# An almost dark galaxy with the mass of the Small Magellanic Cloud ${ }^{\star}$ 

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#### Abstract

Almost dark galaxies are objects that have eluded detection by traditional surveys such as the Sloan Digital Sky Survey (SDSS). The low surface brightness of these galaxies $\left(\mu_{r}(0)>26 \mathrm{mag} \mathrm{arcsec}^{-2}\right.$ ), and hence their low surface stellar mass density (a few solar masses per $\mathrm{pc}^{2}$ or less), suggest that the energy density released by baryonic feedback mechanisms is inefficient in modifying the distribution of the dark matter halos they inhabit. For this reason, almost dark galaxies are particularly promising for probing the microphysical nature of dark matter. In this paper, we present the serendipitous discovery of Nube, an almost dark galaxy with $\left\langle\mu_{V}\right\rangle_{\mathrm{e}} \sim 26.7 \mathrm{mag}_{\mathrm{arcsec}}{ }^{-2}$. The galaxy was identified using deep optical imaging from the IAC Stripe82 Legacy Project. Follow-up observations with the 100 m Green Bank Telescope strongly suggest that the galaxy is at a distance of 107 Mpc . Ultra-deep multi-band observations with the 10.4 m Gran Telescopio Canarias favour an age of $\sim 10 \mathrm{Gyr}$ and a metallicity of $[\mathrm{Fe} / \mathrm{H}] \sim-1.1$. With a stellar mass of $\sim 4 \times 10^{8} M_{\odot}$ and a half-mass radius of $R_{\mathrm{e}}=6.9 \mathrm{kpc}$ (corresponding to an effective surface density of $\langle\Sigma\rangle_{\mathrm{e}} \sim 0.9 M_{\odot} \mathrm{pc}^{-2}$ ), Nube is the most massive and extended object of its kind discovered so far. The galaxy is ten times fainter and has an effective radius three times larger than typical ultradiffuse galaxies with similar stellar masses. Galaxies with comparable effective surface brightness within the Local Group have very low mass (tens of $10^{5} M_{\odot}$ ) and compact structures (effective radius $R_{\mathrm{e}}<1 \mathrm{kpc}$ ). Current cosmological simulations within the cold dark matter scenario, including baryonic feedback, do not reproduce the structural properties of Nube. However, its highly extended and flattened structure is consistent with a scenario where the dark matter particles are ultralight axions with a mass of $m_{B}=\left(0.8_{-0.2}^{+0.4}\right) \times 10^{-23} \mathrm{eV}$.


Key words. galaxies: dwarf - dark matter - galaxies: photometry - galaxies: structure - galaxies: formation

## 1. Introduction

Many cosmological observations at large scales suggest that dark matter can be well described as a cold and collisionless fluid (see e.g. White \& Rees 1978; Blumenthal et al. 1984; Davis et al. 1985; Smoot et al. 1992). Nonetheless, the predictions of this model at galactic scales have faced an increasing number of challenges, such as the "cusp-core" problem, the "missing satellite" problem, and the "too-big-to-fail" problem (see e.g. Boylan-Kolchin et al. 2011; Weinberg et al. 2015; Del Popolo \& Le Delliou 2017). Many of these problems can be mitigated by the effect of baryon feedback on the dark

[^0]matter distribution (see e.g. Davis et al. 1992; Governato et al. 2010; Di Cintio et al. 2014). However, if the number of stars or their spatial density were low enough, it would be difficult to argue that stellar feedback could be responsible for affecting the dark matter distribution, because there would not be enough energy to change the location of the dark matter (see e.g. Peñarrubia et al. 2012; Oñorbe et al. 2015). In parallel, while the direct detection of dark matter particles remains out of reach, other alternatives to the cold dark matter model have gained traction, and are being applied to solve the smallscale challenges. These include the warm dark matter scenario (see e.g. Sommer-Larsen \& Dolgov 2001; Bode et al. 2001), self-interacting dark matter (Spergel \& Steinhardt 2000), and fuzzy dark matter (composed of ultralight axions with masses in the $10^{-23}-10^{-21} \mathrm{eV}$ range; see e.g. Sin 1994; Hu et al. 2000;


Fig. 1. SDSS RGB image of the $20^{\prime} \times 20^{\prime}$ region around Nube. The insets show a zoom in on the galaxy, showing how it appears on images with different limiting surface brightness in the $r$-band ( $26.7,28.6$ and $30.5 \mathrm{mag} \operatorname{arcsec}^{-2} ; 3 \sigma$ on $10^{\prime \prime} \times 10^{\prime \prime}$ boxes). Previously known galaxies in the field of view are labelled.

Matos \& Arturo Ureña-López 2001). For these reasons, the search for objects with extremely low stellar surface densities (where the effect of baryonic feedback is not expected to be relevant) promises to probe the microphysical nature of dark matter, that is, the properties of the dark matter particle.

In recent years, our ability to detect more and more diffuse galaxies with broadband imaging has improved considerably (see e.g. Sandage \& Binggeli 1984; Conselice et al. 2003; van Dokkum et al. 2015; Román \& Trujillo 2017; Lim et al. 2020; Tanoglidis et al. 2021; Trujillo et al. 2021; Marleau et al. 2021; La Marca et al. 2022; Zaritsky et al. 2023), and is expected to continue to improve with the arrival of very deep optical surveys (see e.g. Ivezić et al. 2019). Within this population of very faint galaxies, the so-called "almost dark" galaxies are of particular interest. These faint galaxies are missed in the optical catalogues of wide field surveys such as the Sloan Digital Sky Survey (SDSS, Eisenstein et al. 2011). Although there is no definition of the surface brightness of an almost dark galaxy, given the SDSS surface brightness limit (i.e. $\mu_{r} \sim 26.5-27 \mathrm{mag} \operatorname{arcsec}^{-2} ; 3 \sigma$ in $10^{\prime \prime} \times 10^{\prime \prime}$ boxes), galaxies with central surface brightness of fainter than $\mu_{r}(0) \sim$ $26 \mathrm{mag} \mathrm{arcsec}^{-2}$ are very difficult to detect in SDSS catalogues. Therefore, we can use this value as a rough definition of what an almost dark galaxy should be. These galaxies represent less than $1 \%$ of the galaxies found in blind HI surveys such as Arecibo Legacy Fast Arecibo L-band Feed Array (ALFALFA, Giovanelli et al. 2005), and have HI masses of between $10^{7}$ and
$10^{9} M_{\odot}$. As the definition of an almost dark galaxy depends on whether it is typically detected in the SDSS, its maximum stellar surface mass density is a function of its stellar population properties (age and metallicity). In a scenario where the age of the stellar population is old ( $\sim 10 \mathrm{Gyr}$ ) and metal-poor $([\mathrm{Fe} / \mathrm{H}] \sim-1)$, a surface brightness of $\mu_{r}(0) \sim 26 \mathrm{mag} \operatorname{arcsec}^{-2}$ corresponds to a few $M_{\odot} \mathrm{pc}^{-2}$. This is very close to the expected stellar surface density resulting from the gas density threshold for star formation (Schaye 2004) and is in good agreement with the values found at the edges of galaxies where star formation suddenly drops off (Trujillo et al. 2020; Chamba et al. 2022). Consequently, almost dark galaxies may also be an interesting place to study how galaxy formation occurs at low densities.

In this paper, we describe the structural properties of a very extended almost dark galaxy serendipitously discovered in the IAC Stripe82 Legacy Project (Fliri \& Trujillo 2016; Román \& Trujillo 2018). This object (which we have named Nube ${ }^{1}$ ) was found during a visual inspection of one of the survey fields. The object is not visible in SDSS and appears quite noisy even in deeper images such as those produced by the Stripe82 data (see Fig. 1). Dedicated observations with the Gran Telescopio Canarias (GTC) and the Green Bank Telescope (GBT) have allowed us to characterise its nature in detail. This paper presents a comprehensive analysis of the stellar populations and structural properties of this object. We also consider

[^1]the properties of this very faint galaxy within the cold dark matter and fuzzy dark matter scenarios.

Throughout this work, we adopt a standard cosmological model with the following parameters: $H_{0}=70 \mathrm{~km} \mathrm{~s}^{-1} \mathrm{Mpc}^{-1}$, $\Omega_{\mathrm{m}}=0.3$, and $\Omega_{\Lambda}=0.7$. Based on the probable redshift of Nube, its assumed distance is 107 Mpc , corresponding to a spatial scale of $0.5039 \mathrm{kpc} \operatorname{arcsec}^{-1}$. All magnitudes in this paper are in the AB magnitude system.

## 2. Data

The data used in this paper come from two different facilities: the 10.4 m GTC and the 100 m Robert C. Byrd GBT. The details of each observation are described below.

### 2.1. GTC HiPERCAM images

Deep optical multi-band imaging of Nube was performed using HiPERCAM (Dhillon et al. 2018), which is a quintuplebeam, high-speed astronomical imager capable of simultaneously imaging celestial objects in five different Sloan filters $(u, g, r, i, z)$. The image area of each of the five CCDs is $2048 \times$ 1024 pixels ( $2.7^{\prime} \times 1.4^{\prime} ; 1$ pixel $=0!08$ ), and each is divided into four channels of $1024 \times 512$ pixels. Nube was observed on 9 and 10 January 2019. We followed the dithering strategy described in Trujillo \& Fliri (2016) to reduce the scattered light from the telescope structure as much as possible.

The data reduction of the Nube images follows the steps detailed in Montes et al. (2020) and is briefly described here. The entire reduction process was carried out in a controlled and enclosed environment as described in Akhlaghi et al. (2021). After the standard calibration per CCD channel (bias and flat field), each set of four channels was assembled into a single image. The different exposures that went into the final images were visually inspected, and those with low quality (too noisy and/or very bright background) or strong gradients were discarded. Photometric calibration of these images was performed using SDSS DR12 (Alam et al. 2015). The size of Nube on the sky (effective diameter of $\sim 27^{\prime \prime}$; see Fig. 1) is significantly smaller than the FOV of the camera ( $162^{\prime \prime} \times 84^{\prime \prime}$ ), allowing reliable background subtraction and study of the galaxy using these images. The background subtraction was performed by subtracting a calculated constant value from the masked image. In addition to masking the foreground and background sources, we masked a circular region with a radius of $40^{\prime \prime}$ centred on Nube - which is $\sim 3 R_{\mathrm{e}}$ of this galaxy; see Sect. 3.1 - to avoid over-subtracting the light associated with the outer parts of the galaxy.

The final exposure time on-source is 1 h and 8 min for each band. The limiting surface brightness depths of the final images measured in $10^{\prime \prime} \times 10^{\prime \prime}$ boxes are $30.5,31,30.5,30$, and $29.2 \mathrm{mag}_{\operatorname{arcsec}}{ }^{-2}$ for $u, g, r, i$, and $z$ ( $3 \sigma$ above the background), respectively, which we measured using the method described in Appendix A of Román et al. (2020). The limiting magnitudes for point-like sources are 26.2, 26.7, 26.2, 25.8, and 25 mag from $u$ to $z\left(5 \sigma\right.$ within an aperture of $\left.2^{\prime \prime}\right)$.

### 2.2. GBT HI data

We performed $\sim 12 \mathrm{~h}$ of observations with the $\mathrm{GBT}^{2}$ along the line of sight (LOS) to Nube between January 2019 and

[^2]July 2019 (programmes GBT18B-356 and GBT19A-485). Our observational configuration is identical to that of Karunakaran et al. (2020a), where we used the $L$-band receiver and the VErsatile GBT Astronomical Spectrometer (VEGAS) in Mode 7 (spectral resolution $=3 \mathrm{kHz} \sim 0.7 \mathrm{~km} \mathrm{~s}^{-1}$, bandpass $=100 \mathrm{MHz})$. The use of such a wide bandpass, which is sensitive out to recession velocities of $\sim 14000 \mathrm{~km} \mathrm{~s}^{-1}$, is essential in the blind search for HI in such a faint object, which could be at any distance along the LOS.

These data were reduced using the standard GBTIDL ${ }^{3}$ procedure getps. We removed both narrowband and broadband radio frequency interference (RFI) before searching for potential HI signals. We note that broadband RFI (i.e. GPS L3, $\mu 1.381 \mathrm{GHz}$ ) was particularly prominent, resulting in almost one-third of the data being flagged. The effective integration time of the final calibrated spectrum is therefore $\sim 8.5 \mathrm{~h}$. We also note that the fluxes are scaled up by a factor of $20 \%$ due to the systematic offset in the GBT noise diode calibration values (Goddy et al. 2020). We smoothed the resulting RFI-free spectrum at various spectral resolutions to search for possible HI emission. Figure 2 shows a faint HI detection at $V_{\text {Helio }}=7480 \pm 8 \mathrm{~km} \mathrm{~s}^{-1}$ (red line) with a signal-to-noise ratio of 6.1 , along with the statistically independent XX and YY polarisations (blue and green lines, respectively) that were co-added to produce the total spectrum. Although faint, the line is detected in both XX and YY, and is also consistently detected in spectra constructed from subsets of the data that are selected on observables such as observing date and time. For a number of reasons, which we explain below, we believe that this detection is most likely associated with Nube.

We estimate the systemic heliocentric velocity ( $V_{\text {Helio }}$ ) and velocity width ( $W_{50}$ ) of the HI detection using the methods described in Karunakaran et al. (2020a,b), which are based on those of Springob et al. (2005). Briefly, we fit first-order polynomials to each edge of the HI profile between $15 \%$ and $85 \%$ of the peak flux and find the velocities corresponding to the $50 \%$ flux value. The mean and difference of these velocities give $V_{\text {Helio }}$ and $W_{50}$, respectively. We correct $W_{50}$ for instrumental and cosmological redshift broadening according to Springob et al. (2005) and for ISM turbulence according to Verheijen \& Sancisi (2001, see their Sect. 4; see also Karunakaran et al. 2020a).

We convert $V_{\text {Helio }}$ to a kinematic distance ( $D$ ) using the Hubble-Lemaître law and assuming an uncertainty of 5 Mpc due to peculiar velocities. The derived distance is $107 \pm 5 \mathrm{Mpc}$. We measure the HI flux $\left(S_{\mathrm{HI}}\right)$ by integrating over the line profile, the uncertainties on which are dominated by the noise and the $2 \%$ diode uncertainty (van Zee et al. 1997). The flux value is $S_{\mathrm{HI}}=0.083 \pm 0.02 \mathrm{Jy} \mathrm{km} \mathrm{s}^{-1}$. We use the kinematic distance together with the integrated flux in the standard equation for an optically thin gas (Haynes \& Giovanelli 1984) to calculate the HI mass:

$$
\begin{equation*}
M_{\mathrm{HI}}=2.356 \times 10^{5} D^{2} S_{\mathrm{HI}}\left[M_{\odot}\right] \tag{1}
\end{equation*}
$$

resulting in $M_{\mathrm{HI}}=2.2_{-0.5}^{+0.7} \times 10^{8} M_{\odot}$. All derived properties from the HI spectrum are at a spectral resolution of $25 \mathrm{~km} \mathrm{~s}^{-1}$ and are listed in Table 1.

### 2.3. Possible sources of contamination in HI

Although it is tempting to associate the HI detection with the Nube galaxy, it could be that the measurement corresponds to another galaxy in the LOS. For this reason, a detailed analysis

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Fig. 2. HI detection along the LOS to Nube. We also show different polarisations of the data, i.e. XX and YY, in order to better identify which other HI emission peaks in the spectrum are potentially spurious. The black tick indicates the location of the peak with the strongest signal in our data. We note that other distinct peaks are not always present in the two different polarisations, which reduces their relevance as real signals.
of the velocity distributions of the galaxies in the Nube region is needed to see if the HI detection is compatible with any of them.

Using $\mathrm{NED}^{4}$, we created a catalogue of all objects in this database with spectroscopic redshifts in a region of $100^{\prime} \times 100^{\prime}$ centred on Nube. The total number of galaxies found is 711. All galaxies with recession velocities in the range of $5480<$ $V_{\text {helio }}<9480 \mathrm{~km} \mathrm{~s}^{-1}$, that is 42 galaxies, are plotted in Fig. 3. The galaxies are colour coded according to the relative velocity of the HI detection. In this region, there are many galaxies with relative velocities of $\Delta V<-600 \mathrm{~km} \mathrm{~s}^{-1}$ from the HI detection (red symbols in the figure). There is also a smaller group with velocities of about $500 \mathrm{~km} \mathrm{~s}^{-1}$ greater than the HI detection (purple symbols close to the centre). Interestingly, there is no galaxy with known spectroscopic redshift within a radius of $10^{\prime}$ centred on Nube that has $|\Delta V|<500 \mathrm{~km} \mathrm{~s}^{-1}$. This is important because this is the region where potential contamination from a gas-rich source could mimic the signal we identify as possibly coming from Nube. To better illustrate this point, in Fig. 3 we also show the 2D GBT beam response using the well-characterised pattern at 1.4 GHz from Spekkens et al. (2013). The contours (from blue to black) indicate where the beam efficiency drops by the factor indicated on the label.

The closest galaxy in terms of velocity ( $\Delta V=-17.8 \mathrm{~km} \mathrm{~s}^{-1}$ ) would be the spiral galaxy to the north, UGC 929 , shown in yellow in Fig. 3. This galaxy is 14.7 arcmin or 435 kpc away at the distance corresponding to the HI detection (i.e. 107 Mpc ). This object is close to the first minimum of the beam response - shown by the grey-scale background image in Fig. 3 - which means that the HI emission, if coming from UGC 929, would be suppressed by a factor of $\sim 10^{4}$. If the source of the HI emission is UGC 929, this implies that the HI mass of this galaxy is $\sim 10^{12} M_{\odot}$. Using scaling relations between stellar mass and HI mass, a galaxy with the stellar mass of UGC 929 ( $M_{*}=$ $2.4 \times 10^{10} M_{\odot}$, see Appendix A) is expected to have a mass in HI of $\sim 2 \times 10^{9} M_{\odot}$ and no more than $\sim 2 \times 10^{10} M_{\odot}$ (Feldmann

[^4]2020). This is at least 100 times less gas than what it is needed to reproduce the observations. It is therefore very unlikely that the emission we see at the location of Nube is caused by contamination from UGC 929. To the northwest in Fig. 3, there is another galaxy at a distance of $14^{\prime}$ that has a similar velocity $\left(\Delta V=-27.8 \mathrm{~km} \mathrm{~s}^{-1}\right)$. This galaxy appears to be a low-mass galaxy, probably a satellite of UGC 929. Given it is also located in the first minimum of the beam response and its low mass, it is also very unlikely to be the source of the HI emission.

None of the galaxies with a confirmed spectroscopic redshift appear to be responsible for the detection of HI in the Nube region. However, it is possible that the HI emission is coming from a low-mass galaxy in the region that, due to its faintness, has no spectroscopic redshift. To investigate this, we made a catalogue of all sources with a photometric redshift within a radius of $10^{\prime}$ around Nube. This radius corresponds to the area where the response efficiency of the GBT beam is greater than 5\% (see Fig. 3). We used one of the value-added catalogues from the DR9 DECaLS survey, which provides photometric redshifts of sources down to a $5 \sigma$ depth of $m_{r}=24 \mathrm{mag}$ (Zhou et al. 2021). The list of objects with a photometric redshift compatible with the HI emission (i.e. $z=0.02502$ ) is given in Table C.1. The vast majority of objects with a photometric redshift compatible with the HI detection reported in this work are either stars or pointlike sources. This is to be expected, because at a given magnitude, photometric catalogues are biased towards objects with higher signal-to-noise ratio, and therefore extended objects are either less represented or missing.

Finally, to account for the possibility that some faint and diffuse galaxies might not have a photometric redshift measurement, we visually inspected the inner $10^{\prime}$ around Nube using the Dark Energy Camera Legacy Survey (DECaLS, Dey et al. 2019) DR10 images to find extended faint sources that could potentially be emitting in HI. We found only two very faint galaxies, which are listed in Table C.2. These two galaxies are visually smaller than Nube and therefore potentially less massive at the same distance. Furthermore, these objects are found at a radial distance of $\gtrsim 4.5^{\prime}$ from the central direction of the HI beam. While we cannot completely rule out the possibility that one of these galaxies is responsible for the HI detection, their off-centre location would imply that they are less likely to be associated with the HI detection than Nube.

The analysis in this section reinforces the idea that the detection of an emission peak in HI could be related to the optical detection of Nube. Moreover, the fact that the HI emission observed coincides with the velocity of UGC 929 (even if it is extremely unlikely to be caused by it) is a further argument in favour of the HI detection being genuine and not a fluctuation. Having ruled out contamination of nearby HI sources, there is still the possibility that the optical and HI detections are unrelated, and therefore the distance assumed for the optical counterpart of the galaxy may be incorrect. However, this possibility does come with its own problems. We refer to Appendix B for a detailed discussion of the implications of the distance for Nube. The scenario that follows from the analysis in this section is that Nube could be a satellite of UGC 929 at a distance of 107 Mpc .

## 3. Analysis and results

### 3.1. Radial profiles of Nube

The aim of this paper is to study the properties of Nube in detail in order to determine its origin. A fundamental way to characterise the shape of a galaxy is through its radial surface


Fig. 3. Galaxies found in NED with known spectroscopic redshifts in a range of $100^{\prime} \times 100^{\prime}$ around Nube and a recession velocity of within $2000 \mathrm{~km} \mathrm{~s}^{-1}$ of that of the HI detection (i.e. $V_{\text {helio }}=7480 \mathrm{~km} \mathrm{~s}^{-1}$ ). The galaxies are colour-coded according to their relative velocity with respect to the potential velocity of Nube. The grey-scale background image is the GBT beam response from Spekkens et al. (2013). The contours indicate where the efficiency drops by the factor indicated on the label, from $0.05 \%$ (black) to $50 \%$ (light blue). The top horizontal axis indicates the equivalent size in Mpc at a distance of 107 Mpc .
brightness profiles. In addition, multi-wavelength information provides valuable constraints on galaxy formation processes. We therefore derived the surface brightness radial profiles of this galaxy using the optical, multi-wavelength HiPERCAM data.

To measure the photometry of this galaxy, we first masked out all foreground and background sources in the image. As can be seen in the insets in Fig. 1, the galaxy is very faint and almost transparent. As we cannot determine whether the clumps on top of the galaxy are part of the galaxy or are objects in the line of sight of the galaxy, we masked them all, leaving only the diffuse component for photometry. We used a combination of NoiseChisel (Akhlaghi \& Ichikawa 2015; Akhlaghi 2019) and manual masking after visual inspection. This is a very conservative approach as we are only analysing the diffuse light of the galaxy. Appendix D shows the mask used in this work.

Once we had masked out all the sources of contamination, we derived the radial surface brightness profiles of the galaxy. Given the nearly circular shape of the object, we decided to extract the radial profiles using circular annuli at different radial distances, up to $30^{\prime \prime}$ from the centre of the galaxy. To derive the profiles, we used a custom python code. For each radial bin, the surface brightness was obtained as the $3 \sigma$-clipped median of the pixel values. Figure 4 shows the surface brightness profiles of

Nube for the five bands imaged with HiPERCAM. Some of the profiles are shifted by a constant value - given in the legend for visualisation purposes. The errors are calculated as a combination of the Poisson noise in each annulus and the error in the sky given by the distribution of background pixels in each image. These surface brightness profiles are corrected for the absorption of our Galaxy $(E(B-V)=0.34$, Schlafly \& Finkbeiner 2011).

We can now study the radial variations of the galaxy's stellar populations. The surface brightness profiles are used to determine the radial $g-r$ profile of the galaxy. The top panel in Fig. 5 shows the $g-r$ colour profile of Nube. The colour profile appears flat at all radii, with an average colour of $\sim 0.6$. This almost flat colour profile could be compatible with a similar age and metallicity at all radii, up to $10^{\prime \prime}(15 \mathrm{kpc})$.

The bottom panel of Fig. 5 shows the surface stellar mass density profile for this galaxy. To derive this latter, we follow the procedure given in Bakos et al. (2008) to link the observed surface brightness in the $g$-band to the radial variation of the stellar mass to light $(\mathrm{M} / \mathrm{L})$ ratio. The $\mathrm{M} / \mathrm{L}$ ratio was obtained from the prescriptions given in Roediger \& Courteau (2015) assuming a Chabrier (2003) IMF, using the $g-r$ colour profile. As with the colour profile, the stellar mass density profile of Nube is also relatively flat compared to the stellar mass density profiles


Fig. 4. Radial surface brightness profiles of Nube in the Sloan $u, g, r, i$, and $z$ bands. The profiles in the $i$ and $z$ bands have been shifted vertically for visualisation purposes.
of galaxies with similar stellar masses (e.g. Montes et al. 2021). We expand on this result later in the text.

We used this stellar mass density profile to measure the radius that encloses half the mass of the galaxy, or the half-mass radius ( $R_{\mathrm{e}}$ ). This radius is $R_{\mathrm{e}}=13.7 \pm 1.7 \operatorname{arcsec}(6.9 \pm 0.8 \mathrm{kpc}$ at a distance of 107 Mpc ). The effective surface stellar mass density of Nube is $\langle\Sigma\rangle_{\mathrm{e}}=0.9 \pm 0.1 M_{\odot} \mathrm{pc}^{-2}$. The total stellar mass of the galaxy derived from the surface mass density profile, assuming circular symmetry, is $M_{*}=3.9 \pm 1.0 \times 10^{8} M_{\odot}$. The effective surface brightness is $\left\langle\mu_{V}\right\rangle_{\mathrm{e}}=26.75 \pm 0.02 \mathrm{mag}_{\mathrm{arcsec}}{ }^{-2}$. These values, along with other global properties of Nube, are listed in the Table 1.

To assess the validity of our results, we also fitted Sérsic (1968) models (using the code GALFIT, Peng et al. 2002) to the galaxy in all the different HiPERCAM bands, independently. The values of the half-light effective radius, $r_{\mathrm{e}}$, and the Sérsic index for each band are given in Table 2. The values of $r_{\mathrm{e}}$ from the GALFIT fits are consistent with those derived using the surface brightness profiles.

### 3.2. Stellar populations of Nube

Given the very low surface brightness of this galaxy $\left(\mu_{V}(0) \sim\right.$ $26.2 \mathrm{mag}_{\operatorname{arcsec}}{ }^{-2}$ ), the use of deep multi-wavelength observations is the best, if not the only, way to constrain its stellar populations, namely the age, metallicity, and M/L. To estimate these quantities, we first constructed the spectral energy distribution (SED) of the galaxy (see Fig. 6). The photometry at each band was derived measuring the total flux within a circular aperture with radius equal to the half-mass radius $R_{\mathrm{e}}$ (13.7") of the galaxy. The purpose of using this aperture (instead of the total galaxy) is to ensure enough signal to reliably characterise the SED of Nube. The errors of the photometry of each individual band that compose the SED are a combination of the photometric errors and the zero point uncertainties.

To characterise the stellar population properties of the galaxy, we fitted Bruzual \& Charlot (2003) single stellar population (SSP, instantaneous burst) models to the SED of the galaxy. This is a reasonable assumption given the homogeneous colour distribution of this galaxy (Fig. 5). The parameters to be fitted


Fig. 5. Colour and stellar surface density profile of Nube. Upper panel: HiPERCAM $g-r$ colour profile of Nube. Bottom panel: stellar surface mass density profile of Nube. The upper horizontal axis indicates the equivalent size in kiloparsecs at a distance of 107 Mpc .
are age, metallicity, and luminosity, and we assumed a Chabrier IMF (Chabrier 2003) for the models. As the SED is derived from broadband imaging, it is a good approach to assume an SSP to describe the stellar populations of this galaxy, because the information we have is limited. This assumption gives us the average properties of this galaxy. For the fit, we use the $\chi^{2}$ minimisation approach described in Montes et al. (2014). We derive a most likely age of $10.2_{-2.5}^{+2.0} \mathrm{Gyr}$ and a metallicity $[\mathrm{Fe} / \mathrm{H}]$ of $-1.09_{-0.13}^{+0.09}$; these values are in agreement with those of old diffuse low-mass galaxies (see e.g. Ruiz-Lara et al. 2018; Heesters et al. 2023; Iodice et al. 2023). Uncertainties in the parameters were estimated by marginalising the 1D probability distribution functions obtained during the fitting. As the colour profile of this galaxy is flat at all radii (see Fig. 5), we assume that the stellar populations obtained for the inner $R_{\mathrm{e}}$ are representative of the whole galaxy.

Using the $\mathrm{M} / \mathrm{L}$ derived from the SED fit $\left((\mathrm{M} / \mathrm{L})_{V}=1.8_{-0.5}^{+0.3}\right)^{5}$, the total stellar mass of Nube is $M_{*}=4.4 \pm 0.8 \times 10^{8} M_{\odot}$, which is in agreement with the stellar mass estimate derived in Sect. 3.1. As this stellar mass is derived with the properties inside $R_{\mathrm{e}}$, we prefer to use the previous estimate (Sect. 3.1), as it is more representative of the whole galaxy.

[^5]Table 1. Global properties of Nube.

| RA | Dec | $z$ | Distance <br> $(\mathrm{Mpc})$ | $\log \left(M_{\mathrm{HI}} / M_{\odot}\right)$ | $\log \left(M_{*} / M_{\odot}\right)$ | $R_{\mathrm{e}}$ <br> $(\operatorname{arcsec})$ |
| :--- | :---: | :---: | :---: | :---: | :---: | :---: |
| $01^{\mathrm{h}} 23^{\mathrm{m}} 27.37^{\mathrm{s}}$ | $-00^{\mathrm{d}} 37^{\prime} 27.83^{\prime \prime}$ | $0.02502 \pm 0.00003$ | $107 \pm 5$ | $8.35 \pm 0.12$ | $8.6 \pm 0.1$ | $13.7 \pm 1.7$ |
| $\left(\mathrm{mag} \mathrm{arcsec}^{-2}\right)$ |  |  |  |  |  |  |

Notes. The distance-dependent parameters were calculated assuming that the object is at a distance of 107 Mpc .


Fig. 6. Spectral energy distribution (filled circles) of Nube derived from HiPERCAM photometry within the half-mass radius. The best-fit single stellar population model is shown as the blue line and the best-fit magnitudes (model convolved with the filter response) are the open circles. The age, metallicity, and $\mathrm{M} / \mathrm{L}$ corresponding to the best-fit model are given in the figure.

### 3.3. Dynamical mass of Nube

A rough estimate of the dynamical mass of this galaxy ${ }^{6}$ can be obtained using the following equation from Spekkens \& Karunakaran (2018):
$M_{\mathrm{dyn}}^{3 R_{\mathrm{e}}}=6.96 \times 10^{5} R_{\mathrm{e}}\left(\frac{W_{50}}{2 \times \sin (i)}\right)^{2}\left[M_{\odot}\right]$,
where $R_{\mathrm{e}}$ is the half-mass radius in $\mathrm{kpc}, i$ the inclination of the HI disc in degrees, and $W_{50}$ the turbulence-corrected velocity width of the HI line. We calculate the inclination of the HI disc as follows:
$\cos (i)=(b / a)$.
In Table 1, we provide the axis ratio, $b / a$, of Nube obtained from GALFIT Sérsic fits to the object (see Sect. 3.1). The resulting value $(0.97 \pm 0.01)$ is the average from the axis ratio measured in the $g, r$, and $i$ bands. This results in an inclination value of $13 \pm 3 \mathrm{deg}$. It is worth emphasising that using the axis ratio

[^6]of the stellar distribution to characterise the inclination of the HI disc is a rough approximation. The dynamical mass we derive for Nube is $2.6 \pm 1.7 \times 10^{10} M_{\odot}$ within $3 R_{\mathrm{e}}=20.7 \mathrm{kpc}$.

### 3.4. A search for globular clusters

Globular clusters (GCs) provide a complementary way to estimate the distance to an extragalactic object based on the universality of their luminosity function. The peak of the GC distribution lies at $M_{V}=-7.6 \mathrm{mag}$ for different galaxies (see e.g. Rejkuba 2012), making it a secondary distance indicator. Therefore, the detection of GCs in Nube would help us to validate the distance obtained with the HI detection. We note that the peak of the GC luminosity at the distance of 107 Mpc (i.e. a distance modulus of 35.15 mag ) is expected at $m_{g} \simeq 27.5 \mathrm{mag}$, which is fainter than our point-like $5 \sigma$ limiting magnitudes, meaning that, if there were any, it would be difficult to detect GCs in our images (see Sect. 2.1). We would expect an ultradiffuse galaxy (UDG) of similar stellar mass to Nube to have around about 10 GCs, although the range varies from 0 to 30 GCs. However, only the brightest GCs of the system are expected to be detectable.

We follow a similar procedure to detect GCs as in Montes et al. (2020, 2021). Briefly, we first preselect the GCs based on their morphology (size and ellipticity), and then we refine the selection based on their colours. However, in this case, we do not have the high spatial resolution of the HST images to make a preselection based on the morphology of the GCs, but only the ground-based images. For this reason, we have to impose very strict shape criteria in order to minimise contamination in our selection.

We ran SExtractor on our images in dual mode, with the $r$-band image as the detection image. At a distance of 107 Mpc , GCs will resemble point sources in the images. Therefore, to preselect the GC candidates, we first select stars in our images by imposing that the stellarity parameter in SExtractor ("CLASS_STAR") is greater than 0.98 . We do this to estimate the FWHM of a point-like source in our images, to impose strict conditions on our detections, and to minimise contamination from other objects. We find six point-like objects, three of which were detected in Gaia DR3 (Gaia Collaboration 2023). The FWHMs of these point-like objects are between 10 and 14 pix $\left(0.8^{\prime \prime}\right.$ to $\left.1.12^{\prime \prime}\right)$ and their ellipticities are $<0.1$. Therefore, based on these values, we preselected the GC candidates with an ellipticity of $<0.1,9<F W H M<15$ pix, and $m_{r}<26.2$ mag in order to minimise false detections.

After this initial selection, we also applied the $(u-r)-(r-z)$ colour-colour selection used in Taylor et al. (2017) to narrow down the selection of $\mathrm{GCs}^{7}$. To do this, we defined a colour-

[^7]

Fig. 7. $(u-r)-(r-z)$ colour-colour diagram of the initial sample of GC candidates (blue) of Nube. The green squares are the location of NGC 5128 candidates in Taylor et al. (2017) with a probability of greater than $95 \%$ of being true GCs. The black dashed line indicates the convex hull (envelope) computed as the smallest region containing all the green squares.
colour region based on the position on the diagram of candidate GCs with a probability of $>0.95$ of being true GCs in Taylor et al. (2017). To define this region, we computed the convex hull or convex envelope, that is, the minimum region containing all GCs from Taylor et al. (2017, dashed line in Fig. 7).

Only two sources fall on the area defined by the GCs provided by Taylor et al. (2017) ${ }^{8}$. However, both are very far from Nube ( $>110^{\prime \prime}$ or 54 kpc from its centre), and are therefore unlikely to be associated with the galaxy. We also checked the three sources near the bottom of the region $(u-r \sim 2$, $r-z \sim 0.25$ ), but they are also not spatially associated with Nube. In summary, we are unable to detect any GCs associated with Nube, as expected considering both our point-like limiting magnitudes and the distance to the galaxy of 107 Mpc .

## 4. On the formation mechanism of Nube

### 4.1. Nube in comparison to other low surface brightness galaxies

The galaxy reported in this paper, Nube, has some extreme properties in terms of size and surface brightness. Figure 8 shows the structural properties of Nube compared to other low surface brightness or dwarf galaxies compiled from the literature. The grey open circles represent low-mass Local Group galaxies from McConnachie (2012), the green open squares are field dwarfs, the orange crosses are satellite dwarfs (orbiting a more massive

[^8]galaxy) from Carlsten et al. (2021), the pink " x " symbols mark UDGs galaxies in Coma from van Dokkum et al. (2015), the golden star is UDG32 from Iodice et al. (2021), and the red diamonds are two of the low surface brightness (LSB) galaxies in Virgo from Mihos et al. (2015). Compared to all these objects, Nube appears quite unique given its low central surface brightness and extension for a galaxy of its stellar mass. The closest objects in size, mass, and central surface brightness are the two LSB galaxies in Virgo from Mihos et al. (2015), VLSB-B and VLSB-C ${ }^{9}$, and UDG32 in Hydra I (Iodice et al. 2021). We do not include VLSB-A, as this LSB galaxy is clearly tidally perturbed, as indicated by its visible tidal tails. VLSB-C appears to be quite similar in size and brightness to Nube ( $R_{\mathrm{e}}=5.5 \mathrm{kpc}$ and $\mu_{V}(0)=26.7 \mathrm{mag}_{\mathrm{arcsec}}{ }^{-2}$, Mihos et al. 2015). This galaxy shows no obvious signs of tidal stripping; it is possibly an LSB galaxy in the cluster outskirts or falling into the Virgo cluster for the first time.

We also compare Nube with other galaxies that have been described as "almost dark" in previous works. One example is AGC229385 (also known as Coma P.; Janowiecki et al. 2015; Brunker et al. 2019), which, unlike Nube, appears elongated. It also has a more extreme ratio of HI to stellar mass $\left(M_{\mathrm{HI}} / M_{*}=\right.$ 81; Brunker et al. 2019) compared to Nube ( $M_{\mathrm{HI}} / M_{*} \sim 1$ ). This is also the case for AGC229101 (Leisman et al. 2021). This latter galaxy is an HI source ( $M_{\mathrm{HI}} / M_{*}=98$ ) with a very dim optical counterpart ( $\mu_{g}(0)=26.6 \mathrm{mag}_{\operatorname{arcsec}}{ }^{-2}$ ) and $R_{\mathrm{e}}=3 \mathrm{kpc}$. A further two extended "almost dark" galaxies are AGC229398 and AGC333576 (Gray et al. 2023). Their effective radii are large, 9 and 4.4 kpc respectively, and their stellar masses ( $\sim 10^{8} M_{\odot}$ ) are also close to that of Nube. However, their central surface brightnesses are significantly brighter than that of Nube (23.8 and $24.6 \mathrm{mag} \operatorname{~arcsec}^{-2}$ ).

In addition to comparing the global properties of Nube with other low-mass galaxies, it is also very instructive to compare its surface mass density profile with respect to galaxies of similar stellar mass. In Fig. 9, we show the stellar surface mass density of Nube (blue dots) derived from the $g-r$ colours (Sect. 3.1). For comparison, we plot the profiles of dwarf galaxies of similar stellar mass (grey lines) in the mass range of $1-5 \times 10^{8} M_{\odot}$ from Chamba et al. (2020). This shows how different Nube is from typical dwarf galaxies of similar mass. Even UDGs with large effective radius are not comparable to this galaxy. For example, the iconic UDG Dragonfly 44 (DF44), also shown as the green squares in Fig. 9, has a $r_{\mathrm{e}}=4.3 \pm 0.2 \mathrm{kpc}$ (van Dokkum et al. 2016) ${ }^{10}$, and a central surface brightness of $\mu_{g}(0)=24.5 \mathrm{mag} \mathrm{arcsec}^{-2}$. This is 1.5 times smaller and $1.4 \mathrm{mag} \mathrm{arcsec}^{-2}$ brighter than Nube.

### 4.2. Exploring the nature of Nube

Given the extreme properties of Nube, it is interesting to discuss whether these properties are a result of the original formation of this galaxy or are due to a later evolutionary process caused by the environment in which it is found. In this section, we discuss two possible alternatives. In the first, we explore whether or not Nube can be considered a tidal dwarf galaxy. The second alternative is that Nube was born with a stellar density typical of dwarf galaxies of the same stellar mass and has been deformed

[^9]

Fig. 8. Comparison of the structural properties of Nube with those of other Local Group galaxies from McConnachie (2012, grey open circles), dwarf galaxies in Carlsten et al. (2021, green open squares and orange crosses), UDGs in Virgo from Mihos et al. (2015, red diamonds), UDG32 in Hydra I from Iodice et al. (2021, golden star), and Coma from van Dokkum et al. (2015, pink x). Although the central surface brightness is typical of other dwarf galaxies, the stellar mass and half-mass radius make Nube, at a distance of 107 Mpc , an extreme object, even more so than the UDGs in Virgo and Coma.
by the environment into this peculiar structure. Either of these two possibilities (if correct) should be able to explain the stellar populations of Nube and its present morphology.

### 4.2.1. Is Nube a tidal dwarf galaxy?

There is a population of galaxies whose formation mechanisms would not be associated with what is expected to be the main channel of galaxy formation, namely gravitational collapse of gas within a dark matter halo. These galaxies are known as tidal dwarf galaxies (TDGs). They form from material torn away from larger galaxies by tidal interactions or by harassment (e.g. Duc \& Bournaud 2008) and, as a result of their formation, they exhibit characteristics such as significant gas fractions, low dark matter content, and higher-than-expected metallicity for dwarf galaxies of similar mass (e.g. Duc 2012).

Around $95 \%$ of the known TDG galaxies are located at relatively modest distances ( $\sim 20 \mathrm{kpc}$ ) from their potential progenitors (Kaviraj et al. 2012). However, Nube is very distant $(435 \mathrm{kpc})$ from its likely progenitor UGC 929 . This large distance from its possible parent should not be considered as a factor that counts against the characterisation of Nube as a TDG,
because the selection of TDG galaxies usually requires a visible tidal tail and is therefore biased towards newly formed objects. Some simulations suggest that if the TDG is massive enough ( $M>10^{8} M_{\odot}$ ) and escapes the parent galaxy with sufficient velocity, it could become self-gravitating and avoid falling back into its progenitor system (Bournaud \& Duc 2006). However, it is unclear how long TDGs formed in this way could survive without the protection of a massive dark matter halo, as both axisymmetric and non-axisymmetric instabilities would eventually destroy the system (Sellwood \& Sanders 2022). Assuming that these galaxies manage to survive, they are expected to remain gas-rich and very faint. These old galaxies are therefore good candidates to be identified as almost dark galaxies (see e.g. Cannon et al. 2015; Janowiecki et al. 2015; Leisman et al. 2017; Román et al. 2021).

The metallicity observed in Nube $\left([\mathrm{Fe} / \mathrm{H}]=-1.09_{-0.13}^{+0.09}\right)$ is similar to that of normal dwarf galaxies (e.g. McConnachie 2012). However, it is lower than expected for currently forming TDGs (solar metallicity; e.g. Duc et al. 2000), which are expected to be more metallic as they inherit the metallicity of the more massive parent galaxies. Nonetheless, given the age of the galaxy ( 9 Gyr ), it could have formed from ancient spirals with


Fig. 9. Surface stellar mass density profile of Nube (blue dots) compared with other galaxies of similar stellar mass. The profiles of dwarf galaxies from Chamba et al. (2020) in the same mass range ( $1-5 \times 10^{8} M_{\odot}$ ) are plotted. We also show the profile of DF44 (green squares), as it is an iconic large ultra-diffuse galaxy. The vertical lines indicate the halfmass radius ( $R_{\mathrm{e}}$ ) for Nube (blue solid line) and DF44 (green dashed line).
less enriched gas than the present ones (see e.g. Recchi et al. 2015). Therefore, the relatively low metallicity of Nube should not necessarily be an indication that it did not form as a TDG.

Another characteristic of TDGs is their low dark matter content. The dynamical mass of Nube is $2.6 \pm 1.7 \times 10^{10} M_{\odot}$ within $3 R_{\mathrm{e}}=20.7 \mathrm{kpc}$ (Sect. 3.3), while the stellar mass is $M_{*}=3.9 \pm 1.0 \times 10^{8} M_{\odot}$ (Sect. 3.2). This means that the ratio of dark matter to stellar mass is between 20 and 150 . This amount of dark matter is significantly greater than the amount of dark matter expected if the galaxy formed as a TDG $\left(M_{\mathrm{dyn}} / M_{*}<2\right.$; Gray et al. 2023). In addition, its gas content relative to the stellar content is not very large compared to those found in other TDGs ( $M_{\mathrm{HI}} / M_{*} \sim 10$, Gray et al. 2023). For these reasons, we conclude, with high confidence, that Nube is not a TDG.

### 4.2.2. Nube as a result of environmental processes

If Nube is not a TDG, then it would be interesting to know whether or not its structural properties can be explained as the result of interaction with its environment. It is known that dwarf galaxies can be dynamically heated and "puffed up" by interactions with more massive galaxies (e.g. Liao et al. 2019; Tremmel et al. 2020). This puffing could explain why Nube is similar to other regular dwarf galaxies in terms of age and metallicity, while its physical properties (effective radius and central surface brightness) are extreme. However, against this scenario for the origin of Nube is the fact that the galaxy does not inhabit a particularly dense environment. In fact, its nearest massive neighbour, UGC 929, is at a distance of 435 kpc in projection.

Assuming that Nube and UGC 929 underwent a close encounter in the past, we can make a rough estimate of when this occurred. Considering the current projected distance between the two objects plus a transverse projected velocity of $\sim 100 \mathrm{~km} \mathrm{~s}^{-1}$


Fig. 10. Region of $100^{\prime \prime} \times 100^{\prime \prime}$ around Nube. The figure is a composite of an RGB colour image using the $g, r$, and $i$ HiPERCAM bands and a black and white $g+r$ image for the background.
(i.e. $100 \mathrm{kpc} \mathrm{Gyr}^{-1}$ ), the close encounter would have occurred $4-5 \mathrm{Gyr}$ ago. Given the dynamical time of each galaxy, this is enough time to have erased any signature of the perturbation in the central regions of both objects. We do not find the central region of Nube to be obviously perturbed (see Fig. 10). Regarding UGC 929, we used deep imaging from DECaLS to study its morphology. Figure A. 1 shows the composite DECaLS image of the galaxy. We can see in the image that the galaxy appears to be very symmetric, with no sign of perturbations or interactions. Therefore, if Nube and UGC 929 had met in the past, it would have been long enough ago for the central shape of both galaxies to have already been restored.

However, in the outermost parts of Nube it would be possible to study the effects of this gravitational interaction in the form of tidal distortions. These deformations produce an S-shaped structure (e.g. Johnston et al. 2002; Moreno et al. 2022). Thanks to the deep imaging provided by HiPERCAM, we can investigate whether there is any evidence that this could be the case for Nube. An interaction capable of distorting the galaxy in this way would be clearly visible in our ultradeep imaging ( $\sim 30.5 \mathrm{mag} \mathrm{arcsec}^{-2}$; see Sect. 2.1) even long after the interaction (Moreno et al. 2022). Figure 10 shows an RGB composite of the image of Nube with a black and white background to highlight the fainter outskirts of this galaxy. Upon visual inspection of Nube, we see no obvious signatures of deformation in the outer regions. This can be seen from the lack of excess in the outer parts of the surface brightness and stellar density profiles (Figs. 4 and 5), which is associated with the presence of tidal tails (Johnston et al. 2002; Montes et al. 2020).

To improve the visualisation of the outer parts of Nube, we also modelled Nube with GALFIT. To obtain additional signal-to-noise ratio, we performed a model fit on a $g+r$ image. Figure 11 shows the original (left), model (middle), and residual (right) $g+r$ image. No signs of an excess or deformation can be seen.


Fig. 11. GALFIT model of Nube for the HiPERCAM $g+r$-band. The left panel shows the $g+r$ original image, the middle panel the 2D Sérsic model fit, and the right panel the residuals after subtracting the Sérsic model, after smoothing it with a Gaussian kernel to enhance the residuals.

Table 2. Multi-band effective radii, axis ration, and Sérsic indexes of the GALFIT (Peng et al. 2002) fits to Nube.

| Filter | $r_{\mathrm{e}}$ <br> $(\operatorname{arcsec})$ | $n$ | $b / a$ |
| :--- | :---: | :---: | :---: |
| $u$ | $15.6 \pm 0.6$ | $0.60 \pm 0.03$ | 0.999 |
| $g$ | $15.9 \pm 0.1$ | $0.70 \pm 0.01$ | 0.986 |
| $r$ | $15.3 \pm 0.1$ | $0.65 \pm 0.01$ | 0.962 |
| $i$ | $13.9 \pm 0.1$ | $0.60 \pm 0.01$ | 0.975 |
| $z$ | $16.0 \pm 0.3$ | $0.75 \pm 0.02$ | 0.971 |

## 5. Nube as a test bed to explore the nature of dark matter

### 5.1. Nube within the cold dark matter scenario

Given the extreme properties of Nube, it is interesting to see whether traditional models of cold dark matter are able to reproduce galaxies with these characteristics. In particular, we are interested in knowing if galaxy formation models can produce objects with stellar masses, surface brightnesses, and effective radii like Nube. To answer this question, it is worth looking at simulations that have been able to reproduce the properties of the largest known UDGs. These simulations have sufficient spatial resolution and stellar population feedback recipes to produce galaxies with low surface brightness and large effective radius, such as DF44.

Using the FIRE-2 simulations (Hopkins et al. 2018), Chan et al. (2018) find that isolated field dwarfs, where the effect of stellar feedback (i.e. stellar winds, supernovae, etc.) is large, produce galaxies with surface brightness, effective radius, and stellar masses representative of UDGs (after imposing artificial quenching to simulate the effect of infall into the cluster environment). Simulated galaxies with stellar masses of around $10^{8} M_{\odot}$ and dark matter halos of around $10^{10} M_{\odot}$, such as Nube, have effective radii of less than 5 kpc , which is smaller than what we measure for Nube. NIHAO simulations (Wang et al. 2015) produce galaxies with similar structural properties to UDGs because of episodes of gas outflow associated with star formation (Di Cintio et al. 2017). The dark matter halo masses and stellar masses are consistent with what we find for Nube. However, the effective radii ( $\sim 3 \mathrm{kpc}$ ) of these simulated galaxies are again well below the value we measure for Nube.

Other proposed mechanisms leading to the formation of UDGs include high-z major mergers (Wright et al. 2021)
and a scenario where UDGs populate higher spin halos (Amorisco \& Loeb 2016; Benavides et al. 2023). However, neither of these scenarios can reproduce the observed characteristics of Nube. In Wright et al. (2021), major mergers produce galaxies with $r_{\mathrm{e}}<4 \mathrm{kpc}$. Benavides et al. (2023) explore UDGs in the TNG50 simulation and find that for isolated galaxies of the mass of Nube, the effective radius is $\lesssim 5 \mathrm{kpc}$.

### 5.2. Nube in the fuzzy dark matter framework

As the microphysical nature of the dark matter is still completely unknown, it is worth exploring whether or not extreme objects such as Nube could be compatible with other alternative dark matter models. In particular, models based on ultralight (axionlike) scalar particles have received significant interest (fuzzy dark matter; Schive et al. 2014; May \& Springel 2021). Due to the very small mass of these particles $\left(\sim 10^{-22} \mathrm{eV}\right)$, the quantum effects are expected to appear at the kiloparsec (i.e. galactic) scale. The ultralight particles generate dark matter models with large cores. In particular, the dark matter distribution is expected to generate a central distribution called a soliton, followed by a NFW profile in the outer parts (Schive et al. 2014). The soliton density profile can be well approximated as
$\rho_{\mathrm{s}}(r)=\frac{1.9\left(m_{B} / 10^{-23} \mathrm{eV}\right)^{-2}\left(r_{\mathrm{c}} / \mathrm{kpc}\right)^{-4}}{\left[1+9.1 \times 10^{-2}\left(r / r_{\mathrm{c}}\right)^{2}\right]^{8}} M_{\odot} \mathrm{pc}^{-3}$,
where $m_{B}$ is the mass of the dark matter particle and $r_{\mathrm{c}}$ the core radius where the density has dropped to one-half of its peak value. Equation (29) of Bar et al. (2018) shows that if the core radius and the virial mass of the dark matter halo $M_{\mathrm{h}}$ are known, the dark matter particle can be estimated through the following equation:
$\frac{m_{B}}{10^{-22} \mathrm{eV}}=160\left(\frac{r_{\mathrm{c}}}{\mathrm{pc}}\right)^{-1}\left(\frac{M_{\mathrm{h}}}{10^{12} M_{\odot}}\right)^{-1 / 3}$.
To get a rough estimate of the core radius, we assume the stellar distribution follows the shape of the dark matter soliton. This is a reasonable assumption, as the dark matter is the dominant component generating the global gravitational potential (Sánchez Almeida et al. 2023). Therefore, we project the soliton density profile, obtaining
$\Sigma_{\star}(R)=\frac{\Sigma_{\star}(0)}{\left(1+0.091 \times\left(R / r_{\mathrm{c}}\right)^{2}\right)^{15 / 2}}$.
The stellar mass surface density of Nube fits very well with this equation (see Fig. 12). We obtain a value of $r_{\mathrm{c}}=6.6 \pm 0.4 \mathrm{kpc}$.


Fig. 12. Surface stellar mass density profile of Nube (blue dots) fitted with a projected soliton shape with a core radius of $r_{\mathrm{c}}=6.6 \pm 0.4 \mathrm{kpc}$ (see text for details). The soliton best fit and its errors are indicated by the red line and the red region. The agreement is remarkable when taking into account that the only free parameter is essentially the core radius. The black dashed line is a projected NFW profile for comparison.

The mass of the dark matter particle compatible with this core radius is $m_{B}=\left(0.8_{-0.2}^{+0.4}\right) \times 10^{-23} \mathrm{eV}$.

We also plot a Navarro, Frenk, and White (NFW) profile (dashed black line, Navarro et al. 1996) to show that this commonly used profile to describe the distribution of mass in dark matter halos cannot reproduce the characteristics of Nube, especially at small radii. For Fig. 12, we used the projected NFW profile in Eq. (11) in Wright \& Brainerd (2000), assuming the core radius of the soliton derived above as transition radius ( $r_{\mathrm{s}}$ ).

It is interesting to compare the mass of the axion-like particle derived in this analysis with other studies. Using galactic rotation curves, Bernal et al. (2018) found an average value of $m_{B} \sim 0.5 \times 10^{-23} \mathrm{eV}$, while Bañares-Hernández et al. (2023) derived $m_{B} \sim 2 \times 10^{-23} \mathrm{eV}$. These values are in good agreement with our estimate using only the structural properties of the galaxy. Analysing the velocity dispersion of the Fornax and Sculptor dwarfs, González-Morales et al. (2017) find an upper limit of $m_{B}<4 \times 10^{-23} \mathrm{eV}$. Based on the angular scale of the CMB acoustic peaks and anisotropies, Hložek et al. (2018) derive a lower limit of $m_{B}>0.1 \times 10^{-23} \mathrm{eV}$. These results are consistent with our estimate for the mass of the axion-like particle.

However, the situation is far from clear. Other estimates of the mass of the particle are inconsistent with the above results. To name a few, Chen et al. (2017), using a Jeans analysis of the Milky Way satellite dSphs, find $m_{B}=(17.9 \pm 3.3) \times 10^{-23} \mathrm{eV}$. Dalal \& Kravtsov (2022) suggest that the size and stellar kinematics of ultrafaint dwarfs imply a lower limit of $m_{B}>3 \times$ $10^{-19} \mathrm{eV}$. Analysis of the stellar heating of the Milky Way disc due to the substructure within a fuzzy dark matter halo suggests a lower limit of $m_{B} \geq 4 \times 10^{-23} \mathrm{eV}$ (Chiang et al. 2023). Although fuzzy dark matter could relieve some of the small-scale tensions appearing in the cold dark matter scenario, more work is needed to assess this model.

## 6. Conclusions

In the absence of a laboratory detection of the particles that make up dark matter, the hope for characterising its nature lies in the analysis of astrophysical objects. In particular, galaxies with very low surface brightness (also called "almost dark") offer an interesting possibility to constrain the microphysical properties of dark matter. In these galaxies, baryonic feedback effects are expected to be very moderate, leaving the spatial distribution of their dark matter halo almost intact.

In this paper, we present the serendipitous discovery in the IAC Stripe 82 Legacy Project (Fliri \& Trujillo 2016; Román \& Trujillo 2018) of a very extended $\left(R_{\mathrm{e}}=6.9 \pm\right.$ $0.8 \mathrm{kpc})$ and very low surface brightness $\left(\left\langle\mu_{V}\right\rangle_{\mathrm{e}}=26.75 \pm\right.$ $0.02 \mathrm{mag} \mathrm{arcsec}^{-2}$ ) galaxy of $M_{*}=3.9 \pm 1.0 \times 10^{8} M_{\odot}$ in stellar mass located at 107 Mpc . We call this galaxy Nube. Using the GBT radio telescope, we infer a total halo mass of $2.6 \pm 1.7 \times$ $10^{10} M_{\odot}$ for the galaxy. Current simulations of UDG formation - which take into account baryonic feedback effects and simulate dark matter particles as WIMPs (i.e. cold dark matter) are unable to reproduce objects with the properties of Nube. We investigated the possibility that the object could be reproduced with the predictions for the fuzzy dark matter model. Under the hypothesis that the distribution of stars in Nube is representative of the distribution of the dark matter halo, we find that a solitonshaped profile (typical of fuzzy dark matter) reproduces the observed distribution of stars very well. The mass of the axionlike particle inferred from the fit is $m_{B}=\left(0.8_{-0.2}^{+0.4}\right) \times 10^{-23} \mathrm{eV}$. This value is in good agreement with other astrophysical measurements using the dynamical properties of other low surface brightness galaxies.

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## Appendix A: Properties of UGC929



Fig. A.1. Region of $250^{\prime \prime} \times 250^{\prime \prime}(\sim 123 \mathrm{kpc} \times 123 \mathrm{kpc})$ around the galaxy UGC929. The figure is a composite of an RGB image using the $g, r$, and $z$ from DECaLS and a black and white $g+r+z$ image for the background. This galaxy is at the same redshift as Nube, but is found 14.7 ' to the northwest. The blue arrow indicates the direction of Nube. There is no evidence of morphological disturbance that could point to gas expulsion as the origin of Nube.

Given the difference in distance between both galaxies, we assume that Nube is a satellite of UGC929. UGC929 is at a projected distance of $14.7^{\prime}$ ( 435 kpc at a distance of 107 Mpc ) from Nube. To calculate the stellar mass of UGC929, we use the DECaLS $g$ and $r$ images, obtaining the $\mathrm{M} / \mathrm{L}$ ratio using
the Roediger \& Courteau (2015) prescriptions and assuming a Chabrier (2003). The stellar mass of the galaxy is $M_{*}=2.4 \times$ $10^{10} M_{\odot}$.

## Appendix B: Alternative distances to Nube

Some of the most peculiar properties of Nube depend on the assumed distance to the galaxy. In Section 2.3 we show that it is unlikely that another nearby galaxy could be responsible for the HI emission we detected. However, we cannot completely rule out the possibility that the HI detection is not related to Nube. In that case, the galaxy could be at a different distance from the one assumed in this work. For this reason, it is worth exploring other possible alternative distances.

One possibility is that Nube is a satellite of a galaxy in the field of view other than UGC929. Exploring the nearby galaxies with spectroscopic redshifts (Fig. 1), Nube could be associated with UGC928 ( $8248 \mathrm{~km} / \mathrm{s}$ ). However, if it is associated with UGC928 (projected separation of $\sim 3.8^{\prime}$ ), Nube will be even further away and its properties will be even more extreme. Alternatively, Nube could be closer to us. In Fig. B. 1 we show the inferred properties of Nube at different distances, from 2 Mpc (white) to the assumed distance of 107 Mpc (darker blue). The only galaxy in the field of view with a spectroscopic redshift that is closer to us is UGC931 (Fig. 1) at $\sim 30 \mathrm{Mpc}$ (Springob et al. 2009). If Nube were associated with UGC931, its properties would also be peculiar. Its stellar mass and $R_{e}$ would be similar to those of UDGs, but it would have a significantly lower central surface brightness (third to fourth hexagon from the left, upper right panel in Fig. B.1).

For Nube to be a more normal galaxy, the distance should be between 2 and 10 Mpc . In this case, Nube will be more like the low-mass Local Group galaxies of McConnachie (2012), such as Phoenix ( $\mathrm{M}_{*}=0.7 \times 10^{6} \mathrm{M}_{\odot}, r_{e}=454 \mathrm{pc}$ and $\mu_{V}(0)=25.8 \mathrm{mag} / \mathrm{arcsec}^{2}$ ) or the dwarfs of Carlsten et al. (2021). However, there is no nearby massive galaxy to act as a host for Nube in this case, and therefore it would be a rare example of an isolated very low-mass galaxy.


Fig. B.1. Same as Fig. 8 but showing Nube at different distances: 2, 10, 25, 50, 75, and 107 Mpc .

## Appendix C: Photometric redshifts and other sources around Nube

Table C.1. List of objects with photometric redshifts compatible with the HI detection reported in this paper. The vast majority of these objects are stars or point-like sources (marked with *), another two sources (marked with ${ }^{a}$ ) are part of a known galaxy, UGC931.

| RA <br> $(\mathrm{deg})$ | DEC <br> $(\mathrm{deg})$ | z | $\mathrm{m}_{r}$ <br> $(\mathrm{AB} \mathrm{mag})$ | distance <br> $(\operatorname{arcmin})$ |
| :---: | :---: | :---: | :---: | :---: |
| 20.7455 | -0.6473 | $0.028 \pm 0.03$ | 17.3 | $7.3^{*}$ |
| 20.8104 | -0.6972 | $0.025 \pm 0.03$ | 20.4 | $5.5^{a}$ |
| 20.8105 | -0.7059 | $0.025 \pm 0.02$ | 21.0 | $5.9^{a}$ |
| 20.8349 | -0.6543 | $0.025 \pm 0.1$ | 20.3 | $2.5^{*}$ |
| 20.7561 | -0.5262 | $0.027 \pm 0.05$ | 17.2 | $8.9^{*}$ |
| 20.7682 | -0.5045 | $0.026 \pm 0.2$ | 20.2 | $9.2^{*}$ |
| 20.7984 | -0.6008 | $0.025 \pm 0.1$ | 20.6 | $4.2^{*}$ |
| 20.8535 | -0.5496 | $0.026 \pm 0.1$ | 20.0 | $4.5^{*}$ |
| 20.8571 | -0.5475 | $0.025 \pm 0.01$ | 12.1 | $4.6^{*}$ |
| 20.9294 | -0.4799 | $0.026 \pm 0.09$ | 20.4 | $9.5^{*}$ |
| 21.0163 | -0.6003 | $0.027 \pm 0.15$ | 20.5 | $9.2^{*}$ |

Table C.2. List of faint sources around Nube that we cannot exclude as responsible for the HI emission reported in this work.

| RA <br> $(\mathrm{deg})$ | DEC <br> $(\mathrm{deg})$ | Distance to Nube <br> $(\operatorname{arcmin})$ |
| :---: | :---: | :---: |
| 20.9390 | -0.6132 | 4.5 |
| 20.9092 | -0.5122 | 7.2 |

Given the spatial resolution of the GBT data, in Section 2.3 we discuss the possibility that the HI detection corresponds to another source and not to Nube. Table C. 1 lists all galaxies with photometric redshifts within $10^{\prime}$ of Nube. Table C. 2 lists faint sources that have neither spectroscopic nor photometric redshifts, but could be responsible for the HI emission.

## Appendix D: Masks for photometry and surface brightness profiles

Masking is a key component in obtaining reliable photometry and surface brightness profiles of the galaxy under study. In Section 3.1 we describe how the HiPERCAM mask used to derive the photometry and surface brightness profiles of Nube was created. Figure D. 1 shows the mask used to derive the photometry and surface brightness profiles of the galaxy.


Fig. D.1. Mask (blue regions) applied to a HiPERCAM RGB colour $(\mathrm{g}+\mathrm{r}+\mathrm{i})$ image of a region $100^{\prime \prime} \times 100^{\prime \prime}$ around Nube. The black and white background is a $g+r$ image. The image demonstrates the need for thorough masking in these ultradeep images.


[^0]:    * Reduced images are available at the CDS via anonymous ftp to cdsarc.cds.unistra.fr (130.79.128.5) or via https:// cdsarc.cds.unistra.fr/viz-bin/cat/J/A+A/681/A15

[^1]:    ${ }^{1}$ \noo-beh $\backslash$. Cloud in Spanish.

[^2]:    2 The Green Bank Observatory is a facility of the National Science Foundation operated under cooperative agreement by Associated Universities, Inc.

[^3]:    ${ }^{3}$ http://gbtidl.nrao.edu/

[^4]:    4 The NASA/IPAC Extragalactic Database (NED) is operated by the Jet Propulsion Laboratory, California Institute of Technology, under contract with the National Aeronautics and Space Administration.

[^5]:    5 This M/L value is consistent with the $\mathrm{M} / \mathrm{L}$ ratios derived at each radial distance using the $g-r$ colour profile in Sect. 3.1; a median of 1.7 .

[^6]:    $\overline{6}$ This calculation assumes that Nube is in dynamical equilibrium. This is a reasonable assumption given the large distance to its potential progenitor UGC 929.

[^7]:    7 See Muñoz et al. (2014), Lim et al. (2020), Cantiello et al. (2020) for other examples of using colour-colour diagrams to identify candidate GCs.

[^8]:    8 Taylor et al. (2017) use the GCs of NGC 5128 to generate their catalogue. NGC 5128 is a massive galaxy with intense merger activity. Therefore, the population of GCs in this massive galaxy includes GCs with a wide variety of origins, from those formed in situ in this massive galaxy to those accreted by mergers with a large number of dwarfs. In this sense, the Taylor et al. (2017) catalogue is expected to contain a very heterogeneous and complete sample of GCs.

[^9]:    9 To derive the mass of VLSB-C, we assume a $B-V=0.7 \mathrm{mag}$ for this galaxy, which is typical for dwarf galaxies and the $\mathrm{M} / \mathrm{L}$ colour relation in Bell et al. (2003).
    ${ }^{10}$ Other estimates for $r_{\mathrm{e}}$ for DF44 are $r_{\mathrm{e}}=3.3 \pm 0.3 \mathrm{kpc}$ (Chamba et al. 2020 ) and $3.9 \pm 0.7 \mathrm{kpc}$ (Saifollahi et al. 2021).

