

HYPOELLIPTICITY OF THE $\bar{\partial}$ -NEUMANN PROBLEM AT EXPONENTIALLY DEGENERATE POINTS

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ABSTRACT. We prove local hypoellipticity of the complex Laplacian \square in a domain which has compactness estimates, is of finite type outside a curve transversal to the CR directions and for which the holomorphic tangential derivatives of a defining function are subelliptic multipliers in the sense of Kohn.

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1. INTRODUCTION

For the pseudoconvex domain $\Omega \subset \mathbb{C}^n$ whose boundary is defined in coordinates $z = x + iy$ of \mathbb{C}^n , by

$$(1.1) \quad 2x_n = \exp\left(-\frac{1}{(\sum_{j=1}^{n-1} |z_j|^2)^{\frac{s}{2}}}\right), \quad s > 0,$$

the tangential Kohn Laplacian $\square_b = \bar{\partial}_b \bar{\partial}_b^* + \bar{\partial}_b^* \bar{\partial}_b$ as well as the full Laplacian $\square = \bar{\partial} \bar{\partial}^* + \bar{\partial}^* \bar{\partial}$ show very interesting features especially in comparison with the “tube domain” whose boundary is defined by

$$(1.2) \quad 2x_n = \exp\left(-\frac{1}{(\sum_{j=1}^{n-1} |x_j|^2)^{\frac{s}{2}}}\right), \quad s > 0.$$

(Here z_j have been replaced by x_j at exponent.) Energy estimates are the same for the two domains. For the problem on the boundary $b\Omega$, they come as

$$(1.3) \quad \|(\log \Lambda)^{\frac{1}{s}} u\|_{b\Omega} \lesssim \|\bar{\partial}_b u\|_{b\Omega}^2 + \|\bar{\partial}_b^* u\|_{b\Omega}^2 + \|u\|_{b\Omega}^2$$

for any smooth compact support form $u \in C_c^\infty(b\Omega)^k$ of degree $k \in [1, n-2]$.

Here $\log \Lambda$ is the tangential pseudodifferential operator with symbol $\log(1 + |\xi'|^2)^{\frac{1}{2}}$, $\xi' \in \mathbb{R}^{2n-1}$, the dual real tangent space. As for the problem on the domain Ω , one has simply to replace $\bar{\partial}_b, \bar{\partial}_b^*$ by $\bar{\partial}, \bar{\partial}^*$ and take norms over Ω for forms u in $D_{\bar{\partial}^*}$, the domain of $\bar{\partial}^*$, of degree $1 \leq k \leq n-1$; this can be seen, for instance, in [9]. In particular, these are superlogarithmic (resp. compactness) estimates if $s < 1$ (resp. for

any $s > 0$). A related problem is that of the local hypoellipticity of the Kohn Laplacian \square_b or, with equivalent terminology, the local regularity of the inverse (modulo harmonics) operator $N_b = \square_b^{-1}$. Similar is the notion of hypoellipticity of the Laplacian \square or the regularity of the inverse Neumann operator $N = \square^{-1}$. It has been proved by Kohn in [12] that superlogarithmic estimates suffice for local hypoellipticity of the problem both in the boundary and in the domain. (Note that hypoellipticity for the domain, [12] Theorem 8.3, is deduced from microlocal hypoellipticity for the boundary, [12] Theorem 7.1, but a direct proof is also available, [7] Theorem 5.4.) In particular, for (1.1) and (1.2), there is local hypoellipticity when $s < 1$.

As for the more delicate hypoellipticity, in the uncertain range of indices $s \geq 1$, only the tangential problem has been studied and the striking conclusion is that the behavior of (1.1) and (1.2) split. The first stays always hypoelliptic for any s (Kohn [11]) whereas the second is not for $s \geq 1$ (Christ [4]). When one tries to relate $(\bar{\partial}_b, \bar{\partial}_b^*)$ on $b\Omega$ to $(\bar{\partial}, \bar{\partial}^*)$ on Ω , estimates go well through (Kohn [12] Section 8 and Khanh [7] Chapter 4) but not regularity. In particular, the two conclusions about tangential hypoellipticity of \square_b for (1.1) and non-hypoellipticity for (1.2) when $s \geq 1$, cannot be automatically transferred from $b\Omega$ to Ω . Now, for the non-hypoellipticity in Ω in case of the tube (1.2) we have obtained with Baracco in [1] a result of propagation which is not equivalent but intimately related. The real lines x_j are propagators of holomorphic extendibility from Ω across $b\Omega$. What we prove in the present paper is hypoellipticity in Ω for (1.1) when $s \geq 1$.

Theorem 1.1. *Let Ω be a pseudoconvex domain of \mathbb{C}^n in a neighborhood of $z_o = 0$ and assume that the $\bar{\partial}$ -Neumann problem satisfies the following properties*

- (i) *there are local compactness estimates,*
- (ii) *there are subelliptic estimates for $(z_1, \dots, z_{n-1}) \neq 0$,*
- (iii) *$\partial_{z_j} r$, $j = 1, \dots, n-1$, are subelliptic multipliers (cf. [10]).*

Then \square is locally hypoelliptic at z_o .

The proof follows in Section 2. It consists in relating the system on Ω to the tangential system on $b\Omega$ along the guidelines of [12] Section 8, and then in using the argument of [11] simplified by the additional assumption (i).

Remark 1.2. The domain with boundary (1.1), but not (1.2), satisfies the hypotheses of Theorem 1.1 for any $s > 0$: (i) is obvious, and (ii) and (iii) are the content of [11] Section 4.

Notice that $\partial\Omega$ is given only locally in a neighborhood of z_o . We can continue $\partial\Omega$ leaving it unchanged in a neighborhood of z_o , making it strongly pseudoconvex elsewhere, in such a way that it bounds a relatively compact domain $\Omega \subset\subset \mathbb{C}^n$ (cf. [14]). In this situation \square is hypoelliptic at every boundary point. Also, it is well defined a H^0 inverse Neumann operator $N = \square^{-1}$, and, by Theorem 1.1, the $\bar{\partial}$ -Neumann solution operator $\bar{\partial}^*N$ preserves $C^\infty(\bar{\Omega})$ -smoothness. It even preserves the exact Sobolev class H^s according to Theorem 2.7 below. In other words, the canonical solution $u = \bar{\partial}^*Nf$ of $\bar{\partial}u = f$ for $f \in \text{Ker } \bar{\partial}$ is H^s exactly at the points of $b\Omega$ where f is H^s . The Bergman projection B also preserves $C^\infty(\bar{\Omega})$ -smoothness on account of Kohn's formula $B = \text{Id} - \bar{\partial}^*N\bar{\partial}$.

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2. HYPOELLIPTICITY OF \square AND EXACT HYPOELLIPTICITY OF $\bar{\partial}^*N$

We state properly hypoellipticity and exact hypoellipticity of a general system (P_j) .

Definition 2.1. (i) The system (P_j) is locally hypoelliptic at $z_o \in b\Omega$ if

$$P_j u \in C^\infty(\bar{\Omega})_{z_o}^k \text{ for any } j \text{ implies } u \in C^\infty(\bar{\Omega})_{z_o}^k,$$

where $C^\infty(\bar{\Omega})_{z_o}^k$ denotes the set of germs of k -forms smooth at z_o .

(ii) The system (P_j) is exactly locally hypoelliptic at $z_o \in b\Omega$ when there is a neighborhood U of z_o such that for any pair of cut-off functions ζ and ζ' in $C_c^\infty(U)$ with $\zeta'|_{\text{supp}(\zeta)} \equiv 1$ we have for any s and for suitable c_s

$$(2.1) \quad \|\zeta u\|_s^2 \leq c_s \left(\sum_j \|\zeta' P_j u\|_s^2 + \|u\|_0^2 \right), \quad u \in C^\infty(\bar{\Omega})^k \cap D_{(P_j)}.$$

If (P_j) happens to have an inverse, this is said to be locally regular and locally exactly regular in the situation of (i) and (ii) respectively.

Remark 2.2. By Kohn-Nirenberg [13] the assumption $u \in C^\infty$ can be removed from (2.1). Precisely, by the elliptic regularization, one can prove that if $\zeta' P_j u \in H^s$ and $\zeta' u \in H^0$, then $\zeta u \in H^s$ and satisfies (2.1). This motivates the word ‘‘exact’’, that is, Sobolev exact. Not only the local C^∞ - but also the H^s -smoothness passes from $P_j u$ to u .

Let ϑ be the formal adjoint of $\bar{\partial}$ and $\Delta = \bar{\partial}\vartheta + \vartheta\bar{\partial}$ the Laplacian; it acts on forms by the action of the usual Laplacian on its coefficients.

If $u \in D_{\square}$, then $\square u = \Delta u$. We first prove exact hypoellipticity of the system $(\bar{\partial}, \bar{\partial}^*, \Delta)$; hypoellipticity of \square itself will follow by the method of Boas-Straube.

Theorem 2.3. *In the situation of Theorem 1.1, we have, for a neighborhood U of z_o and for any couple of cut-off ζ and ζ' with $\zeta'|_{\text{supp}\zeta} \equiv 1$*

$$(2.2) \quad \|\zeta u\|_s^2 \lesssim \|\zeta' \bar{\partial} u\|_s^2 + \|\zeta' \bar{\partial}^* u\|_s^2 + \|\zeta' \Delta u\|_{s-2}^2 + \|u\|_0^2, \quad u \in D_{\bar{\partial}^*}.$$

In particular, the system $(\bar{\partial}, \bar{\partial}^, \Delta)$ is exactly locally hypoelliptic at $z_o = 0$.*

Remark 2.4. The hypoellipticity of \square_b under (ii) and (iii) of Theorem 1.1 is proved by Kohn in [11]. It does not require (i) but it is not exact hypoellipticity (the neighborhood U of (2.1) depends on s). However, inspection of his proof shows that, if (i) is added, then in fact (2.1) holds for $(P_j) = \square_b$. Our proof consists in a reduction to the tangential system.

Proof. We proceed in several steps which are highlighted in two intermediate propositions. We use the standard notation $Q(u, u)$ for $\|\bar{\partial} u\|_0^2 + \|\bar{\partial}^* u\|_0^2$ and some variants as, for an operator Op , $Q_{Op}(u, u) := \|Op \bar{\partial} u\|_0^2 + \|Op \bar{\partial}^* u\|_0^2$; most often, in our paper, Op is chosen as $\Lambda^s \zeta'$. We decompose a form u as

$$\begin{cases} u = u^\tau + u^\nu, \\ u^\tau = u^{\tau+} + u^{\tau-} + u^{\tau 0}, \end{cases}$$

where the first is the decomposition in tangential and normal component and the second is the microlocal decomposition $u^{\tau \pm 0} = \Psi^{\pm 0} u^\tau$ in which $\Psi^{\pm 0}$ are the tangential pseudodifferential operators whose symbols $\psi^{\pm 0}$ are a conic decomposition of the unity in the space dual to \mathbb{R}^{2n-1} the real orthogonal to ∂r (cf. Kohn [12]). We begin our proof by remarking that any of the forms $u^\# = u^\nu, u^{\tau-}, u^{\tau 0}$ enjoys elliptic estimates

$$(2.3) \quad \|\zeta u^\#\|_s^2 \lesssim \|\zeta' \bar{\partial} u^\#\|_{s-1}^2 + \|\zeta' \bar{\partial}^* u^\#\|_{s-1}^2 + \|u^\#\|_0^2 \quad s \geq 2.$$

We refer to [6] formula (1) of Main theorem as a general reference but also give an outline of the proof. For this, we have to call into play the tangential s -Sobolev norm which is defined by $\|u\|_s = \|\Lambda^s u\|_0$. We start from

$$(2.4) \quad \|\zeta u^\#\|_1^2 \lesssim Q(\zeta u^\#, \zeta u^\#) + \|u^\#\|_0^2;$$

this is the basic estimate for u^ν (which vanishes at $b\Omega$) whereas it is [12] Lemma 8.6 for $u^{\tau-}$ and $u^{\tau 0}$. Applying (2.4) to $\zeta'\Lambda^{s-1}\zeta u^\#$ one gets the estimate of tangential norms for any s . Finally, by non-characteristicity of $(\bar{\partial}, \bar{\partial}^*)$ one passes from tangential to full norms along the guidelines of [16] Theorem 1.9.7. The version of this argument for \square can be found in [12] second part of p. 245. Because of (2.3), it suffices to prove (2.2) for the only $u^{\tau+}$. We further decompose

$$u^{\tau+} = u^{\tau+(h)} + u^{\tau+(0)},$$

where $u^{\tau+(h)}$ is the ‘‘harmonic extension’’ in the sense of Kohn [12] and $u^{\tau+(0)}$ is just the complementary part. We denote by $\bar{\partial}^\tau$ the extension of $\bar{\partial}_b$ from $b\Omega$ to Ω which stays tangential to the level surfaces $r \equiv \text{const}$. It acts on tangential forms u^τ and it is defined by $\bar{\partial}^\tau u^\tau = (\bar{\partial} u^\tau)^\tau$. We denote by $\bar{\partial}^{\tau*}$ its adjoint; thus $\bar{\partial}^{\tau*} u^\tau = \bar{\partial}^*(u^\tau)$. We use the notations \square^τ and Q^τ for the corresponding Laplacian and energy. We notice that over a tangential form u^τ we have a decomposition

$$(2.5) \quad Q = Q^\tau + \|\bar{L}_n u^\tau\|_0^2.$$

The proof of (2.2) for $u^{\tau+}$ requires two crucial technical results. Here is the first which is the most central

Proposition 2.5. *For the harmonic extension $u^{\tau+(h)}$ we have*

$$(2.6) \quad \|\|\zeta u^{\tau+(h)}\|\|_s^2 \lesssim Q_{\Lambda^s \zeta'}^\tau(u^{\tau+(h)}, u^{\tau+(h)}) + \|u^{\tau+(h)}\|_0^2.$$

Proof. We apply compactness estimates (cf. e.g. [7] Section 6) for $\zeta'\Lambda^s \zeta u^{\tau+(h)}$,

$$(2.7) \quad \|\|\zeta'\Lambda^s \zeta u^{\tau+(h)}\|\|^2 \leq \epsilon Q(\zeta'\Lambda^s \zeta u^{\tau+(h)}, \zeta'\Lambda^s \zeta u^{\tau+(h)}) + c_\epsilon \|\|\zeta'\Lambda^s \zeta u^{\tau+(h)}\|\|_{-1}^2.$$

We decompose Q according to (2.5). We calculate Q^τ over $\zeta'\Lambda^s \zeta u^{\tau+(h)}$ and compute errors coming from commutators $[Q^\tau, \zeta'\Lambda^s \zeta]$. In this calculation we assume that the cut off functions are of product type $\zeta(z')\zeta(t)$ where z' (resp. t) are complex (resp. totally real) tangential coordinates in $T_{z_0} b\Omega$. We have

$$(2.8) \quad \begin{aligned} & Q^\tau(\zeta'\Lambda^s \zeta u^{\tau+(h)}, \zeta'\Lambda^s \zeta u^{\tau+(h)}) \\ & \lesssim Q_{\zeta'\Lambda^s \zeta}^\tau(u^{\tau+(h)}, u^{\tau+(h)}) + \|\|\zeta u^{\tau+(h)}\|\|_s^2 + \|\|\zeta' u^{\tau+(h)}\|\|_{s-1}^2 \\ & + \left(\|\|(|\dot{\zeta}(z')| + |\dot{\zeta}'(z')|)\Lambda^s u^{\tau+(h)}\|\|_0^2 + \|\| \sum_{j=1}^{n-1} |r_{z_j}| (|\dot{\zeta}(t)| + |\dot{\zeta}'(t)|) \Lambda^s u^{\tau+(h)}\|\|_0^2 \right). \end{aligned}$$

We explain (2.8). First, the commutators $[\bar{\partial}^\tau, \zeta' \Lambda^s \zeta]$ (and similarly as for $[\bar{\partial}^{\tau*}, \zeta' \Lambda^s \zeta]$) are decomposed by Jacobi identity as

$$[\bar{\partial}^\tau, \zeta' \Lambda^s \zeta] = [\bar{\partial}^\tau, \zeta'] \Lambda^s \zeta + \zeta' [\bar{\partial}^\tau, \Lambda^s] \zeta + \zeta' \Lambda^s [\bar{\partial}^\tau, \zeta].$$

The central commutator $[\bar{\partial}^\tau, \Lambda^s]$ produces the error term $\|\zeta u^{\tau+(h)}\|_s^2$. As for the two others, we have

$$[\bar{\partial}^\tau, \zeta(z') \zeta(t)] = [\bar{\partial}^\tau, \zeta(z')] \zeta(t) + \zeta(z') [\bar{\partial}^\tau, \zeta(t)],$$

and similarly for ζ replaced by ζ' and $\bar{\partial}^\tau$ by $\bar{\partial}^{\tau*}$. Now,

$$(2.9) \quad [\bar{\partial}^\tau, \zeta(z')] \sim \dot{\zeta}(z').$$

On the other hand, we first notice that it is not restrictive to assume that $\partial_{z_1}, \dots, \partial_{z_{n-1}}$ are a basis of $T_0^{1,0} b\Omega$ for otherwise, owing to (iii), we have subelliptic estimates from which local regularity readily follows. Thus, each \bar{L}_j , $j = 1, \dots, n-1$, is of type $\bar{L}_j = r_{\bar{z}_j} \partial_{\bar{z}_n} - r_{\bar{z}_n} \partial_{\bar{z}_j}$, and then

$$(2.10) \quad \begin{aligned} [\bar{\partial}^\tau, \zeta(t)] &\sim \sum_{j=1}^{n-1} [\bar{L}_j, \zeta(t)] \\ &\sim \sum_{j=1}^{n-1} r_{\bar{z}_j} \dot{\zeta}(t). \end{aligned}$$

By combining (2.9) with (2.10) (and using the analogous for ζ' and $\bar{\partial}^{\tau*}$), we get the last line of (2.8). This establishes (2.8). Next, since $(\bar{\partial}^\tau, \bar{\partial}^{\tau*})$ has subelliptic estimates, say η -subelliptic, for $z' \neq 0$ and hence in particular over $\text{supp } \dot{\zeta}(z')$ and $\text{supp } \dot{\zeta}'(z')$ and since the $r_{\bar{z}_j}$ are, say, η -subelliptic multipliers even at $z' = 0$, then the last line of (2.8) is estimated by $\|\zeta'' \Lambda^{s-\eta} \zeta' u^{\tau+(h)}\|^2$ where $\zeta'' \equiv 1$ over $\text{supp } \zeta'$. This shows, using iteration over increasing k such that $k\eta > s$ and over decreasing j from $s-1$ to 0, that (2.7) and (2.8) imply (2.6) provided that we add on the right side the extra term $\|\bar{L}_n \zeta' \Lambda^s \zeta u^{\tau+(h)}\|^2$. Note that, as a result of the inductive process, we have to replace $Q_{\zeta' \Lambda^s \zeta}$ in (2.8) by $Q_{\Lambda^s \zeta'}$ in (2.6).

Up to this point the argument is the same as in [11] and does not make any use of the specific properties of the harmonic extension $u^{\tau+(h)}$. We start the new part which is dedicated to prove that $\|\bar{L}_n \zeta' \Lambda^s \zeta u^{\tau+(h)}\|^2$ can be removed from the right of (2.6). For this we have to use the main property of this extension expressed by [12] Lemma 8.5, that is,

$$(2.11) \quad \|\bar{L}_n \zeta u^{\tau+(h)}\|_0^2 \lesssim \sum_{j=1}^{n-1} \|\bar{L}_j \zeta u_b^{\tau+}\|_{b, -\frac{1}{2}}^2 + \|u^{\tau+}\|_0^2.$$

Note that (2.11) differs from [12] Lemma 8.5 by $[\bar{L}_n, \Psi^+]$; but this is an error term which can be taken care of by $u^{\tau+0}$ to which elliptic estimates apply. Applying (2.11) to $\zeta' \Lambda^s \zeta u^{\tau+(h)}$ (for the first inequality below), and using the classical inequality $\|\cdot\|_{b, -\frac{1}{2}}^2 \leq c_\epsilon \|\cdot\|_0^2 + \epsilon \|\partial_r \cdot\|_{-1}^2$ (cf. e.g. [8] (1.10)) together with the splitting $\partial_r = \bar{L}_n + Tan$ (for the second), we get

$$\begin{aligned}
(2.12) \quad \|\bar{L}_n \zeta' \Lambda^s \zeta u^{\tau+(h)}\|_0^2 &\lesssim \sum_{j=1}^{n-1} \|\bar{L}_j \zeta' \Lambda^s \zeta u_b^{\tau+}\|_{b, -\frac{1}{2}}^2 + \|\zeta' \Lambda^s \zeta u^{\tau+}\|_0^2 \\
&\stackrel{\text{by (2.11)}}{\lesssim} c_\epsilon \sum_{j=1}^{n-1} \|\bar{L}_j \zeta' \Lambda^s \zeta u^{\tau+(h)}\|_0^2 + \epsilon \sum_{j=1}^{n-1} \|[\bar{L}_n, \bar{L}_j] \zeta' \Lambda^s \zeta u^{\tau+(h)}\|_{-1}^2 \\
&\quad + \epsilon \sum_{j=1}^{n-1} \|Tan \bar{L}_j \zeta' \Lambda^s \zeta u^{\tau+(h)}\|_{-1}^2 + \|\zeta' \Lambda^s \zeta u^{\tau+(h)}\|_0^2.
\end{aligned}$$

The first term on the right of the last inequality is controlled by $\sum_{j=1}^{n-1} \|\zeta' \Lambda^s \zeta \bar{L}_j u^{\tau+(h)}\|^2 + \|\zeta u^{\tau+(h)}\|_s^2 + \|\zeta'' u^{\tau+(h)}\|_{s-1}^2$ by the first part of the proposition; moreover, we have the immediate estimate $\sum_{j=1}^{n-1} \|\zeta' \Lambda^s \zeta \bar{L}_j u^{\tau+(h)}\|^2 \lesssim Q_{\Lambda^s \zeta'}^\tau(u^{\tau+(h)}, u^{\tau+(h)})$. The term which carries ϵTan , after Tan has been annihilated by the Sobolev norm of index -1 , has the same estimate as the first term. It remains to control the second term in the right which involves $\epsilon \bar{L}_n$. We rewrite $\bar{L}_n \bar{L}_j = \bar{L}_j \bar{L}_n + [\bar{L}_n, \bar{L}_j]$; when \bar{L}_j moves in first position, it is annihilated by -1 and what remains is absorbed in the left. As for the commutator, we have

$$\begin{aligned}
\|[\bar{L}_n, \bar{L}_j] \zeta' \Lambda^s \zeta u^{\tau+(h)}\|_{-1}^2 &\lesssim \|\zeta u^{\tau+(h)}\|_s^2 + \|\partial_r \zeta' \Lambda^s \zeta u^{\tau+(h)}\|_{-1}^2 \\
&\lesssim \|\zeta u^{\tau+(h)}\|_s^2 + \|\bar{L}_n \zeta' \Lambda^s \zeta u^{\tau+(h)}\|_{-1}^2,
\end{aligned}$$

where we have used the splitting $\partial_r = Tan + \bar{L}_n$ in the second inequality. Again, the term with \bar{L}_n , which now comes in -1 norm, is absorbed in the left of (2.12). Summarizing up, we have got

$$\begin{aligned}
(2.13) \quad \|\bar{L}_n \zeta' \Lambda^s \zeta u^{\tau+(h)}\|_0^2 &\lesssim c_\epsilon Q_{\Lambda^s \zeta'}^\tau(u^{\tau+(h)}, u^{\tau+(h)}) \\
&\quad + \|\zeta u^{\tau+(h)}\|_s^2 + \|\zeta'' u^{\tau+(h)}\|_{s-1}^2.
\end{aligned}$$

But $\|\bar{L}_n \cdot\|_{-1}^2$ comes with a factor ϵ of compactness and hence the term in s -norm in the last line can be absorbed in the left of the initial

inequalities (2.7) or (2.6). Finally, we use an inductive argument to go down from $s - 1$ to 0. This concludes the proof of the proposition. \square

We remark now that

$$\begin{aligned}
(2.14) \quad & \|\zeta u^{\tau+(h)}\|_0^2 \lesssim \|\zeta u_b^{\tau+}\|_{b, -\frac{1}{2}}^2 \\
& \lesssim \|\zeta u^{\tau+}\|_0^2 + \|\partial_r \zeta u^{\tau+}\|_{-1}^2 \\
& \lesssim \|\zeta u^{\tau+}\|_0^2 + \|\bar{L}_n \zeta u^{\tau+}\|_{-1}^2 + \|Tan \zeta u^{\tau+}\|_{-1}^2 \\
& \lesssim Q_{\Lambda^{-1}\zeta}(u^{\tau+}, u^{\tau+}) + \|\zeta u^{\tau+}\|_0^2.
\end{aligned}$$

The same inequality also holds for $u^{\tau+(h)}$ replaced by $u^{\tau+(0)}$ on account of the identity $u^{\tau+(0)} = u^{\tau+} + u^{\tau+(h)}$. We need another preparation result

Proposition 2.6. *We have*

$$(2.15) \quad Q_{\Lambda^s \zeta'}(u^{\tau+(h)}, u^{\tau+(h)}) \lesssim Q_{\Lambda^s \zeta'}(u^{\tau+}, u^{\tau+}) + Q_{\partial_r \Lambda^{s-1} \zeta'}(u^{\tau+}, u^{\tau+})$$

and

$$\begin{aligned}
(2.16) \quad & \|\zeta u^{\tau+(0)}\|_s^2 \lesssim Q_{\Lambda^{s-1} \zeta'}(u^{\tau+}, u^{\tau+}) + Q_{\partial_r \Lambda^{s-2} \zeta'}(u^{\tau+}, u^{\tau+}) \\
& + \|\zeta' \Delta u^{\tau+}\|_{s-2}^2 + \|u^{\tau+}\|_0^2.
\end{aligned}$$

Proof. The proof of (2.15) is an immediate combination of the formulas $\|\zeta' u^{\tau+(h)}\|_0 \lesssim \|\zeta' u_b^{\tau+}\|_{b, -\frac{1}{2}}$ and $\|\zeta' u^{\tau+}\|_{b, -\frac{1}{2}} \lesssim \|\zeta' u^{\tau+}\|_0 + \|\partial_r \zeta' u^{\tau+}\|_{-1}^2$.

We prove now (2.16). By elliptic estimate for $u^{\tau+(0)}$ (which vanishes at $b\Omega$) with respect to the order 2 elliptic operator Δ , we have

$$(2.17) \quad \|\zeta u^{\tau+(0)}\|_s^2 \lesssim \|\zeta' \Delta u^{\tau+(0)}\|_{s-2}^2 + \|u^{\tau+(0)}\|_0^2.$$

This result of Sobolev regularity at the boundary is very classical: it is formulated, for functions in H_0^1 such as the coefficients of $u^{\tau+(0)}$, e.g. in Evans [5] Theorem 5 p. 323. Owing to the identity $\Delta u^{\tau+(0)} = \Delta u^{\tau+} + P^1 u^{\tau+(h)}$ for a 1-order operator P^1 (cf. [12] p. 241), we can replace $\Delta u^{\tau+(0)}$ by $\Delta u^{\tau+}$ on the right side of (2.17) putting the contribution of P^1 into an error term of type $\|\zeta' u^{\tau+(h)}\|_{s-1} + \|\zeta' \partial_r u^{\tau+(h)}\|_{s-2}$, which can be estimated, on account of the splitting $\partial_r = \bar{L}_n + Tan$, by $\|\zeta' u^{\tau+(h)}\|_{s-1} + \|\zeta'' u^{\tau+(h)}\|_{s-2} + Q_{\Lambda^{s-2} \zeta'}(u^{\tau+(h)}, u^{\tau+(h)})$. We write the terms of order $s - 1$ and $s - 2$ as a common $\|\zeta'' u^{\tau+(h)}\|_{s-1}$ that we can estimate, using (2.6) and (2.15), by

$$\|\zeta'' u^{\tau+(h)}\|_{s-1}^2 \lesssim Q_{\Lambda^{s-1} \zeta'''}(u^{\tau+}, u^{\tau+}) + Q_{\Lambda^{s-2} \partial_r \zeta'''}(u^{\tau+}, u^{\tau+}).$$

This brings down from $s - 1$ to 0 the Sobolev index in the error term. This 0-order term $\|u^{\tau+(h)}\|_0^2$, together with its companion $\|u^{\tau+(0)}\|_0^2$ in the right of (2.17), is estimated, because of (2.14), by $\|u^{\tau+}\|_0^2$ up to a term $Q_{\Lambda^{-1}\zeta}$ which is controlled by the right side of (2.16). This concludes the proof of (2.16). \square

End of proof of Theorem 2.3. We prove (2.2) for $u^{\tau+}$; this implies the conclusion in full generality according to the first part of the proof. We have

$$(2.18) \quad \begin{aligned} \|\zeta u^{\tau+(h)}\|_s^2 &\underset{\sim}{\leq} Q_{\Lambda^s \zeta'}(u^{\tau+(h)}, u^{\tau+(h)}) + \|u^{\tau+(h)}\|_0^2 \\ &\underset{\text{by (2.6)}}{\lesssim} \\ &\underset{\text{by (2.15) and (2.14)}}{\lesssim} Q_{\Lambda^s \zeta'}(u^{\tau+}, u^{\tau+}) + Q_{\partial_r \Lambda^{s-1} \zeta'}(u^{\tau+}, u^{\tau+}) + \|u^{\tau+}\|_0^2. \end{aligned}$$

We combine (2.18) with (2.16); what we get is

$$(2.19) \quad \begin{aligned} \|\zeta u^{\tau+}\|_s^2 &\leq \|\zeta u^{\tau+(h)}\|_s^2 + \|\zeta u^{\tau+(0)}\|_s^2 \\ &\underset{\sim}{\leq} \|\zeta' \bar{\partial} u^{\tau+}\|_s^2 + \|\zeta' \bar{\partial}^* u^{\tau+}\|_s^2 + \|\zeta' \Delta u^{\tau+}\|_{s-2}^2 + \|u^{\tau+}\|_0^2. \end{aligned}$$

By the non-characteristicity of Q , we can replace the tangential norm $\|\cdot\|_s$ by the full norm $\|\cdot\|_s$ in the left of (2.19). (The explanation of this point can be found, for example, in [12] second part of p. 245.) This proves (2.2) for $u^{\tau+}$ and thus also for a general u . \square

We modify $b\Omega$ outside a neighborhood of z_o where it satisfies the hypotheses of Theorem 1.1 so that it is strongly pseudoconvex in the modified portion and bounds a relatively compact domain; in particular, there is well defined the H^0 inverse N of \square in this domain. There is an immediate crucial consequence of Theorem 2.3.

Theorem 2.7. *We have that*

$$(2.20) \quad \bar{\partial}^* N \text{ is exactly regular over } \text{Ker } \bar{\partial}$$

and

$$(2.21) \quad \bar{\partial} N \text{ is exactly regular over } \text{Ker } \bar{\partial}^*.$$

Proof. As for (2.20), we put $u = \bar{\partial}^* N f$ for $f \in \text{Ker } \bar{\partial}$. We get

$$\begin{cases} \bar{\partial} u = f, \\ \bar{\partial}^* u = 0, \\ \Delta u = (\vartheta \bar{\partial} + \bar{\partial} \vartheta) \bar{\partial}^* N f \\ \quad = \vartheta (\bar{\partial} \bar{\partial}^* + \bar{\partial}^* \bar{\partial}) N f + \bar{\partial} \vartheta \bar{\partial}^* N f \\ \quad = \vartheta \square N f = \vartheta f. \end{cases}$$

Thus, by (2.2)

$$(2.22) \quad \begin{aligned} \|\zeta u\|_s^2 &\lesssim \|\zeta' f\|_s^2 + \|\zeta' \vartheta f\|_{s-2}^2 + \|u\|_0^2 \\ &\lesssim \|\zeta' f\|_s^2 + \|u\|_0^2. \end{aligned}$$

To prove (2.21), we put $u = \bar{\partial} N f$ for $f \in \text{Ker } \bar{\partial}^*$. We have a similar calculation as above which leads to the same formula as (2.22) (with the only difference that ϑ is replaced by $\bar{\partial}$ in the intermediate inequality). Thus from (2.22) applied both for $\bar{\partial}^* N$ and $\bar{\partial} N$ on $\text{Ker } \bar{\partial}$ and $\text{Ker } \bar{\partial}^*$ respectively, we conclude that these operators are exactly regular. \square

We are ready for the proof of Theorem 1.1. This follows from Theorem 2.7 by the method of Boas-Straube.

Proof of Theorem 1.1. From the regularity of $\bar{\partial}^* N$ it follows that the Bergman projection B is also regular. (Notice that exact regularity is perhaps lost by taking $\bar{\partial}$ in B .) We exploit formula (5.36) in [15] in unweighted norms, that is, for $t = 0$:

$$\begin{aligned} N_q &= B_q(N_q \bar{\partial})(\text{Id} - B_{q-1})(\bar{\partial}^* N_q) B_q \\ &\quad + (\text{Id} - B_q)(\bar{\partial}^* N_{q+1}) B_{q+1}(N_{q+1} \bar{\partial})(\text{Id} - B_q). \end{aligned}$$

Now, in the right side, the $\bar{\partial} N$'s and $\bar{\partial}^* N$'s are evaluated over $\text{Ker } \bar{\partial}^*$ and $\text{Ker } \bar{\partial}$ respectively; thus they are exactly regular. The B 's are also regular and therefore such is N . This concludes the proof of Theorem 1.1. \square

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