-l 2^{-n}\right) \quad \text { for all } t \in\left(l 2^{-n}, t_{\gamma_{x, a}^{k, n}}^{+}\right)
$$

If instead $\gamma_{x, a}^{k, n}\left(l 2^{-n}\right) \in E_{1}$, then we extend $\gamma_{x, a}^{k, n}$ on the whole interval $\left(l 2^{-n},(l+1) 2^{-n}\right]$ by setting

$$
\gamma_{x, a}^{k, n}(t)=\gamma_{2^{-n}, \gamma_{x, a}^{k, n}\left(l 2^{-n}\right)}\left(t-l 2^{-n}\right) \quad \text { for all } t \in\left(l 2^{-n},(l+1) 2^{-n}\right]
$$

The extension of the curves $\gamma_{x, a}^{0, n}$ for $(x, a) \in E_{1}$ is defined by the same procedure described above for the curves $y_{x, a}^{k, n}$ for $(x, a) \in E_{3}$ with $k=1$.

### 3.4 Approximate Lagrangian representation

The approximate characteristics built in the previous section belong to the space

$$
\tilde{\Gamma}:=\left\{\left(\gamma, t_{y}^{-}, t_{\gamma}^{+}\right): 0 \leq t_{\gamma}^{-} \leq t_{\gamma}^{+} \leq 1, \gamma \in \mathrm{BV}\left(\left(t_{y}^{-}, t_{\gamma}^{+}\right) ; B_{R} \times[0, M]\right)\right\}
$$

For every $n \in \mathbb{N}$ sufficiently large, we define $\omega_{n} \in \mathcal{M}(\tilde{\Gamma})$ by
where the curves $\gamma_{x, a}^{k, n}$ are defined in Section 3.3.2.
Lemma 3.7. Let $\omega_{n}$ be defined in (3.9). Then the following estimates hold:

$$
\begin{array}{ll}
e_{h}(n):=\int_{\tilde{\Gamma}} \sup _{t \in\left(t_{y}^{-}, t_{y}^{+}\right)}\left|y_{x}(t)-\gamma_{x}\left(t_{\gamma}^{-}\right)-\int_{t_{\bar{\gamma}}^{-}}^{t} i e^{i y_{a}(s)} d s\right| d \omega_{n}(\gamma)=o(1) & \text { as } n \rightarrow \infty, \\
e_{\nu}(n):=\int_{\tilde{\Gamma}} \operatorname{Tot}^{V^{2}}{ }_{\left(t_{y}^{-}, t_{y}^{+}\right)} \gamma_{a} d \omega_{n}(\gamma) \leq v\left(B_{R}\right)+o(1) & \text { as } n \rightarrow \infty . \tag{3.11}
\end{array}
$$

Proof. Since for $\omega_{n}$-a.e. $\left(\gamma, t_{\gamma}^{-}, t_{\gamma}^{+}\right) \in \tilde{\Gamma}$ it holds

$$
\dot{\gamma}_{x}(t)=i e^{i y_{a}(t)} \quad \text { for all } t \in\left(y, t_{\gamma}^{-}, t_{\gamma}^{+}\right) \backslash 2^{-n} \mathbb{N}
$$

we have

$$
\begin{align*}
&\left.\sup _{t \in\left(t_{y},\right.}, t_{y}^{+}\right) \\
&\left|\gamma_{x}(t)-\gamma_{x}\left(t_{\gamma}^{-}\right)-\int_{t_{y}^{-}}^{t} i e^{i y_{a}(s)} d s\right| \leq \sum_{l^{-(\gamma)}}^{l^{+}(\gamma)}\left|\gamma_{x}\left(l 2^{-n}\right)-\gamma_{x}\left(l 2^{-n}-\right)\right|  \tag{3.12}\\
&=\sum_{l^{-}(\gamma)}^{l^{+}(\gamma)} \mid T_{x}\left(\gamma\left(l 2^{-n}-\right)-\gamma_{x}\left(l 2^{-n}-\right) \mid\right.
\end{align*}
$$

where

$$
\begin{aligned}
& l^{-}(y)=2^{n} \inf \left(2^{-n} \mathbb{Z} \cap\left(t_{\gamma}^{-}, t_{\gamma}^{+}\right)\right) \\
& l^{+}(y)=2^{n} \sup \left(2^{-n} \mathbb{Z} \cap\left(t_{\gamma}^{-}, t_{\gamma}^{+}\right)\right)
\end{aligned}
$$

Integrating (3.12) with respect to $\omega_{n}$, it follows by Corollary 3.5 with $\bar{t}=2^{-n}$ that

$$
\begin{align*}
e_{h}(n) & \leq \sum_{l=1}^{2^{n}-1} \int_{X}\left|T_{x}(x, a)-x\right| d\left(e_{l 2^{-n}-}\right)_{\sharp} \omega_{n} \\
& \leq \sum_{l=1}^{2^{n}-1}\left(\int_{X}\left|T_{x}(x, a)-x\right| \rho_{\bar{t}}^{2}(x, a) d x d a+2 R\left\|\left(\left(e_{l 2^{-n_{-}}}\right)_{\sharp} \omega_{n}-\rho_{\bar{t}}^{2} \mathscr{L}^{3}\right)^{+}\right\|\right) \\
& \leq \frac{2^{n}}{L_{2^{-n}}}\left(2^{-n}+2^{\frac{-3 n}{2}}\right) v\left(B_{R}\right)+\varepsilon_{2^{-n}}^{\frac{1}{2}}\left(2 R+\varepsilon_{2^{-n}}^{\frac{1}{2}} M\right)+2 R \sum_{l=1}^{2^{n}-1}\left\|\left(\left(e_{l 2^{-n-}}\right)_{\sharp} \omega_{n}-\rho_{\bar{t}}^{2} \mathscr{L}^{3}\right)^{+}\right\|, \tag{3.13}
\end{align*}
$$

where $X=B_{R} \times[0, M]$ and $e_{t-}: \tilde{\Gamma}(t) \rightarrow X$ is defined by $e_{t_{-}}(\gamma)=\lim _{t^{\prime} \rightarrow t_{-}} \gamma\left(t^{\prime}\right)$. Since, by construction,

$$
\left(e_{l 2^{-n}-}\right)_{\sharp} \omega_{n} \leq \chi^{2}(\bar{t}) \mathscr{L}^{3} \quad \text { and } \quad \rho_{\bar{t}}^{2} \leq \chi^{2}(\bar{t})
$$

with

$$
\left\|\left(\chi^{2}(\bar{t})-\rho_{\bar{t}}^{2}\right) \mathscr{L}^{3}\right\| \leq 2^{-n} \varepsilon_{2^{-n}}
$$

for every $l=1, \ldots, 2^{n}-1$ it holds

$$
\begin{equation*}
\left\|\left(\left(e_{l 2^{-n}-}\right)_{\sharp} \omega_{n}-\rho_{\bar{t}}^{2} \mathscr{L}^{3}\right)^{+}\right\| \leq 2^{-n} \varepsilon_{2^{-n}} . \tag{3.14}
\end{equation*}
$$

Plugging (3.14) into (3.13), we immediately get (3.10).

Now, we prove (3.11). Since $\gamma_{a}$ is constant in each connected component of $\left(t_{\gamma}^{-}, t_{\gamma}^{+}\right) \backslash 2^{-n} \mathbb{N}$ for $\omega_{n}$-a.e. $\gamma$, it follows by Corollary 3.5 that

$$
\begin{aligned}
\int_{\tilde{\Gamma}} \operatorname{Totgar}_{\left(t_{\nu}^{-}, t_{y}^{+}\right)} \gamma_{a} d \omega_{n}(\gamma) & =\sum_{l=1}^{2^{n}-1} \int_{\tilde{\Gamma}\left(l 2^{-n}\right)}\left|T_{a}\left(y\left(l 2^{-n}-\right)\right)-\gamma_{a}\left(l 2^{-n}-\right)\right| d \omega_{n}(\gamma) \\
& =\sum_{l=1}^{2^{n}-1} \int_{X}\left|T_{a}(x, a)-a\right| d\left(e_{l 2^{-n}-}\right)_{\sharp} \omega_{n} \\
& \leq \sum_{l=1}^{2^{n}-1} \int_{X}\left|T_{a}(x, a)-a\right| \rho_{\bar{t}}^{2}(x, a) d x d a+M\left\|\left(\left(e_{l 2^{-n}-}\right)_{\sharp} \omega_{n}-\rho_{\bar{t}}^{2} \mathscr{L}^{3}\right)^{+}\right\| \\
& \leq 2^{n}\left[\left(2^{-n}+2^{\frac{-3 n}{2}}\right) v\left(B_{R}\right)+\varepsilon_{2^{-n}}^{\frac{1}{2}} \bar{t}\left(2 R+\varepsilon_{2^{-n}}^{\frac{1}{2}} M\right)\right]+M \varepsilon_{2^{-n}},
\end{aligned}
$$

which implies (3.11).
Now, we show that $\left(e_{t}\right)_{\sharp} \omega_{n}$ approximates $\chi \mathscr{L}^{3}$ in the strong topology of measures for every $t \in 2^{-n} \mathbb{N} \cap[0,1)$. This property and the weak continuity estimate provided in Proposition 3.4 will guarantee Definition 3.1 (1).

Lemma 3.8. For every $l=0, \ldots, 2^{n}-1$, it holds

$$
\begin{equation*}
\left\|\left(e_{l 2^{-n}}\right)_{\sharp} \omega_{n}-\chi \mathscr{L}^{3}\right\| \leq 2^{-n+1} l \varepsilon_{2^{-n}} . \tag{3.15}
\end{equation*}
$$

Moreover, for every $t \in\left[l 2^{-n},(l+1) 2^{-n}\right)$ and every $\psi \in C_{c}^{\infty}\left(B_{R} \times[0, M]\right)$, it holds

$$
\begin{equation*}
\left|\int_{X} \psi d\left(e_{t}\right)_{\sharp} \omega_{n}-\int_{X} \psi d\left(e_{l 2^{-n}}\right)_{\sharp} \omega_{n}\right| \leq 2^{-n}\left(2 M \mathscr{H}^{1}\left(\partial B_{R}\right)\|\psi\|_{L^{\infty}}+\|\nabla \psi\|_{L^{\infty}} \mathscr{L}^{3}\left(H_{\phi}\right)\right) . \tag{3.16}
\end{equation*}
$$

Proof. First, the case $l=0$ follows by the definition of $\omega_{n}$. In order to get (3.15), we prove that for every $l=0, \ldots, 2^{n}-2$ it holds

$$
\left\|\left(e_{(l+1) 2^{-n}}\right)_{\sharp} \omega_{n}-\chi \mathscr{L}^{3}\right\| \leq\left\|\left(e_{l 2^{-n}}\right)_{\sharp} \omega_{n}-\chi \mathscr{L}^{3}\right\|+2^{-n+1} \varepsilon_{2^{-n}} .
$$

Indeed,

$$
\begin{aligned}
\left\|\left(e_{(l+1) 2^{-n}}\right)_{\sharp} \omega_{n}-\chi \mathscr{L}^{3}\right\| & \leq\left\|\left(e_{(l+1) 2^{-n}}\right)_{\sharp} \omega_{n}-\rho_{\bar{t}}^{1} \mathscr{L}^{3}\right\|+\left\|\rho_{\bar{t}}^{1} \mathscr{L}^{3}-\chi \mathscr{L}^{3}\right\| \\
& =\left\|T_{\sharp}\left(e_{(l+1) 2^{-n}-}\right)_{\sharp} \omega_{n}-T_{\sharp}\left(\rho_{\bar{t}}^{2} \mathscr{L}^{3}\right)\right\|+\left\|\rho_{\bar{t}}^{1} \mathscr{L}^{3}-\chi \mathscr{L}^{3}\right\| \\
& \leq\left\|\left(e_{(l+1) 2^{-n}-}\right)_{\sharp} \omega_{n}-\rho_{\bar{t}}^{2} \mathscr{L}^{3}\right\|+2^{-n} \varepsilon_{2^{-n}} \\
& \left.\leq \|\left(e_{(l+1) 2^{-n}-}\right)_{\sharp} \omega_{n}-\chi^{2}(\bar{t}) \mathscr{L}^{3}\right)\|+\|\left(\chi^{2}(\bar{t})-\rho_{\bar{t}}^{2} \mathscr{L}^{3}\right) \|+2^{-n} \varepsilon_{2^{-n}} \\
& \leq\left\|\left(e_{l 2^{-n}}\right)_{\sharp} \omega_{n}-\chi \mathscr{L}^{3}\right\|+2 \cdot 2^{-n} \varepsilon_{2^{-n}} .
\end{aligned}
$$

Inequality (3.16) follows by

$$
\begin{aligned}
\left|\int_{X} \psi d\left(e_{t}\right)_{\sharp} \omega_{n}-\int_{X} \psi d\left(e_{l 2^{-n}}\right)_{\sharp} \omega_{n}\right| \leq & \left|\int_{X} \psi d\left(e_{t}\right)_{\sharp} \omega_{n\llcorner }\left\{t_{y}^{-}>l 2^{-n}\right\}\right|+\left|\int_{X} \psi d\left(e_{l 2^{-n}}\right)_{\sharp} \omega_{n L}\left\{t_{\gamma}^{+}<t\right\}\right| \\
& +\left|\int_{X} \psi d\left(e_{t}\right)_{\sharp} \omega_{n L}\left\{t_{y}^{-} \leq l 2^{-n}\right\}-\int_{X} \psi d\left(e_{\left.l 2^{-n}\right)}\right)_{\sharp} \omega_{n L}\left\{t_{\gamma}^{+}>t\right\}\right| \\
\leq & 2 \cdot 2^{-n} \mathscr{H}^{1}\left(\partial B_{R}\right) M\|\psi\|_{L^{\infty}}+\|\nabla \psi\|_{L^{\infty}} 2^{-n} \omega_{n}(\tilde{\Gamma}(t)) \\
\leq & 2^{-n+1} \mathscr{H}^{1}\left(\partial B_{R}\right) M\|\psi\|_{L^{\infty}}+\|\nabla \psi\|_{L^{\infty}} 2^{-n} \mathscr{L}^{3}\left(H_{\phi}\right),
\end{aligned}
$$

as desired.

### 3.5 Compactness of $\omega_{n}$ and existence of a Lagrangian representation

We consider on $\tilde{\Gamma}$ the topology $\tau$ that induces the following convergence: $\left(\gamma_{n}, t_{\gamma_{n}}^{-}, t_{\gamma_{n}}^{+}\right)$converges to $\left(\gamma, t_{\gamma}^{-}, t_{\gamma}^{+}\right)$if $t_{\gamma_{n}}^{ \pm} \rightarrow t_{\gamma}^{ \pm}$with respect to the Euclidean topology in $\mathbb{R}$ and there exist extensions $\tilde{\gamma}, \tilde{\gamma}_{n}$ of $\gamma, \gamma_{n}$ defined on $(0,1)$
such that the horizonal components $\tilde{\gamma}_{n, x}$ converge to $\tilde{\gamma}_{x}$ uniformly and the vertical components $\tilde{y}_{n, a}$ converge to $\tilde{\gamma}_{a}$ in $L^{1}(0,1)$.
Lemma 3.9. The sequence of measures $\omega_{n}$ defined in (3.9) is bounded and tight in $\mathcal{M}(\tilde{\Gamma})$, namely for every $\varepsilon>0$ there exists a compact $K_{\varepsilon} \subset \tilde{\Gamma}$ such that for every $n \in \mathbb{N}$ it holds

$$
\omega_{n}\left(\tilde{\Gamma} \backslash K_{\varepsilon}\right)<\varepsilon .
$$

Proof. We prove first that the sequence $\omega_{n}$ is bounded: for every $n$ it holds

$$
\left|E_{3, n} \cap H_{\phi}\right| \leq\left|E_{3, n}\right| \leq M \mathscr{H}^{1}\left(\partial B_{R}\right) 2^{-n} .
$$

In particular,

$$
\underset{n \rightarrow \infty}{\limsup }\left|\omega_{n}\right|(\tilde{\Gamma})=\limsup _{n \rightarrow \infty} \mathscr{L}^{3}\left(H_{\phi}\right)+2^{n}\left|E_{3, n} \cap H_{\phi}\right| \leq \mathscr{L}^{3}\left(H_{\phi}\right)+M \mathscr{H}^{1}\left(\partial B_{R}\right) .
$$

In order to prove the tightness of the sequence $\omega_{n}$, we consider for every $n \in \mathbb{N}$ and $C>0$ the set of curves $\left(\gamma, t_{y}^{-}, t_{y}^{+}\right) \in \tilde{\Gamma}_{n, c} \subset \tilde{\Gamma}$ satisfying the following properties:
(i) $\operatorname{Totarar}_{\left(t_{\nu}, t_{\nu}^{+}\right)} \gamma_{a} \leq C$.
(ii) One has

$$
\sum_{k=l-(y)}^{l^{+}(y)}\left|y_{x}\left(2^{-n} k\right)-y_{x}\left(2^{-n} k-\right)\right| \leq C e_{h}(n)^{1 / 2}
$$

where $e_{h}(n)$ is defined in Lemma 3.7, and

$$
l^{-}(y):=2^{n} \inf 2^{-n} \mathbb{Z} \cap\left(t_{\gamma}^{-}, t_{\gamma}^{+}\right) \quad \text { and } \quad l^{+}(y):=2^{n} \sup 2^{-n} \mathbb{Z} \cap\left(t_{\gamma}^{-}, t_{y}^{+}\right) .
$$

(iii) Lip $y_{\llcorner }\left(\left[(k-1) 2^{-n}, k 2^{-n}\right) \leq 1\right.$ for every $k=l^{-}(y), \ldots, l^{+}(\gamma)$.

Since $e_{h}(n)$ tends to 0 as $n \rightarrow \infty$, for every $C>0$ the space

$$
\tilde{\Gamma}(C):=\bigcup_{n=1}^{\infty} \tilde{\Gamma}_{n, C}
$$

is compact with respect to the topology $\tau$ introduced above. Moreover, it follows by Lemma 3.7 and the Chebychev inequality that for every $\varepsilon>0$ there exists $C>0$ sufficiently large such that for every $n \in \mathbb{N}$,

$$
\omega_{n}(\tilde{\Gamma} \backslash \tilde{\Gamma}(C)) \leq \varepsilon
$$

By Theorem 2.5, it follows that the sequence $\omega_{n}$ is precompact with respect to the narrow convergence. We show in the next lemma that every limit point of $\omega_{n}$ is a Lagrangian representation of the hypograph of $\phi$ on $B_{R}$.
Lemma 3.10. Every limit point $\omega$ of the sequence $\omega_{n}$ is a Lagrangian representation of the hypograph of $\phi$ on $B_{R}$.

Proof. We need to check that the three conditions in Definition 3.1 are satisfied and that $\omega \in \mathcal{M}_{+}(\Gamma)$, namely that $\omega$ is concentrated on $\Gamma$.

Condition (1). We prove that for every $t \in(0,1)$ the following two limits hold in the sense of distributions:

$$
\begin{equation*}
\lim _{n \rightarrow \infty}\left(e_{t}\right)_{\sharp} \omega_{n}=\mathscr{L}^{3}{ }_{\llcorner } H_{\phi} \quad \text { and } \quad \lim _{n \rightarrow \infty}\left(e_{t}\right)_{\sharp} \omega_{n}=\left(e_{t}\right)_{\sharp} \omega \text {. } \tag{3.17}
\end{equation*}
$$

For every $t=2^{-k} \mathbb{N} \cap(0,1)$ for some $k \in \mathbb{N}$, the first limit holds true thanks to Lemma 3.8 since $\chi \mathscr{L}^{3}=\mathscr{L}^{3}{ }_{L} H_{\phi}$ by the definition of $\chi$. The continuity in time stated in (3.16) implies that the limit holds true therefore for every $t \in(0,1)$ in the sense of distributions. We observe that the second limit in (3.17) is not trivial since $e_{t}$ is not continuous on $\tilde{\Gamma}$ with respect to the topology introduced above. In order to establish it, we need to check that for every $\psi \in C_{c}^{\infty}\left(B_{R} \times[0, M]\right)$ it holds

$$
\lim _{n \rightarrow \infty} \int_{\tilde{\Gamma}(t)} \psi(\gamma(t)) d \omega_{n}=\int_{\Gamma}(t)
$$

Let $I \subset(0,1)$ be a non-empty open interval. Then consider the continuous and bounded function $T_{\psi, I}: \tilde{\Gamma} \rightarrow \mathbb{R}$ defined by

$$
T_{\psi, I}\left(\gamma, t_{\gamma}^{-}, t_{\gamma}^{+}\right):=\int_{I \cap\left(t_{\gamma}^{-}, t_{y}^{+}\right)} \psi(\gamma(t)) d t
$$

By the definition of narrow convergence and by the Fubini theorem, it follows that

$$
\lim _{n \rightarrow \infty} \int_{I} \int_{\tilde{\Gamma}(t)} \psi(\gamma(t)) d \omega_{n} d t=\lim _{n \rightarrow \infty} \int_{\tilde{\Gamma}} T_{\psi, I} d \omega_{n}=\int_{\tilde{\Gamma}} T_{\psi, I} d \omega=\int_{I} \int_{\tilde{\Gamma}(t)} \psi(\gamma(t)) d \omega d t
$$

This proves that the second limit in (3.17) holds for $\mathscr{L}^{1}$-a.e. $t \in(0,1)$. In order to prove that the limit is valid for every $t \in(0,1)$, we observe that $\omega$ is concentrated on curves with endpoints in $\partial B_{R}$ and for every $t \in(0,1)$ it holds

$$
\omega\left(\left\{\left(\gamma, t_{\gamma}^{-}, t_{\gamma}^{+}\right) \in \tilde{\Gamma}: t \in\left(t_{\gamma}^{-}, t_{\gamma}^{+}\right) \text {and } \gamma_{a}(t-) \neq \gamma_{a}(t+)\right\}\right)=0
$$

Indeed, assume by contradiction that there exist $\bar{t} \in(0,1)$ and $\varepsilon>0$ such that

$$
\begin{equation*}
\omega\left(\left\{\left(\gamma, t_{\gamma}^{-}, t_{\gamma}^{+}\right) \in \tilde{\Gamma}: \bar{t} \in\left(t_{\gamma}^{-}, t_{\gamma}^{+}\right) \text {and }\left|\gamma_{a}(\bar{t}-)-\gamma_{a}(\bar{t}+)\right|>\varepsilon\right\}\right)>\varepsilon \tag{3.18}
\end{equation*}
$$

Inequality (3.18) implies that for every $t_{1}, t_{2} \in(0,1)$ such that $t_{1}<\bar{t}<t_{2}$ it holds

$$
\begin{equation*}
e_{V}\left(t_{1}, t_{2}\right):=\int_{\tilde{\Gamma}} \operatorname{Tot}^{\left.\operatorname{Var}_{\left(t_{y}^{-}\right.}, t_{\gamma}^{+}\right) \cap\left(t_{1}, t_{2}\right)} \gamma_{a} d \omega(\gamma) \geq \varepsilon^{2} \tag{3.19}
\end{equation*}
$$

On the other hand, by localizing in $\left(t_{1}, t_{2}\right)$ the same argument as in the proof of Lemma 3.7 to obtain (3.11), we have that

$$
\begin{equation*}
e_{\nu}\left(n, t_{1}, t_{2}\right):=\int_{\tilde{\Gamma}} \operatorname{Tot}^{\operatorname{Var}_{\left(t_{y}^{-}, t_{y}^{+}\right) n\left(t_{1}, t_{2}\right)} y_{a} d \omega_{n}(\gamma) \leq\left(t_{2}-t_{1}\right) v\left(B_{R}\right)+o(1) \quad \text { as } n \rightarrow \infty . . . . . .} \tag{3.20}
\end{equation*}
$$

By choosing

$$
t_{2}-t_{1}<\frac{\varepsilon^{2}}{1+v\left(B_{R}\right)}
$$

the two conditions in (3.19) and (3.20) contradict each other.
In particular, $t \mapsto\left(e_{t}\right)_{\sharp} \omega$ is continuous in the sense of distributions on $B_{R} \times[0, M]$, and therefore the second limit in (3.17) holds for every $t \in(0,1)$.
Condition (2). The function $g: \tilde{\Gamma} \rightarrow \mathbb{R}$ defined by

$$
g\left(y, t_{y}^{-}, t_{y}^{+}\right):=\sup _{t \in\left(t_{y}^{-}, t_{y}^{+}\right)}\left|\gamma_{x}(t)-\gamma_{x}\left(t_{\gamma}^{-}\right)-\int_{t_{\gamma}^{-}}^{t} i e^{i y_{a}(s)} d s\right|
$$

is lower semicontinuous. Therefore,

$$
\int_{\tilde{\Gamma}} g(\gamma) d \omega \leq \lim _{n \rightarrow \infty} \int_{\tilde{\Gamma}} g(\gamma) d \omega_{n}
$$

which is equal to 0 by (3.10).
Condition (3). This follows similarly from (3.11). In particular, $\omega$ is concentrated on $\Gamma$ and this concludes the proof.

### 3.6 Representation of the defect measure and good curves selection

In the following proposition, we show that the kinetic measure $U_{\phi}$ can be decomposed along the characteristic trajectories detected by the Lagrangian representation $\omega_{h}$.

Proposition 3.11. Let $\omega_{h}$ be a Lagrangian representation of the hypograph of $\phi$ on $B_{R}$ obtained as limit point of $\omega_{n}$ as in the previous section. Then

$$
\mathscr{L}^{1} \times U_{\phi}=\int_{\Gamma} \mu_{\gamma} d \omega_{h}(\gamma) \quad \text { and } \quad \mathscr{L}^{1} \times\left|U_{\phi}\right|=\int_{\Gamma}\left|\mu_{\gamma}\right| d \omega_{h}(\gamma) .
$$

Proof. Let

$$
\bar{\psi}(t, x, a)=\varphi(t) \psi(x, a) \in C_{c}^{\infty}\left((0,1) \times B_{R} \times[0, M]\right)
$$

Then

$$
\begin{align*}
-\int_{(0,1) \times X} \varphi \partial_{a} \psi d U_{\phi} d t & =\int_{(0,1) \times H_{\phi}} i e^{i a} \cdot \nabla_{x} \psi \varphi d x d a d t \\
& =\int_{0}^{1} \int_{\Gamma(t)} i e^{i y_{a}(t)} \cdot \nabla_{x} \psi(\gamma(t)) \varphi(t) d \omega_{h}(\gamma) d t \\
& =\int_{0}^{1} \int_{\Gamma(t)} \dot{\gamma}_{x}(t) \cdot \nabla_{x} \psi(\gamma(t)) d \omega_{h}(\gamma) \varphi(t) d t \\
& =\int_{\Gamma}^{t_{\nu}^{+}} \int_{t_{\gamma}^{-}} \dot{\gamma}_{x}(t) \cdot \nabla_{x} \psi(\gamma(t)) \varphi(t) d t d \omega_{h}(\gamma) \tag{3.21}
\end{align*}
$$

For every $\gamma \in \Gamma$, we consider the map $\psi_{\gamma}:=\psi \circ \gamma:\left(t_{\gamma}^{-}, t_{\gamma}^{+}\right) \rightarrow B_{R} \times[0, M]$. Since $\omega_{h}$-a.e. $\gamma \in \Gamma$ has bounded variation on its domain, also $\psi_{y} \in \mathrm{BV}\left(\left(t_{\gamma}^{-}, t_{\gamma}^{+}\right) ; \mathbb{R}\right)$, and we have the following chain rule:

$$
\begin{align*}
D_{t} \psi_{y} & =\nabla \psi(\gamma(t)) \cdot \tilde{D}_{t} y+\sum_{t_{j} \in J_{y}}\left(\psi\left(\gamma\left(t_{j}+\right)\right)-\psi\left(\gamma\left(t_{j}-\right)\right)\right) \delta_{t_{j}} \\
& =\nabla_{x} \psi(\gamma(t)) \cdot \tilde{D}_{t} \gamma_{x}+\partial_{a} \psi(\gamma(t)) \tilde{D}_{t} \gamma_{a}+\sum_{t_{j} \in J_{y}}\left(\psi\left(\gamma\left(t_{j}+\right)\right)-\psi\left(\gamma\left(t_{j}-\right)\right)\right) \delta_{t_{j}} \tag{3.22}
\end{align*}
$$

Since for $\omega$-a.e. $\gamma$ it holds $D_{t} \gamma_{x}=\dot{\gamma}_{x}(t) \mathscr{L}^{1}$, plugging (3.22) into (3.21), we obtain

$$
\begin{aligned}
-\int_{(0,1) \times X} \varphi \partial_{a} \psi d U_{\phi} d t & =\int_{\Gamma}\left(\int_{\left(t_{y}^{+}, t_{y}^{+}\right)} \varphi\left(D_{t} \psi_{y}-\partial_{a} \psi(\gamma(t)) \tilde{D}_{t} \gamma_{a}-\sum_{t_{j} \in J_{y}}\left(\psi_{y}\left(t_{j}+\right)-\psi_{y}\left(t_{j}-\right)\right) \delta_{t_{j}}\right)\right) d \omega_{h} \\
& =\int_{\Gamma}\left(\int_{\left(t_{y}^{-}, t_{y}^{+}\right)} \varphi\left(D_{t} \psi_{y}-\partial_{a} \psi(\gamma(t)) \tilde{D}_{t} \gamma_{a}\right)-\sum_{t_{j} \in J_{y}} \varphi\left(t_{j}\right)\left(\psi_{y}\left(t_{j}+\right)-\psi_{y}\left(t_{j}-\right)\right)\right) d \omega_{h}
\end{aligned}
$$

We observe that, by construction, if $t_{\gamma}^{-}>0$, then $\gamma\left(t_{\gamma}^{-}\right) \in \partial B_{R} \times[0, M]$, and therefore $\psi\left(\gamma\left(t_{\gamma}^{-}\right)\right)=0$. Similarly, if $t_{\gamma}^{+}<1$, then $\psi\left(\gamma\left(t_{\gamma}^{+}\right)\right)=0$. Therefore,

$$
\begin{aligned}
\int_{\Gamma} \int_{\left(t_{\gamma}^{-}, t_{\gamma}^{+}\right)} \varphi(t) D_{t} \psi_{y} d \omega_{h}(\gamma) & =-\int_{\Gamma} \int_{\left(t_{\gamma}^{+}, t_{y}^{+}\right)} \varphi^{\prime}(t) \psi(\gamma(t)) d t d \omega_{h}(\gamma) \\
& =\int_{0}^{1} \int_{H_{\phi}} \varphi^{\prime}(t) \psi(x, a) d x d a d t=0
\end{aligned}
$$

Since for $\omega_{h}$-a.e. $\gamma$ and every $t_{j} \in J_{y}$ it holds $\gamma_{x}\left(t_{j}+\right)=\gamma_{x}\left(t_{j}-\right)$, it follows from the definition of $\mu_{y}$ that

$$
\begin{aligned}
-\int_{(0,1) \times X} \varphi \partial_{a} \psi d U_{\phi} d t & =\int_{\Gamma}\left(\int_{\left(t_{y}, t_{\nu}^{+}\right)}\left(-\varphi(t) \partial_{a} \psi(\gamma(t)) \tilde{D}_{t} \gamma_{a}\right)-\sum_{t_{j} \in J_{y}} \varphi\left(t_{j}\right)\left(\psi\left(\gamma\left(t_{j}+\right)\right)-\psi\left(\gamma\left(t_{j}-\right)\right)\right)\right) d \omega_{h} \\
& =-\int_{\Gamma} \int_{(0,1) \times X} \varphi \partial_{a} \psi d \mu_{y} d \omega_{h}(\gamma)
\end{aligned}
$$

This proves the first equality in the statement when tested with functions of the form $\varphi \partial_{a} \psi$ for two test functions $\varphi, \psi$. Since both $U_{\phi}$ and $\int \mu_{\gamma} d \omega_{h}$ are supported on $[0,1] \times B_{R} \times[0, M]$, the equality holds true for every test function.

The inequality

$$
\mathscr{L}^{1} \times\left|U_{\phi}\right| \leq \int_{\Gamma}\left|\mu_{y}\right| d \omega_{h}
$$

follows immediately from the already proved first equality in the statement. In order to prove the opposite inequality, it is enough to prove the global inequality

$$
\left(\mathscr{L}^{1} \times\left|U_{\phi}\right|\right)\left((0,1) \times B_{R} \times[0, M]\right) \geq \int_{\Gamma}\left|\mu_{\gamma}\right|\left((0,1) \times B_{R} \times[0, M]\right) d \omega_{h} .
$$

We observe that $\left|\mu_{\gamma}\right|\left((0,1) \times B_{R} \times[0, M]\right)=\operatorname{TotVar}_{\left(t_{\nu}^{-}, t_{\nu}^{+}\right)} \gamma_{a}$ and that the map

$$
\left(\gamma, t_{\gamma}^{-}, t_{\gamma}^{+}\right) \mapsto \operatorname{Tot}^{\left.\operatorname{Var}_{\left(t_{\gamma}^{-}\right.}, t_{y}^{+}\right)} \gamma_{a}
$$

is lower semicontinuous on $\tilde{\Gamma}$. Therefore, it follows from (3.11) that

$$
\begin{aligned}
& \int_{\Gamma}\left|\mu_{y}\right|\left(B_{R} \times[0, M]\right) d \omega_{h}=\int_{\Gamma} \operatorname{Tot}^{\operatorname{Var}} \\
&\left(t_{\gamma}^{-}, t_{\gamma}^{+}\right) \\
& \leq y_{n \rightarrow \infty} d \omega_{h} \\
& \leq\left(\mathscr{L}^{1} \times\left|U_{\phi}\right|\right)\left((0,1) \times B_{R} \times[0, M]\right)
\end{aligned}
$$

With the result above, the proof of the part of Theorem 3.2 concerning the hypograph of $\phi$ is complete; the statement for the epigraph of $\phi$ can be proven in the same way.

The following lemma is an application of the Tonelli theorem and it is already proven in [19], to which we refer for the details.

Lemma 3.12. For $\omega_{h}$-a.e. $\gamma \in \Gamma$ and for $\mathscr{L}^{1}$-a.e. $t \in\left(t_{\gamma}^{-}, t_{\gamma}^{+}\right)$, the following assertions hold:
(i) $\gamma_{x}(t)$ is a Lebesgue point of $\phi$.
(ii) $\gamma_{a}(t)<\phi\left(\gamma_{x}(t)\right)$.

We denote by $\Gamma_{h}$ the set of curves $\gamma \in \Gamma$ such that the two properties above hold. Similarly, for $\omega_{e}-$ a.e. $\gamma \in \Gamma$ and for $\mathscr{L}^{1}$-a.e. $t \in\left(t_{y}^{-}, t_{\gamma}^{+}\right)$, the following assertions hold:
(i) $\gamma_{x}(t)$ is a Lebesgue point of $\phi$.
(ii) $\gamma_{a}(t)>\phi\left(\gamma_{x}(t)\right)$

We denote the set of these curves by $\Gamma_{e}$.

## 4 Rectifiability of the measure $v$

In this section, we prove that the measure $v:=\left(p_{x}\right)_{\sharp}\left|U_{\phi}\right|$ is concentrated on a 1-rectifiable set. The rectifiability of $v$ is equivalent to the rectifiability of both the measures $\left(p_{x}\right)_{\sharp} U_{\phi}^{-}$and $\left(p_{x}\right)_{\sharp} U_{\phi}^{+}$. Since these two cases are analogous, we only provide the proof of the rectifiability of $\left(p_{x}\right)_{\sharp} U_{\phi}^{-}$.

### 4.1 Pairing between $\omega_{h}$ and $\omega_{e}$ and its decomposition

In the following lemma, we introduce a pairing between the two representations $\omega_{h} \otimes \mu_{\gamma}^{-}$and $\omega_{e} \otimes \mu_{\gamma}^{+}$of the negative part of the defect measure $\mathscr{L}^{1} \times U_{\phi}^{-}$. We will denote by $X$ the set $B_{R} \times[0, M]$.
Lemma 4.1. Denote by $p_{1}, p_{2}:(\Gamma \times[0,1] \times X)^{2} \rightarrow \Gamma \times[0,1] \times X$ the standard projections. Then there exists a plan $\pi^{-} \in \mathcal{M}\left((\Gamma \times[0,1] \times X)^{2}\right)$ with marginals

$$
\left\{\begin{array}{l}
\left(p_{1}\right)_{\sharp} \pi^{-}=\omega_{h} \otimes \mu_{\gamma}^{-},  \tag{4.1}\\
\left(p_{2}\right)_{\sharp} \pi^{-}=\omega_{e} \otimes \mu_{\gamma}^{+},
\end{array}\right.
$$

## concentrated on the set

$$
\begin{gathered}
\mathcal{G}:=\left\{\left(\left(\gamma, t_{\gamma}^{-}, t_{y}^{+}, t, x, a\right),\left(y^{\prime}, t_{\gamma}^{-\prime}, t_{y}^{+\prime}, t^{\prime}, x^{\prime}, a^{\prime}\right)\right) \in(\Gamma \times X)^{2}: t \in\left(t_{y}^{-}, t_{\gamma}^{+}\right), t^{\prime} \in\left(t_{y}^{-\prime}, t_{y}^{+\prime}\right), t=t^{\prime},\right. \\
\left.\gamma_{x}(t)=x=x^{\prime}=y_{x}^{\prime}\left(t^{\prime}\right), a=a^{\prime}, a \in\left[\gamma_{a}(t+), \gamma_{a}(t-)\right] \cap\left[\gamma_{a}^{\prime}\left(t^{\prime}-\right), y_{a}^{\prime}\left(t^{\prime}+\right)\right]\right\} .
\end{gathered}
$$

Proof. First, we observe that, by definition, $\omega_{h} \otimes \mu_{\gamma}^{-}$is concentrated on the set

$$
\mathcal{G}_{h}^{-}:=\left\{\left(\gamma, t_{\gamma}^{-}, t_{\gamma}^{+}, t, x, a\right) \in \Gamma \times[0,1] \times X: t \in\left(t_{\gamma}^{-}, t_{\gamma}^{+}\right), \gamma_{x}(t)=x, a \in\left[\gamma_{a}(t+), \gamma_{a}(t-)\right]\right\}
$$

and $\omega_{e} \otimes \mu_{\gamma}^{+}$is concentrated on the set

$$
\mathcal{G}_{e}^{+}:=\left\{\left(\gamma, t_{\gamma}^{-}, t_{y}^{+}, t, x, a\right) \in \Gamma \times[0,1] \times X: t \in\left(t_{y}^{-}, t_{y}^{+}\right), \gamma_{x}(t)=x, a \in\left[\gamma_{a}(t-), \gamma_{a}(t+)\right]\right\} .
$$

By denoting by $p_{2,3}: \Gamma \times[0,1] \times X \rightarrow[0,1] \times X$ the standard projection, it follows from (3.4) that

$$
\left(p_{2,3}\right)_{\sharp}\left(\omega_{h} \otimes \mu_{\gamma}^{-}\right)=\mathscr{L}^{1} \times U_{\phi}^{-}=\left(p_{2,3}\right)_{\sharp}\left(\omega_{e} \otimes \mu_{\gamma}^{+}\right) .
$$

By the disintegration theorem (see, for example, [2]), there exist two measurable families of probability measures

$$
\left(\mu_{t, x, a}^{-, h}\right)_{(t, x, a) \in X},\left(\mu_{t, x, a}^{+, e}\right)_{(t, x, a) \in X} \in \mathcal{P}(\Gamma \times[0,1] \times X)
$$

such that

$$
\begin{equation*}
\omega_{h} \otimes \mu_{\gamma}^{-}=\int_{[0,1] \times X} \mu_{t, x, a}^{-, h} d \mathscr{L}^{1} \times U_{\phi}^{-} \quad \text { and } \quad \omega_{e} \otimes \mu_{\gamma}^{+}=\int_{[0,1] \times X} \mu_{t, x, a}^{+, e} d \mathscr{L}^{1} \times U_{\phi}^{-} \tag{4.2}
\end{equation*}
$$

and for $\mathscr{L}^{1} \times U_{\phi}^{-}$-a.e. $(t, x, a)$ the measures $\mu_{t, x, a}^{-, h}$ and $\mu_{t, x, a}^{+, e}$ are concentrated on the set

$$
p_{2,3}^{-1}(\{t, x, a\})=\left\{\left(y, t_{\gamma}^{-}, t_{y}^{+}, t^{\prime}, x^{\prime}, a^{\prime}\right) \in \Gamma \times[0,1] \times X: t^{\prime}=t, x^{\prime}=x, a^{\prime}=a\right\}
$$

Moreover, since $\omega_{h} \otimes \mu_{\gamma}^{-}$is concentrated on the set $\mathcal{G}_{h}^{-}$and $\omega_{e} \otimes \mu_{\gamma}^{+}$is concentrated on the set $\mathcal{G}_{e}^{+}$, we have that for $\mathscr{L}^{1} \times U_{\phi}^{-}$-a.e. $(t, x, a)$ the measure $\mu_{t, x, a}^{-, h}$ is concentrated on $p_{2,3}^{-1}(\{t, x, a\}) \cap \mathcal{G}_{h}^{-}$and $\mu_{t, x, a}^{+, e}$ is concentrated on $p_{2,3}^{-1}(\{t, x, a\}) \cap \mathcal{G}_{e}^{+}$. Eventually, we set

$$
\pi^{-}:=\int_{[0,1] \times X}\left(\mu_{t, x, a}^{-, h} \otimes \mu_{t, x, a}^{+, e}\right) d\left(\mathscr{L}^{1} \times U_{\phi}^{-}\right) .
$$

From (4.2), directly (4.1) follows. By the above discussion, for $\mathscr{L}^{1} \times U_{\phi}^{-}$-a.e. $(t, x, a) \in[0,1] \times X$ the measure $\mu_{t, x, v}^{-, h} \otimes \mu_{t, x, v}^{+, e}$ is concentrated on $\left(p_{2,3}^{-1}(\{t, x, a\}) \cap \mathcal{G}_{h}^{-}\right) \times\left(p_{2,3}^{-1}(\{t, x, a\}) \cap \mathcal{G}_{e}^{+}\right)$. Therefore, $\pi^{-}$is concentrated on

$$
\bigcup_{(t, x, a) \in[0,1] \times X}\left(p_{2,3}^{-1}(\{t, x, a\}) \cap \mathcal{G}_{h}^{-}\right) \times\left(p_{2,3}^{-1}(\{t, x, a\}) \cap \mathcal{G}_{e}^{+}\right)=\mathcal{G}
$$

and this concludes the proof.
Now, we split the set $\mathcal{G}$ introduced in Lemma 4.1 in finitely many components. First, we set

$$
\begin{aligned}
& \mathcal{G}_{h, \text { jump }}^{-}:=\left\{\left(\gamma, t_{\gamma}^{-}, t_{\gamma}^{+}, t, x, a\right) \in \mathcal{G}_{h}^{-}: \gamma_{a}(t+)<\gamma_{a}(t-)\right\}, \\
& \mathcal{G}_{e, \text { jump }}^{+}:=\left\{\left(\gamma, t_{\gamma}^{-}, t_{\gamma}^{+}, t, x, a\right) \in \mathcal{G}_{e}^{+}: \gamma_{a}(t-)<\gamma_{a}(t+)\right\} .
\end{aligned}
$$

Moreover, we consider the following covering with overlaps of $[0, M]$. Let $L=\left\lfloor\frac{2 M}{\pi}\right\rfloor$ and for every $l=0, \ldots, L$ set

$$
I_{l}=\left(l \frac{\pi}{2}-\frac{\pi}{8},(l+1) \frac{\pi}{2}+\frac{\pi}{8}\right)
$$

and

$$
\begin{aligned}
& \mathcal{G}_{h, l}^{-}:=\left\{\left(\gamma, t_{\gamma}^{-}, t_{\gamma}^{+}, t, x, a\right) \in \mathcal{G}_{h}^{-}: \gamma_{a}(t+), \gamma_{a}(t-) \in I_{l}\right\}, \\
& \mathcal{G}_{e, l}^{+}:=\left\{\left(\gamma, t_{\gamma}^{-}, t_{\gamma}^{+}, t, x, a\right) \in \mathcal{G}_{e}^{+}: \gamma_{a}(t-), \gamma_{a}(t+) \in I_{l}\right\} .
\end{aligned}
$$

Then we define

$$
\pi_{l}^{-}:=\pi_{\llcorner }^{-}\left(\mathcal{G}_{h, l}^{-} \times \mathcal{G}_{e, l}^{+}\right), \quad \pi_{\mathrm{jump}}^{-}=\pi_{\llcorner }^{-}\left(\left(\mathcal{G}_{h, \mathrm{jump}}^{-} \times \mathcal{G}_{e}^{+}\right) \cup\left(\mathcal{G}_{h}^{-} \times \mathcal{G}_{e, \mathrm{jump}}^{+}\right)\right)
$$

We prove separately that $v_{\text {jump }}^{-}:=\left(p_{x}^{1}\right)_{\sharp} \pi_{\text {jump }}^{-}$is 1-rectifiable and that $v_{l}^{-}:=\left(p_{x}^{1}\right)_{\sharp} \pi_{l}^{-}$is rectifiable for every $l=0, \ldots, L$.


Figure 1: The blue curve $\gamma$ represents an element of $G_{\mathrm{cr}}\left(\bar{\gamma}, \tilde{t}_{\bar{\gamma}}^{-}, \tilde{t}_{\bar{\gamma}}^{+}\right)$: in this particular case, $l=4 k$ for some $k \in \mathbb{N}$.

### 4.2 Rectifiability of $v_{l}^{-}$

The proof of the rectifiability of $v_{l}^{-}$follows the strategy used in [19]. In particular, the first step is to identify a countable family of Lipschitz curves where we will prove that $v_{l}^{-}$is concentrated.

### 4.2.1 Shock curves

For shortness, we set

$$
e_{l}:=i e^{i\left(l \frac{\pi}{2}+\frac{\pi}{4}\right)} \quad \text { and } \quad e_{l}^{\perp}:=i e_{l} .
$$

The following proposition establishes the intuitive fact that a curve of the epigraph cannot cross from below a curve of the hypograph. Since the same proposition and the following corollary were proven in [19] in the case of the Burgers equation, we only sketch the arguments here.

Proposition 4.2. Let $\left(\bar{\gamma}, t_{\bar{\gamma}}^{-}, t_{\bar{\gamma}}^{+}\right) \in \Gamma_{h}$ and let $\left(\tilde{t}_{\bar{\gamma}}^{-}, \tilde{t}_{\bar{\gamma}}^{+}\right) \subset\left(t_{\bar{\gamma}}^{-}, t_{\bar{\gamma}}^{+}\right)$be such that

$$
\bar{\gamma}_{a}\left(\left(\tilde{t}_{\bar{\gamma}}^{-}, \tilde{t}_{\bar{\gamma}}^{+}\right)\right) \subset I_{l}
$$

We denote by $G_{\mathrm{cr}}\left(\bar{\gamma}, \tilde{t}_{\bar{\gamma}}^{-}, \tilde{t}_{\bar{\gamma}}^{+}\right)$the set of curves

$$
\left(y, t_{\gamma}^{-}, t_{\gamma}^{+}\right) \in \Gamma_{e}
$$

for which there exist $\bar{t}_{1}, \bar{t}_{2} \in\left(\tilde{t}_{\bar{\gamma}}^{-}, \tilde{t}_{\bar{\gamma}}^{+}\right)$and $t_{1}, t_{2} \in\left(t_{\gamma}^{-}, t_{\gamma}^{+}\right)$such that the following conditions are satisfied (see Figure 1):
(i) $t_{1}<t_{2}$ and $\bar{t}_{1}<\bar{t}_{2}$.
(ii) $\gamma_{a}\left(\left(t_{1}, t_{2}\right)\right) \subset I_{l}$.
(iii) $\gamma_{x}\left(t_{1}\right) \cdot e_{l}=\bar{\gamma}_{x}\left(\bar{t}_{1}\right) \cdot e_{l}$ and $\gamma_{x}\left(t_{1}\right) \cdot e_{l}^{\perp}>\bar{\gamma}_{x}\left(\bar{t}_{1}\right) \cdot e_{l}^{\perp}$.
(iv) $\gamma_{x}\left(t_{2}\right) \cdot e_{l}=\bar{\gamma}_{x}\left(\bar{t}_{2}\right) \cdot e_{l}$ and $\gamma_{x}\left(t_{2}\right) \cdot e_{l}^{\perp}<\bar{\gamma}_{x}\left(\bar{t}_{2}\right) \cdot e_{l}^{\perp}$.

Then

$$
\omega_{e}\left(G_{\mathrm{cr}}\left(\bar{\gamma}, \tilde{t}_{\bar{\gamma}}^{-}, \tilde{t}_{\bar{\gamma}}^{+}\right)\right)=0
$$

Proof. Let

$$
s^{-}:=\bar{\gamma}_{x}\left(\tilde{t}_{\bar{\gamma}}^{-}\right) \cdot e_{l} \quad \text { and } \quad s^{+}:=\bar{\gamma}_{x}\left(\tilde{t}_{\bar{\gamma}}^{+}\right) \cdot e_{l} .
$$

Since $\bar{\gamma}_{a}\left(\left(\tilde{t}_{\bar{\gamma}}^{-}, \tilde{t}_{\bar{\gamma}}^{+}\right)\right) \subset I_{l}$ and $\dot{\bar{\gamma}}_{x}(t)=i e^{i \bar{\gamma}_{a}(t)}$ for $\mathscr{L}^{1}$-a.e. $t \in\left(\tilde{t}_{\bar{\gamma}}^{-}, \tilde{t}_{\bar{\gamma}}^{+}\right)$, the map

$$
\begin{aligned}
h_{\bar{\gamma}}:\left(\tilde{t}_{\bar{\gamma}}^{-}, \tilde{t}_{\bar{\gamma}}^{+}\right) & \rightarrow\left(s^{-}, s^{+}\right), \\
t & \mapsto \bar{\gamma}_{x}(t) \cdot e_{l}
\end{aligned}
$$

is bi-Lipschitz. For every $s \in\left(s^{-}, s^{+}\right)$, we set $g_{\bar{\gamma}}(s)=y_{x}\left(h_{\bar{\gamma}}^{-1}(t)\right) \cdot e_{l}^{\perp}$. Let $\delta>0$ and let $\psi_{\delta}: \mathbb{R} \rightarrow \mathbb{R}$ be the Lipschitz approximation of the Heaviside function defined by $\psi_{\delta}(v)=0 \vee(v / \delta \wedge 1)$. Let us consider a measurable selection of $t_{1}, t_{2}$ in $G_{\mathrm{cr}}\left(\bar{\gamma}, \tilde{t}_{\bar{\gamma}}^{-}, \tilde{t}_{\bar{\gamma}}^{+}\right)$and let us set

$$
G_{\mathrm{cr}}\left(\bar{\gamma}, \tilde{t}_{\bar{\gamma}}^{-}, \tilde{t}_{\bar{\gamma}}^{+}, \delta\right):=\left\{\left(\gamma, t_{\gamma}^{-}, t_{\gamma}^{+}\right) \in G_{\mathrm{cr}}\left(\bar{\gamma}, \tilde{t}_{\bar{\gamma}}^{-}, \tilde{t}_{\bar{\gamma}}^{+}\right): \gamma_{x}\left(t_{1, \gamma}\right) \cdot e_{l}^{\perp}-g_{\bar{\gamma}}\left(h_{\bar{\gamma}}\left(\gamma_{x}(t) \cdot e_{l}\right)\right)>\delta\right\} .
$$

For every $t \in(0,1)$ and $\gamma \in G_{\mathrm{cr}}\left(\bar{\gamma}, \tilde{t}_{\bar{\gamma}}^{-}, \tilde{t}_{\bar{\gamma}}^{+}, \delta\right)$, set

$$
f(\gamma, t):= \begin{cases}0 & \text { if } t<t_{1, \gamma} \\ 1-\psi_{\delta}\left(\gamma_{x}(t) \cdot e_{l}^{\perp}-g_{\bar{\gamma}}\left(y_{x}(t) \cdot e_{l}\right)\right) & \text { if } t \in\left(t_{1, \gamma}, t_{2, \gamma}\right) \\ 1 & \text { if } t>t_{2, \gamma}\end{cases}
$$

Finally, we consider the functional

$$
\Psi_{\delta}(t):=\int_{\Gamma_{\mathrm{cr}}\left(\bar{y}, \tilde{t}_{\bar{\gamma}}^{-}, \tilde{t}_{\bar{y}}^{+}, \delta\right)} f(\gamma, t) d \omega_{e}(\gamma)
$$

A straightforward computation shows that

$$
\begin{equation*}
\Psi_{\delta}^{\prime}(t) \leq \frac{C}{\delta} \int_{G(\delta, t)}\left[\bar{\gamma}_{a}\left(h_{\bar{\gamma}}^{-1}\left(\gamma_{x}(t) \cdot e_{l}\right)\right)-\gamma_{a}(t)\right]^{+} d \omega_{e}(\gamma) \tag{4.3}
\end{equation*}
$$

where

$$
\begin{array}{r}
G(\delta, t)=\left\{\left(\gamma, t_{\gamma}^{-}, t_{\gamma}^{+}\right) \in G_{\mathrm{cr}}\left(\bar{\gamma}, \tilde{t}_{\bar{\gamma}}^{-}, \tilde{t}_{\bar{\gamma}}^{+}, \delta\right): t \in\left(t_{1, \gamma}, t_{2, \gamma}\right)\right. \text { and } \\
\left.\gamma_{x}(t) \cdot e_{l}^{\perp} \in\left(g_{\bar{\gamma}}\left(\gamma_{x}(t) \cdot e_{l}\right), g_{\bar{\gamma}}\left(\gamma_{x}(t) \cdot e_{l}\right)+\delta\right)\right\} .
\end{array}
$$

Let us set

$$
S_{\delta}:=\left\{x \in B_{R}: x \cdot e_{l} \in\left(g_{\bar{y}}\left(x \cdot e_{l}\right), g_{\bar{y}}\left(x \cdot e_{l}\right)+\delta\right)\right\}
$$

Since

$$
\left(e_{t}\right)_{\sharp} \omega_{e\llcorner } G(\delta, t) \leq \mathcal{L}_{\llcorner }^{3}\left(E_{\phi} \cap\left(S_{\delta} \times[0, M]\right)\right)
$$

and for $\mathscr{L}^{1}$-a.e. $t \in\left(\tilde{t}_{\bar{\gamma}}^{-}, \tilde{t}_{\bar{\gamma}}^{+}\right)$the point $\bar{\gamma}_{x}(t)$ is a Lebesgue point of $\phi$ with value larger than $\bar{\gamma}_{a}(t)$, we obtain from (4.3) that $\Psi_{\delta}^{\prime}(t) \leq o(1)$ as $\delta \rightarrow 0$. By the definition of the functional $\Psi_{\delta}$ it holds

$$
\omega_{e}\left(G_{\mathrm{cr}}\left(\bar{\gamma}, \tilde{t}_{\bar{\gamma}}^{-}, \tilde{t}_{\bar{\gamma}}^{+}\right)\right) \leq \liminf _{\delta \rightarrow 0} \Psi_{\delta}(1)=0
$$

Corollary 4.3. Let $\bar{x} \in B_{R}$ and denote by $\Gamma_{l}^{-}(\bar{x})$ the set of curves $\left(\gamma, t_{\gamma}^{-}, t_{\gamma}^{+}\right) \in \Gamma_{h}$ for which there exists $t_{1} \in\left(t_{\gamma}^{-}, t_{\gamma}^{+}\right)$ such that

$$
\gamma_{x}\left(t_{1}\right) \cdot e_{l}=\bar{x} \cdot e_{l} \quad \text { and } \quad \gamma_{x}\left(t_{1}\right) \cdot e_{l}^{\perp}<\bar{x} \cdot e_{l}^{\perp}
$$

Similarly, let $\Gamma_{l}^{+}(\bar{x})$ be the set of curves $\left(\gamma, t_{\gamma}^{-}, t_{\gamma}^{+}\right) \in \Gamma_{e}$ for which there exists $t_{1}^{\prime} \in\left(t_{\gamma}^{-}, t_{\gamma}^{+}\right)$such that

$$
\gamma_{x}\left(t_{1}^{\prime}\right) \cdot e_{l}=\bar{x} \cdot e_{l} \quad \text { and } \quad \gamma_{x}\left(t_{1}^{\prime}\right) \cdot e_{l}^{\perp}>\bar{x} \cdot e_{l}^{\perp}
$$

Then there exists a Lipschitz function $f_{\bar{\chi}, l}:\left[\bar{x} \cdot e_{l},+\infty\right) \rightarrow \mathbb{R}$ such that

$$
\left\{\begin{array}{l}
\omega_{h}\left(\left\{\left(\gamma, t_{\gamma}^{-}, t_{\gamma}^{+}\right) \in \Gamma_{l}^{-}(\bar{x}): \exists t_{2} \in\left(t_{1}, t_{\gamma}^{+}\right) \text {s.t. } \gamma_{a}\left(\left(t_{1}, t_{2}\right)\right) \subset I_{l} \text { and } \gamma_{x}\left(t_{2}\right) \cdot e_{l}^{\perp}>f_{\bar{x}, l}\left(\gamma_{x}\left(t_{2}\right) \cdot e_{l}\right)\right\}\right)=0  \tag{4.4}\\
\omega_{e}\left(\left\{\left(\gamma, t_{\gamma}^{-}, t_{\gamma}^{+}\right) \in \Gamma_{l}^{+}(\bar{x}): \exists t_{2}^{\prime} \in\left(t_{1}^{\prime}, t_{\gamma}^{+}\right) \text {s.t. } \gamma_{a}\left(\left(t_{1}^{\prime}, t_{2}^{\prime}\right)\right) \subset I_{l} \text { and } \gamma_{x}\left(t_{2}^{\prime}\right) \cdot e_{l}^{\perp}<f_{\bar{x}, l}\left(\gamma_{x}\left(t_{2}^{\prime}\right) \cdot e_{l}\right)\right\}\right)=0
\end{array}\right.
$$

where $t_{1}, t_{1}^{\prime}$ are as above.
Proof. Let $I \subset\left[\bar{x} \cdot e_{l},+\infty\right)$ be the set of values $y$ for which there exist $\gamma \in \Gamma_{l}^{-}(\bar{x})$ and $t \in\left(t_{1}, t_{2}\right)$ such that $\gamma_{x}(t) \cdot e_{l}=y$, where $t_{1}<t_{2}$ are such that $\gamma_{a}\left(\left(t_{1}, t_{2}\right)\right) \subset I_{l}$. Let $\tilde{f}_{\bar{x}, l}$ be defined on $I$ by

$$
\begin{equation*}
\tilde{f}_{\bar{x}, l}(s):=\sup \left\{y_{x}(t) \cdot e_{l}^{\perp}: \gamma \in \Gamma_{l}^{-}(\bar{x}), t \in\left(t_{1}, t_{2}\right), \gamma_{x}(t) \cdot e_{l}=s\right\}=\sup g_{\gamma}(s) \tag{4.5}
\end{equation*}
$$

where we used the same notation as in the proof of Proposition 4.2: we set $g_{\gamma}(s):=\gamma_{x}(t) \cdot e_{l}^{\perp}$, where $t$ is the unique value in $\left(t_{1}, t_{2}\right)$ for which $\gamma_{x}(t) \cdot e_{l}=s$. Since $\gamma_{a}\left(\left(t_{1}, t_{2}\right)\right) \subset I_{l}$, it is straightforward to check that $g_{y}$ is Lipschitz with Lipschitz constant bounded by $\tan \left(\frac{3 \pi}{8}\right)$. The function $f_{\bar{x}, l}$ is then defined as the smallest biggest $C$-Lipschitz function such that $f_{\bar{x}, l} \geq \tilde{f}_{\bar{x}, l}$ on $I$ and $f_{\bar{x}, l}\left(\bar{x} \cdot e_{l}\right)=\bar{x} \cdot e_{l}^{\perp}$, where $C>\tan (3 \pi / 8)$. The first equation in (4.4) follows from the fact that $f_{\bar{x}, l} \geq \tilde{f}_{\bar{x}, l}$ on $I$. Now, we prove the second equation in (4.4): given $\left(\gamma, t_{\gamma}^{-}, t_{\gamma}^{+}\right) \in \Gamma_{l}^{-}(\bar{x})$ and $t, t_{1}, t_{2}$ as above, let us consider the set $\Gamma^{\prime}(\gamma, \bar{x})$ of curves $\left(\gamma^{\prime}, t_{\gamma}^{\prime-}, t_{\gamma}^{\prime+}\right) \in \Gamma$ for which there exist $t_{1}^{\prime}<t_{2}^{\prime}$ in $\left(t_{y}^{\prime-}, t_{y}^{\prime+}\right)$ such that the following conditions are satisfied:
(i) $\gamma_{a}^{\prime}(t) \in I_{l}$ for every $t \in\left(t_{1}^{\prime}, t_{2}^{\prime}\right)$.
(ii) $\gamma_{x}^{\prime}\left(t_{1}^{\prime}\right) \cdot e_{l}=\bar{x} \cdot e_{l}$ and $\gamma_{x}^{\prime}\left(t_{1}^{\prime}\right) \cdot e_{l}^{\perp}>\bar{x} \cdot e_{l}^{\perp}$.
(iii) $\gamma_{x}^{\prime}\left(t_{2}^{\prime}\right) \cdot e_{l}<\gamma_{x}\left(t_{2}\right) \cdot e_{l}$.

By Proposition 4.2, it follows that

$$
\begin{equation*}
\omega_{e}\left(\Gamma^{\prime}(y, \bar{x})\right)=0 \tag{4.6}
\end{equation*}
$$

Since the functions $g_{\gamma}$ in (4.5) are equi-Lipschitz, the supremum in (4.5) can be realized by taking only countably many curves in $\Gamma_{l}^{+}(\bar{x})$. Therefore, it follows from (4.6) that

$$
\begin{equation*}
\omega_{e}\left(\left\{\left(\gamma, t_{\gamma}^{-}, t_{\gamma}^{+}\right) \in \Gamma_{l}^{+}(\bar{x}): \exists t_{2}^{\prime} \in\left(t_{1}^{\prime}, t_{\gamma}^{+}\right) \text {such that } \gamma_{a}\left(\left(t_{1}^{\prime}, t_{2}^{\prime}\right)\right) \subset I_{l} \text { and } \gamma_{x}\left(t_{2}^{\prime}\right) \cdot e_{l}^{\perp}<\bar{f}_{\bar{x}, l}\left(\gamma_{x}\left(t_{2}^{\prime}\right) \cdot e_{l}\right)\right\}\right)=0 \tag{4.7}
\end{equation*}
$$

Finally, since for every $\left(\gamma^{\prime}, t_{\gamma}^{\prime-}, t_{\gamma}^{\prime+}\right) \in \Gamma^{\prime}(\gamma, \bar{x})$ the associated map $g_{\gamma^{\prime}}:\left(y_{x}^{\prime}\left(t_{1}^{\prime}\right) \cdot e_{l}, \gamma_{x}^{\prime}\left(t_{2}^{\prime}\right) \cdot e_{l}\right) \rightarrow \mathbb{R}$ is Lipschitz with Lipschitz constant bounded by $\tan \left(\frac{3 \pi}{8}\right)$, we can replace $\bar{f}_{\bar{x}, l}$ with $f_{\bar{x}, l}$ in (4.7). This gives the second equation in (4.4) and it concludes the proof.

The following elementary lemma is about functions of bounded variation of one variable: we refer to [2] for the theory of BV functions.

Lemma 4.4. Let $v:(a, b) \rightarrow \mathbb{R}$ be $a \mathrm{BV}$ function and denote by $D^{-} v$ the negative part of the measure $D v$. Then for $\tilde{D}^{-} v$-a.e. $\bar{x} \in(a, b)$ there exists $\delta>0$ such that

$$
\bar{v}(x)>\bar{v}(\bar{x}) \quad \text { for all } x \in(\bar{x}-\delta, \bar{x}), \quad \text { and } \quad \bar{v}(x)<\bar{v}(\bar{x}) \quad \text { for all } x \in(\bar{x}, \bar{x}+\delta) .
$$

We are now in the position to prove the rectifiability of $v_{l}^{-}$.
Proposition 4.5. The measure $v_{l}^{-}$is concentrated on the set

$$
\bigcup_{\bar{x} \in \mathbb{Q}^{2} \cap B_{R}} C_{f_{\bar{x}, l}}, \quad \text { where } \quad C_{f_{\bar{x}, l}}:=B_{R} \cap \bigcup_{s>\bar{x} \cdot e_{l}^{\perp}}\left\{s e_{l}^{\perp}+f_{\bar{x}, l}(s) e_{l}\right\} .
$$

Proof. We prove this proposition in four steps.
Step 1. For every $\bar{x} \in B_{R} \cap \mathbb{Q}^{2}$ and every $\left(\gamma, t_{\gamma}^{-}, t_{\gamma}^{+}\right) \in \Gamma_{h}$, we consider the open set $I_{\bar{x}, l, y}^{+} \subset\left(t_{\gamma}^{-}, t_{\gamma}^{+}\right)$defined by the following property: we say that $t \in I_{\bar{x}, l, y}^{+}$if there exists $t^{\prime} \in\left(t_{\gamma}^{-}, t\right)$ such that

$$
\gamma_{a}\left(\left(t^{\prime}, t\right)\right) \subset I_{l}, \quad \gamma_{x}\left(t^{\prime}\right) \cdot e_{l}^{\perp}=\bar{x} \cdot e_{l}^{\perp}, \quad \gamma_{x}\left(t^{\prime}\right) \cdot e_{l}>\bar{x} \cdot e_{l} .
$$

Moreover, we set

$$
\mathcal{G}_{\bar{x}, l}^{>}:=\left\{\left(\gamma, t_{\gamma}^{-}, t_{\gamma}^{+}, t, x, a\right) \in \Gamma_{h} \times(0,1) \times B_{R} \times[0, M]: t \in I_{\bar{x}}^{+}, l, y\right\}
$$

Similarly, for every $\left(\gamma, t_{\gamma}^{-}, t_{\gamma}^{+}\right) \in \Gamma_{e}$, we let $I_{\bar{x}, l, \gamma}^{-} \subset\left(t_{\gamma}^{-}, t_{\gamma}^{+}\right)$be the set of $t$ for which there exists $t^{\prime} \in\left(t_{\gamma}^{-}, t\right)$ such that

$$
y_{a}\left(\left(t^{\prime}, t\right)\right) \subset I_{l}, \quad y_{x}\left(t^{\prime}\right) \cdot e_{l}^{\perp}=\bar{x} \cdot e_{l}^{\perp}, \quad y_{x}\left(t^{\prime}\right) \cdot e_{l}<\bar{x} \cdot e_{l}
$$

and we set

$$
\mathcal{G}_{\bar{\chi}, l}^{<}:=\left\{\left(\gamma, t_{\gamma}^{-}, t_{\gamma}^{+}, t, x, a\right) \in \Gamma_{e} \times(0,1) \times B_{R} \times[0, M]: t \in I_{\bar{\chi}}^{-}, l, \gamma\right\} .
$$

We consider

$$
\pi_{\bar{\chi}, l}^{-}:=\pi_{\llcorner }^{-}\left(\mathcal{G}_{\bar{\chi}, l}^{>} \times \mathcal{G}_{\bar{\chi}, l}^{<}\right)
$$

and we prove that $\left(p_{x}^{1}\right)_{\sharp} \pi_{\bar{x}, l}^{-}$is concentrated on $C_{f_{\bar{x}, l}}$, where

$$
\begin{aligned}
p_{x}^{1}:\left(\Gamma \times(0,1) \times B_{R} \times[0, M]\right)^{2} & \rightarrow B_{R}, \\
\left(\gamma, t_{\gamma}^{-}, t_{\gamma}^{+}, t, x, a, \gamma^{\prime}, t_{\gamma}^{-\prime}, t_{\gamma}^{+\prime}, t^{\prime}, x^{\prime}, a^{\prime}\right) & \mapsto x .
\end{aligned}
$$

Trivially, it holds

$$
\left(p_{x}^{1}\right)_{\sharp} \pi_{\bar{\chi}, l}^{-} \leq\left(p_{x}^{1}\right)_{\sharp}\left[\pi_{\llcorner }^{-}\left(\mathcal{G}_{\bar{\chi}, l}^{>} \times\left(\Gamma \times(0,1) \times B_{R} \times[0, M]\right)\right)\right] .
$$

From Corollary 4.3 it follows that for $\omega_{h}$-a.e. $\left(\gamma, t_{y}^{-}, t_{y}^{+}\right) \in \Gamma_{h}$ it holds

$$
\gamma_{x}(t) \cdot e_{l}^{\perp}>\bar{x} \cdot e_{l}^{\perp} \quad \text { and } \quad \gamma_{x}(t) \cdot e_{l} \geq f_{\bar{x}, l}\left(\gamma_{x}(t) \cdot e_{l}^{\perp}\right) \quad \text { for all } t \in I_{\bar{x}, l, \gamma}
$$

Therefore,

$$
\begin{equation*}
\left(p_{x}^{1}\right)_{\sharp} \pi_{\bar{x}, l}^{-}\left(\left\{x \in B_{R}: x \cdot e_{l}^{\perp} \leq \bar{x} \cdot e_{l}^{\perp}\right\} \cup\left\{x \in B_{R}: x \cdot e_{l}^{\perp}>\bar{x} \cdot e_{l}^{\perp} \text { and } x \cdot e_{l}<f_{\bar{x}, l}\left(x \cdot e_{l}^{\perp}\right)\right\}\right)=0 . \tag{4.8}
\end{equation*}
$$

In the same way, we get

$$
\begin{equation*}
\left(p_{x}^{2}\right)_{\sharp} \pi_{\bar{x}, l}^{-}\left(\left\{x \in B_{R}: x \cdot e_{l}^{\perp} \leq \bar{x} \cdot e_{l}^{\perp}\right\} \cup\left\{x \in B_{R}: x \cdot e_{l}^{\perp}>\bar{x} \cdot e_{l}^{\perp} \text { and } x \cdot e_{l}>f_{\bar{x}, l}\left(x \cdot e_{l}^{\perp}\right)\right\}\right)=0, \tag{4.9}
\end{equation*}
$$

where

$$
\begin{aligned}
p_{x}^{2}:\left(\Gamma \times(0,1) \times B_{R} \times[0, M]\right)^{2} & \rightarrow B_{R} \\
\left(\gamma, t_{\gamma}^{-}, t_{\gamma}^{+}, t, x, a, y^{\prime}, t_{y}^{-\prime},{t_{\gamma}^{+\prime}}^{+\prime}, t^{\prime}, x^{\prime}, a^{\prime}\right) & \mapsto x^{\prime}
\end{aligned}
$$

Finally, since $\pi^{-}$is concentrated on $\mathcal{G}$,

$$
\left(p_{x}^{1} \otimes p_{x}^{2}\right)_{\sharp} \pi^{-} \in \mathcal{M}\left(([0, T] \times \mathbb{R})^{2}\right)
$$

is concentrated on the graph of the identity on $B_{R}$ and in particular $\left(p_{x}^{1}\right)_{\sharp} \pi_{\bar{\chi}, l}^{-}=\left(p_{x}^{2}\right)_{\sharp} \pi_{\bar{\chi}, l}^{-}$. Therefore, it follows from (4.8) and (4.9) that $\left(p_{x}^{1}\right)_{\sharp} \pi_{\bar{x}, l}^{-}$is concentrated on

$$
\left\{x \in B_{R}: x \cdot e_{l}>\bar{x} \cdot e_{l} \text { and } x \cdot e_{l}^{\perp}=f_{\bar{x}, l}\left(x \cdot e_{l}\right)\right\}=C_{f_{\bar{x}, l}} .
$$

Step 2. We prove that for $\pi_{l}^{-}$-a.e.

$$
z=\left(\gamma, t_{\gamma}^{-}, t_{\gamma}^{+}, t, x, a, y^{\prime}, t_{\gamma}^{-\prime}, t_{\gamma}^{+^{\prime}}, t^{\prime}, x^{\prime}, a^{\prime}\right) \in\left(\Gamma \times(0,1) \times B_{R} \times[0, M]\right)^{2}
$$

there exists $\delta>0$ such that for every $s \in(t-\delta, t)$ and $s^{\prime} \in\left(t^{\prime}-\delta, t^{\prime}\right)$ the following properties hold:
(i) $\gamma_{a}(s) \in I_{l}$ and $\gamma_{a}(s)>a$.
(ii) $\gamma_{a}^{\prime}\left(s^{\prime}\right) \in I_{l}$ and $\gamma_{a}^{\prime}\left(s^{\prime}\right)<a^{\prime}$.

It is sufficient to prove the properties in (i) since the ones in (ii) can be shown analogously. The statement is trivial for elements $z$ for which $\gamma_{a}(t-)>a$, and it follows immediately by Lemma 4.4 applied to $\gamma_{a}$ if $\gamma_{a}$ is continuous at $t$. Since $\pi^{-}$is concentrated on points $Z$ for which $\gamma_{a}(t+) \leq \gamma_{a}(t-)$, it is sufficient to check that

$$
\pi_{l}^{-}\left(\left\{z \in\left(\Gamma \times(0,1) \times B_{R} \times[0, M]\right)^{2}: \gamma_{a}(t-)=a>\gamma_{a}(t+)\right\}\right)=0 .
$$

This follows immediately from the facts that for $\omega_{h}$-a.e. $\gamma$ the measure $\mu_{\gamma}^{-}$has no atoms and the set

$$
(t, x, a) \in(0,1) \times B_{R} \times[0, M]
$$

for which $\gamma_{x}(t)=x$ and $\gamma_{a}(t-)=a>\gamma_{a}(t+)$ is at most countable.
Step 3. We prove that for $\pi_{l}^{-}$-a.e. $z \in\left(\Gamma \times(0,1) \times B_{R} \times[0, M]\right)^{2}$ there exists $\bar{x} \in \mathbb{Q}^{2} \cap B_{R}$ such that

$$
z \in \mathcal{G}_{\bar{x}, l}^{>} \times \mathcal{G}_{\bar{x}, l}^{<}
$$

Let us consider $\delta>0$ from step 2. From property (i) and (3.1), it follows that for every $s \in(t-\delta, t)$ it holds

$$
\begin{equation*}
\gamma_{x}(s) \cdot e_{l}>\gamma_{x}(t) \cdot e_{l}-i e^{i a} \cdot e_{l}\left(\gamma_{x}(t) \cdot e_{l}^{\perp}-\gamma_{x}(s) \cdot e_{l}^{\perp}\right) \tag{4.10}
\end{equation*}
$$

and similarly for every $s^{\prime} \in\left(t^{\prime}-\delta, t^{\prime}\right)$,

$$
\begin{equation*}
\gamma_{x}^{\prime}\left(s^{\prime}\right) \cdot e_{l}<\gamma_{x}^{\prime}\left(t^{\prime}\right) \cdot e_{l}-i e^{i a^{\prime}} \cdot e_{l}\left(\gamma_{x}^{\prime}\left(t^{\prime}\right) \cdot e_{l}^{\perp}-\gamma_{x}^{\prime}\left(s^{\prime}\right) \cdot e_{l}^{\perp}\right) \tag{4.11}
\end{equation*}
$$

Since $\pi^{-}$is concentrated on $\mathcal{G}$, for $\pi_{l}^{-}$-a.e.

$$
z \in\left(\Gamma \times(0,1) \times B_{R} \times[0, M]\right)^{2}
$$

it also holds $a=a^{\prime}$ and $\gamma_{x}(t)=x=\gamma_{x}^{\prime}\left(t^{\prime}\right)$. Let us consider

$$
y \in\left(y_{x}(t) \cdot e_{l}^{\perp}-\frac{\delta}{100}, \gamma_{x}(t) \cdot e_{l}^{\perp}\right) \cap \sqrt{2} \mathbb{Q} .
$$

Then there exist $s \in(t-\delta, t)$ and $s^{\prime} \in\left(t^{\prime}-\delta, t^{\prime}\right)$ such that $\gamma_{x}(s) \cdot e_{l}^{\perp}=y=\gamma_{x}^{\prime}\left(s^{\prime}\right) \cdot e_{l}^{\perp}$. It follows from (4.10) and (4.11) that

$$
\begin{aligned}
y_{x}^{\prime}\left(s^{\prime}\right) \cdot e_{l} & <\gamma_{x}^{\prime}\left(t^{\prime}\right) \cdot e_{l}-i e^{i a^{\prime}} \cdot e_{l}\left(y_{x}^{\prime}\left(t^{\prime}\right) \cdot e_{l}^{\perp}-\gamma_{x}^{\prime}\left(s^{\prime}\right) \cdot e_{l}^{\perp}\right) \\
& =x \cdot e_{l}-i e^{i a^{\prime}} \cdot e_{l}\left(x \cdot e_{l}^{\perp}-y\right) \\
& =\gamma_{x}(t) \cdot e_{l}-i e^{i a} \cdot e_{l}\left(\gamma_{x}(t) \cdot e_{l}^{\perp}-\gamma_{x}(s) \cdot e_{l}^{\perp}\right) \\
& <\gamma_{x}(s) \cdot e_{l} .
\end{aligned}
$$

Let $z \in\left(\gamma_{x}^{\prime}\left(s^{\prime}\right) \cdot e_{l}, \gamma_{x}(s) \cdot e_{l}\right) \cap \sqrt{2} \mathbb{Q}$ and set $\bar{x}=z e_{l}+y e_{l}^{\perp}$. By construction, it holds

$$
\mathcal{Z} \in \mathcal{G}_{\bar{x}, l}^{>} \times \mathcal{G}_{\bar{x}, l}^{<}
$$

Since $e_{l}, e_{l}^{\perp} \in(\sqrt{2} \mathbb{Q})^{2}$, we obtain $\bar{x} \in \mathbb{Q}^{2}$.
Step 4. It follows by step 3 that

$$
\begin{equation*}
\pi_{l}^{-} \leq \pi_{\llcorner }^{-}\left(\bigcup_{\bar{x} \in \mathbb{Q}^{2} \cap B_{R}} \mathcal{G}_{\bar{x}, l}^{>} \times \mathcal{G}_{\bar{x}, l}^{<}\right) \tag{4.12}
\end{equation*}
$$

Since by step 1 we have that $\left(p_{x}^{1}\right)_{\sharp} \pi_{\bar{\chi}, l}^{-}$is concentrated on $C_{f_{\bar{\chi}, l}}$, the statement of the proposition follows from (4.12).

### 4.3 Rectifiability of $v_{\text {jump }}^{-}$

In the next lemma, we prove a regularity density estimate at a point $\bar{x}$ provided that the entropy dissipation measure decays faster than in a shock point.

Lemma 4.6. Let $\left(\bar{\gamma}, t_{\bar{\gamma}}^{-}, t_{\bar{\gamma}}^{+}\right) \in \Gamma_{h}, \bar{t} \in\left(t_{\bar{\gamma}}^{-}, t_{\bar{\gamma}}^{+}\right)$, and set $\bar{x}=\bar{\gamma}_{x}(\bar{t})$ and $\bar{a}=\bar{\gamma}_{a}(\bar{t}-) \vee \bar{\gamma}_{a}(\bar{t}+)$. Then there exists an absolute constant $c>0$ such that for every $\delta \in(0, \pi / 2)$ at least one of the following holds true:

$$
\begin{align*}
\liminf _{r \rightarrow 0} \frac{\mathcal{L}^{2}\left(\left\{x \in B_{r}(\bar{x}): \phi(x) \geq \bar{a}-\delta\right\}\right)}{r^{2}} & \geq c \delta, \\
\limsup _{r \rightarrow 0} \frac{\nu\left(B_{r}(\bar{x})\right)}{r} & \geq c \delta^{3} . \tag{4.13}
\end{align*}
$$

Proof. We assume without loss of generality that $\bar{a}=\bar{\gamma}_{a}(\bar{t}-)$, and we let $\delta_{1}>0$ be such that for every $t \in\left(\bar{t}-\delta_{1}, \bar{t}\right)$ it holds

$$
\bar{\gamma}_{a}(t) \in\left(\bar{a}-\frac{\delta}{5}, \bar{a}+\frac{\delta}{5}\right) .
$$

Moreover, we set $\bar{r}=\delta_{1} / 2$ so that for every $r \in(0, \bar{r})$ there exists $t_{r} \in\left(\bar{t}-\delta_{1}, \bar{t}\right)$ such that $\bar{\gamma}_{x}\left(t_{r}\right) \in \partial B_{r}(\bar{x})$ and $\gamma_{x}(t) \in B_{r}(\bar{x})$ for every $t \in\left(t_{r}, \bar{t}\right)$. Since $\bar{\gamma} \in \Gamma_{h}$ and $\bar{\gamma}_{a}(t) \geq \bar{a}-\delta / 5$ for every $t \in\left(t_{r}, \bar{t}\right)$, there exists $\varepsilon>0$ (possibly depending on $r$ ) such that

$$
\mathscr{L}^{2}\left(\left\{x \in S_{\varepsilon, r}: \phi(x) \geq \bar{a}-\frac{\delta}{5}\right\}\right) \geq \varepsilon r, \quad \text { where } \quad S_{\varepsilon, r}:=\bar{\gamma}_{x}\left(\left(t_{r}, \bar{t}\right)\right)+B_{\varepsilon}(0)
$$

For every $\left(\gamma, t_{\gamma}^{-}, t_{\gamma}^{+}\right) \in \Gamma$, we consider the nontrivial interiors $\left(t_{\gamma, i}^{-}, t_{\gamma, i}^{+}\right)_{i=1}^{N_{\gamma}}$ of the connected components of $\gamma_{a}^{-1}\left((\bar{a}-\delta), \bar{a}-\frac{2}{5} \delta\right)$, which intersect

$$
\gamma^{-1}\left(S_{\varepsilon, r} \times\left(\bar{a}-\frac{4}{5} \delta, \bar{a}-\frac{3}{5} \delta\right)\right) .
$$

Notice that we have the estimate

$$
N_{y} \leq 1+\frac{5}{\delta} \operatorname{Tot} \operatorname{Var} \gamma_{a}
$$

For every $i \in \mathbb{N}$, we consider

$$
\Gamma_{i}:=\left\{\left(y, t_{\gamma}^{-}, t_{\gamma}^{+}\right) \in \Gamma: N_{\gamma} \geq i\right\}
$$

and the measurable restriction map

$$
\begin{aligned}
R_{i}: \Gamma_{i} & \rightarrow \Gamma \\
\left(\gamma, t_{\gamma}^{-}, t_{\gamma}^{+}\right) & \mapsto\left(\gamma, t_{\gamma, i}^{-}, t_{\gamma, i}^{+}\right)
\end{aligned}
$$

We finally consider the measure

$$
\tilde{\omega}_{h}:=\sum_{i=1}^{\infty}\left(R_{i}\right)_{\sharp}\left(\omega_{h\llcorner } \Gamma_{i}\right) .
$$

We observe that $\tilde{\omega}_{h} \in \mathcal{M}_{+}(\Gamma)$ since for every $N>0$,

$$
\left\|\sum_{i=1}^{N}\left(R_{i}\right)_{\sharp}\left(\omega_{h\llcorner } \Gamma_{i}\right)\right\| \leq \int_{\Gamma} N_{y} d \omega_{h} \leq \int_{\Gamma}\left(1+\frac{5}{\delta} \operatorname{Tot} \operatorname{Var} \gamma_{a}\right) d \omega_{h}(\gamma)<\infty .
$$

The advantage of using the restrictions introduced above is in the following estimate: by an elementary transversality argument, there exists an absolute constant $\tilde{c}>0$ such that for $\tilde{\omega}_{h}$-a.e. $\left(\gamma, t_{\gamma}^{-}, t_{\gamma}^{+}\right) \in \Gamma$ it holds

$$
\begin{equation*}
\mathscr{L}^{1}\left(\left\{t \in\left(t_{\gamma}^{-}, t_{\gamma}^{+}\right): \gamma(t) \in S_{\varepsilon, r} \times\left(\bar{a}-\frac{4}{5} \delta, \bar{a}-\frac{3}{5} \delta\right)\right\}\right) \leq \tilde{c} \frac{\varepsilon}{\delta} . \tag{4.14}
\end{equation*}
$$

By construction, we have that for every $t \in(0,1)$ it holds

$$
\begin{equation*}
\left(e_{t}\right)_{\sharp} \tilde{\omega}_{h} \geq \mathscr{L}^{3}\left\llcorner\left\{(x, a) \in S_{\varepsilon, r} \times\left(\bar{a}-\frac{4}{5} \delta, \bar{a}-\frac{3}{5} \delta\right): \phi(x) \geq a\right\} .\right. \tag{4.15}
\end{equation*}
$$

Since the measure of this set is at least $\varepsilon r \delta / 5$, it follows by (4.14) and (4.15) that

$$
\begin{equation*}
\tilde{\omega}_{h}(\Gamma) \geq \varepsilon r \frac{\delta}{5} \cdot \frac{\delta}{\tilde{c} \varepsilon}=\frac{\delta^{2}}{5 \tilde{c}} r \tag{4.16}
\end{equation*}
$$

We consider $\Gamma=\Gamma_{1} \cup \Gamma_{2}$, where

$$
\Gamma_{1}:=\left\{\left(\gamma, t_{\gamma}^{-}, t_{\gamma}^{+}\right) \in \Gamma: t_{\gamma}^{+}-t_{\gamma}^{-} \geq r\right\} \quad \text { and } \quad \Gamma_{2}:=\left\{\left(\gamma, t_{\gamma}^{-}, t_{\gamma}^{+}\right) \in \Gamma: t_{\gamma}^{+}-t_{\gamma}^{-}<r\right\}
$$

For $\tilde{\omega}_{h}$-a.e. $\left(\gamma, t_{\gamma}^{-}, t_{\gamma}^{+}\right) \in \Gamma_{1}$ it holds

$$
\mathscr{L}^{1}\left(\left\{t \in\left(t_{\gamma}^{-}, t_{\gamma}^{+}\right): y(t) \in B_{2 r}(\bar{x}) \times\left(\bar{a}-\delta, \bar{a}-\frac{2}{5} \delta\right)\right\}\right) \geq r
$$

while for $\tilde{\omega}_{h}$-a.e. $\left(\gamma, t_{\gamma}^{-}, t_{\gamma}^{+}\right) \in \Gamma_{2}$ we have

$$
\gamma_{x}\left(t_{y}^{-}, t_{y}^{+}\right) \subset B_{2 r}(\bar{x}) \quad \text { and } \quad \operatorname{TotVar} \gamma_{a} \geq \frac{2}{5} \delta
$$

It follows from (4.16) that at least one of the following holds:

$$
\begin{equation*}
\tilde{\omega}_{h}\left(\Gamma_{1}\right) \geq \frac{\delta^{2}}{10 \tilde{c}} r \quad \text { or } \quad \tilde{\omega}_{h}\left(\Gamma_{2}\right) \geq \frac{\delta^{2}}{10 \tilde{c}} r . \tag{4.17}
\end{equation*}
$$

If the second condition holds, then we have that

$$
v\left(B_{2 r}(\bar{x})\right) \geq\left|U_{\phi}\right|\left(B_{2 r}(\bar{x}) \times(\bar{a}-\delta, \bar{a})\right) \geq \frac{\delta^{3}}{25 \tilde{c}} r
$$

so that the second condition in the statement is satisfied. Otherwise, we assume that the first condition in (4.17) holds: since for every $t \in(0,1)$,

$$
\left(e_{t}\right)_{\sharp} \tilde{\omega}_{h} \leq \chi \mathscr{L}^{3}
$$

it follows from (4.17) and the Fubini theorem that

$$
\mathscr{L}^{2}\left(\left\{x \in B_{2 r}(\bar{x}): \phi(x) \geq \bar{a}-\delta\right\}\right) \geq \frac{\delta^{2}}{10 \tilde{c}} r \cdot \frac{5 r}{3 \delta}=\frac{\delta}{6 \tilde{c}} r^{2}
$$

so that the first condition in the statement holds true.
Remark 4.7. We observe that the third power in (4.13) is optimal; this is related to the fact that the optimal regularity of $\phi$ is $B_{\infty, \text { loc }}^{1 / 3,3}(\Omega)$; see [12].
We also state the same result for curves in $\Gamma_{e}$, whose proof is analogous to the one of Lemma 4.6.

Lemma 4.8. Let $\left(y, t_{\gamma}^{-}, t_{y}^{+}\right) \in \Gamma_{e}, t \in\left(t_{y}^{-}, t_{\gamma}^{+}\right)$, and set $\bar{x}=\gamma_{x}(t)$ and $\bar{a}=\gamma_{a}(t-) \wedge \gamma_{a}(t+)$. Then there exists an absolute constant $c>0$ such that for every $\delta \in(0, \pi / 2)$ at least one of the following holds true:

$$
\begin{aligned}
\liminf _{r \rightarrow 0} \frac{\mathcal{L}^{2}\left(\left\{x \in B_{r}(\bar{x}): \phi(x) \leq \bar{a}+\delta\right\}\right)}{r^{2}} & \geq c \delta, \\
\limsup _{r \rightarrow 0} \frac{v\left(B_{r}(\bar{x})\right)}{r} & \geq c \delta^{3} .
\end{aligned}
$$

The main result of this section is the following proposition.
Proposition 4.9. For $v_{\text {jump }}^{-}$a.e. $x \in B_{R}$,

$$
\begin{equation*}
\limsup _{r \rightarrow 0} \frac{v\left(B_{r}(x)\right)}{r}>0 . \tag{4.18}
\end{equation*}
$$

Proof. For $v_{\text {jump }}^{-}$-a.e. $\bar{x} \in B_{R}$, one of the following assertions holds:
(i) There exist $\left(\gamma, t_{\gamma}^{-}, t_{\gamma}^{+}, t, x, a\right) \in \mathcal{G}_{h, \text { jump }}^{-}$and $\left(y^{\prime}, t_{\gamma}^{-\prime}, t_{\gamma}^{+\prime}, t^{\prime}, x^{\prime}, a^{\prime}\right) \in \Gamma_{e}$ such that

$$
x=x^{\prime}=\bar{x} \quad \text { and } \quad \gamma_{a}^{\prime}\left(t^{\prime}+\right) \leq a^{\prime}=a<\gamma_{a}(t-) .
$$

(ii) There exist $\left(y^{\prime}, t_{\gamma}^{-\prime}, t_{\gamma}^{+\prime}, t^{\prime}, x^{\prime}, a^{\prime}\right) \in \mathcal{G}_{e, \text { jump }}^{+}$and $\left(\gamma, t_{\gamma}^{-}, t_{\gamma}^{+}, t, x, a\right) \in \Gamma_{e}$ such that

$$
x=x^{\prime}=\bar{x} \quad \text { and } \quad y_{a}^{\prime}\left(t^{\prime}+\right)<a^{\prime}=a \leq \gamma_{a}(t-) .
$$

Since the two cases are equivalent, we consider only the first one. We apply Lemma 4.6 to the curve $\gamma$ and Lemma 4.8 to the curve $\gamma^{\prime}$ with $\delta=\left(\gamma_{a}(t-)-a\right) / 3$. If condition (ii) holds in at least one of the two cases, then the statement follows; otherwise both of the following inequalities are satisfied:

$$
\liminf _{r \rightarrow 0} \frac{\left\{x \in B_{r}(\bar{x}): \phi(x) \geq \gamma_{a}(t-)-\delta\right\}}{r^{2}} \geq c \delta^{2}, \quad \liminf _{r \rightarrow 0} \frac{\left\{x \in B_{r}(\bar{x}): \phi(x) \leq \gamma_{a}(t-)-2 \delta\right\}}{r^{2}} \geq c \delta^{2} .
$$

This condition excludes that $\bar{x}$ is a point of vanishing mean oscillation of $\phi$. Therefore, $\bar{x} \in \Sigma$ by Theorem 1.2, i.e. (4.18) holds true.

### 4.4 Conclusion

Collecting the results in Sections 4.2 and 4.3, we obtain the rectifiability of the measure $\left(p_{x}\right)_{\sharp} U_{\phi}^{-}$.
Proposition 4.10. The measure $\left(p_{x}\right)_{\sharp} U_{\phi}^{-}$is 1-rectifiable.
Proof. First, we observe that since $\pi^{-}$is concentrated on $\mathcal{G}$ and

$$
\mathcal{G} \subset\left(\mathcal{G}_{h, \text { jump }}^{-} \times \mathcal{G}_{e}^{+}\right) \cup\left(\mathcal{G}_{h}^{-} \times \mathcal{G}_{e, \text { jump }}^{+}\right) \cup\left(\bigcup_{l=0}^{L}\left(\mathcal{G}_{h, l}^{-} \times \mathcal{G}_{e, l}^{+}\right)\right),
$$

it follows from the definitions of $\pi_{l}^{-}$and $\pi_{\mathrm{jump}}^{-}$that

$$
\pi^{-} \leq \pi_{\mathrm{jump}}^{-}+\sum_{l=0}^{L} \pi_{l}^{-}
$$

In particular,

$$
\left(p_{x}\right)_{\sharp} U_{\phi}^{-}=\left(p_{x}^{1}\right)_{\sharp} \pi^{-} \leq\left(p_{x}^{1}\right)_{\sharp} \pi_{\mathrm{jump}}^{-}+\sum_{l=0}^{L}\left(p_{x}^{1}\right)_{\sharp} \pi_{l}^{-} .
$$

Since $\left(p_{x}^{1}\right)_{\sharp} \pi_{l}^{-}$is 1-rectifiable for every $l=0, \ldots, L$ by Proposition 4.5 and $\left(p_{x}^{1}\right)_{\sharp} \pi_{\mathrm{jump}}^{-}$is 1-rectifiable by Proposition 4.9 and Theorem 1.2, also $\left(p_{x}\right)_{\sharp} U_{\phi}^{-}$is 1-rectifiable.
As mentioned at the beginning of this section, the rectifiability of the positive part $\left(p_{x}\right)_{\sharp} U_{\phi}^{+}$can be proven following the same procedure. Therefore, this concludes the proof of Theorem 1.3.

Funding: The author has been supported by the SNF Grant 182565.

## References

[1] L. Ambrosio and G. Crippa, Existence, uniqueness, stability and differentiability properties of the flow associated to weakly differentiable vector fields, in: Transport Equations and Multi-D Hyperbolic Conservation Laws, Lect. Notes Unione Mat. Ital. 5, Springer, Berlin (2008), 3-57.
[2] L. Ambrosio, N. Fusco and D. Pallara, Functions of Bounded Variation and Free Discontinuity Problems, Oxford Math. Monogr., The Clarendon Press, New York, 2000.
[3] L. Ambrosio and N. Gigli, A user's guide to optimal transport, in: Modelling and Optimisation of Flows on Networks, Lecture Notes in Math. 2062, Springer, Heidelberg (2013), 1-155.
[4] L. Ambrosio, B. Kirchheim, M. Lecumberry and T. Rivière, On the rectifiability of defect measures arising in a micromagnetics model, in: Nonlinear Problems in Mathematical Physics and Related Topics. II, Int. Math. Ser. (N. Y.) 2, Kluwer/Plenum, New York (2002), 29-60.
[5] L. Ambrosio, M. Lecumberry and T. Rivière, A viscosity property of minimizing micromagnetic configurations, Comm. Pure Appl. Math. 56 (2003), no. 6, 681-688.
[6] P. Aviles and Y. Giga, A mathematical problem related to the physical theory of liquid crystal configurations, in: Miniconference on Geometry and Partial Differential Equations. 2 (Canberra 1986), Proc. Centre Math. Anal. Austral. Nat. Univ. 12, Australian National University, Canberra (1987), 1-16.
[7] S. Bianchini, P. Bonicatto and E. Marconi, A lagrangian approach to multidimensional conservation laws, preprint (2017), SISSA 36/MATE.
[8] S. Bianchini and S. Daneri, On Sudakov's type decomposition of transference plans with norm costs, Mem. Amer. Math. Soc. 251 (2018), no. 1197, 1-112.
[9] P. Billingsley, Convergence of Probability Measures, 2nd ed., Wiley Ser. Probab. Stat., John Wiley \& Sons, New York, 1999.
[10] C. De Lellis and F. Otto, Structure of entropy solutions to the eikonal equation, J. Eur. Math. Soc. (JEMS) 5 (2003), no. 2, 107-145.
[11] C. De Lellis, F. Otto and M. Westdickenberg, Structure of entropy solutions for multi-dimensional scalar conservation laws, Arch. Ration. Mech. Anal. 170 (2003), no. 2, 137-184.
[12] F. Ghiraldin and X. Lamy, Optimal Besov differentiability for entropy solutions of the eikonal equation, Comm. Pure Appl. Math. 73 (2020), no. 2, 317-349.
[13] P.-E. Jabin and B. Perthame, Compactness in Ginzburg-Landau energy by kinetic averaging, Comm. Pure Appl. Math. 54 (2001), no. 9, 1096-1109.
[14] X. Lamy and F. Otto, On the regularity of weak solutions to Burgers' equation with finite entropy production, Calc. Var. Partial Differential Equations 57 (2018), no. 4, Article ID 94.
[15] M. Lecumberry, Geometric structure of micromagnetic walls and shock waves in scalar conservation laws, PhD thesis, Université de Nantes, 2004.
[16] M. Lecumberry, Geometric structure of magnetic walls, in: Journées "Équations aux Dérivées Partielles", École Polytechnique, Palaiseau (2005), Exp. No. I.
[17] P.-L. Lions, B. Perthame and E. Tadmor, A kinetic formulation of multidimensional scalar conservation laws and related equations, J. Amer. Math. Soc. 7 (1994), no. 1, 169-191.
[18] E. Marconi, On the structure of weak solutions to scalar conservation laws with finite entropy production, preprint (2019), https://arxiv.org/abs/1909.07257.
[19] E. Marconi, The rectifiability of the entropy defect measure for Burgers equation, preprint (2020), https://arxiv.org/abs/2004.09932.
[20] T. Rivière and S. Serfaty, Limiting domain wall energy for a problem related to micromagnetics, Comm. Pure Appl. Math. 54 (2001), no. 3, 294-338.
[21] T. Rivière and S. Serfaty, Compactness, kinetic formulation, and entropies for a problem related to micromagnetics, Comm. Partial Differential Equations 28 (2003), no. 1-2, 249-269.
[22] C. Villani, Optimal Transport. Old and New, Grundlehren Math. Wiss. 338, Springer, Berlin, 2009.


[^0]:    *Corresponding author: Elio Marconi, EPFL B, Station 8, CH-1015 Lausanne, Switzerland, e-mail: elio.marconi@epfl.ch. https://orcid.org/0000-0003-2989-2594

