# ALGEBRAIC RICCATI EQUATION AND $J$-SPECTRAL FACTORIZATION FOR $\mathcal{H}_{\infty}$ SMOOTHING AND DECONVOLUTION* 

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#### Abstract

This paper deals with a general steady-state estimation problem in the $\mathcal{H}_{\infty}$ setting. The existence of the stabilizing solution of the related algebraic Riccati equation (ARE) and of the solution of the associated $J$-spectral factorization problem is investigated. The existence of such solutions is well established if the prescribed attenuation level $\gamma$ is larger than $\gamma_{f}$ (the infimum of the values of $\gamma$ for which a causal estimator with attenuation level $\gamma$ exists). We consider the case when $\gamma \leq \gamma_{f}$ and show that the stabilizing solution of the ARE still exists (except for a finite number of values of $\gamma$ ) as long as a fixed-lag acausal estimator (smoother) does. The stabilizing solution of the ARE may be employed to derive a state-space realization of a minimum-phase $J$-spectral factor of the $J$-spectrum associated with the estimation problem. This $J$-spectral factor may be used, in turn, to compute the minimum-lag smoothing estimator. Some of the aspects of the $J$-spectral factorization problem and the properties of its solutions are discussed in correspondence to the (finite number of) values of $\gamma$ for which the stabilizing solution of the ARE does not exist.


Key words. estimation problems, $J$-spectral factorization, algebraic Riccati equation, filtering, deconvolution, smoothing

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1. Introduction and problem statement. Consider the discrete-time linear system

$$
\begin{align*}
x_{k+1} & =A x_{k}+B w_{k},  \tag{1.1a}\\
y_{k} & =C x_{k}+D w_{k},  \tag{1.1b}\\
z_{k} & =L x_{k}+M w_{k}, \tag{1.1c}
\end{align*}
$$

where $x_{k} \in \mathbb{R}^{n}$ and $y_{k} \in \mathbb{R}^{p}$ are the state and the measurement output vector, respectively, and $w_{k} \in \mathbb{R}^{m}$ is the vector of inputs and disturbances. The to-beestimated signal $z_{k} \in \mathbb{R}^{l}$ is an unaccessible linear combination of the state and the input. This is very general and includes, as particular cases, the state-filtering problem ( $L=I$ and $M=0$ ) and the deconvolution problem $(L=0$ and $M=I)$. See section 1.3 for a brief discussion and references on these problems.

Let $F(z)$ be the transfer function of a causal filter driven by the observations $y_{k}$ whose output $\hat{z}_{k}$ is an estimate of $z_{k}$, and let $e_{k}:=z_{k}-\hat{z}_{k}$. The infinite-horizon $\mathcal{H}_{\infty}$ filtering problem consists in designing an estimator $F(z)$ guaranteeing a prescribed level of attenuation $\gamma$ between the $\ell_{2}$-norm of $w_{k}$ and the $\ell_{2}$-norm of the estimation error $e_{k}$. Introducing the transfer functions

$$
\begin{equation*}
H_{1}(z):=C(z I-A)^{-1} B+D \tag{1.2a}
\end{equation*}
$$

[^0]\[

$$
\begin{equation*}
H_{2}(z):=L(z I-A)^{-1} B+M \tag{1.2b}
\end{equation*}
$$

\]

this problem is equivalent to that of finding a stable causal transfer function $F(z)$ such that

$$
\begin{equation*}
F(z) H_{1}(z)-H_{2}(z) \in \mathcal{R} \mathcal{H}_{\infty}, \quad\left\|F(z) H_{1}(z)-H_{2}(z)\right\|_{\infty}<\gamma \tag{1.3}
\end{equation*}
$$

where $\mathcal{R H}_{\infty}$ denotes the space of real rational matrix functions whose poles lie all in $\{z \in \mathbb{C}:|z|<1\} . \mathcal{R} \mathcal{H}_{\infty}$ is endowed with the infinity norm associating with any $G(z) \in \mathcal{R} \mathcal{H}_{\infty},\|G(z)\|_{\infty}=\sup _{|z|=1}\|G(z)\|$ with $\|G(z)\|$ being the largest singular value of $G(z)$ for each given $z$.

The $\mathcal{H}_{\infty}$ fixed-lag smoothing problem with preview horizon of length $N$ (or $N$ lag smoothing problem) may be defined as the problem of estimating (in the $\mathcal{H}_{\infty}$ framework) the signal $z_{k}$ given the observations up to time $k+N, N \geq 0$. In other words, the only difference with the filtering problem is that the transfer function of the estimator is not required to be causal but has to be of the form $E(z)=z^{N} F(z)$ with $F(z)$ being stable and causal. This may be easily reformulated as the problem of finding a stable causal transfer function $F(z)$ such that

$$
\begin{equation*}
F(z) H_{1}(z)-H_{2, N}(z) \in \mathcal{R} \mathcal{H}_{\infty}, \quad\left\|F(z) H_{1}(z)-H_{2, N}(z)\right\|_{\infty}<\gamma \tag{1.4}
\end{equation*}
$$

with

$$
\begin{equation*}
H_{2, N}(z):=z^{-N} H_{2}(z) \tag{1.5}
\end{equation*}
$$

For a general overview on $\mathcal{H}_{\infty}$ estimation see to [19, 21] and references therein. For the continuous time version of the $\mathcal{H}_{\infty}$ smoothing problem see [18].

Assumptions. Let ${ }^{1}$

$$
\begin{equation*}
B_{1}:=B\left(I-D^{\top}\left(D D^{\top}\right)^{-1} D\right), \quad F:=A-B D^{\top}\left(D D^{\top}\right)^{-1} C \tag{1.6}
\end{equation*}
$$

We make the following standard assumptions:

$$
\begin{equation*}
D D^{\top}>0 \tag{1.7a}
\end{equation*}
$$

$(A, C)$ is detectable, $\left(F, B_{1}\right)$ is reachable,
$F$ is invertible.
Moreover, if $M D^{\top} \neq 0$, we define the auxiliary signal $\xi_{k}:=z_{k}-M D^{\top}\left(D D^{\top}\right)^{-1} y_{k}=$ $\widetilde{L} x_{k}+\widetilde{M} w_{k}$ with $\widetilde{L}:=L-M D^{\top}\left(D D^{\top}\right)^{-1} C$ and $\widetilde{M}:=M-M D^{\top}\left(D D^{\top}\right)^{-1} D$ and observe that, since $\xi$ is obtained from $z$ by subtracting a linear combination of the observed signal $y$, the estimations of $z$ and $\xi$ are equivalent problems (i.e., the estimation error is the same). In addition, a simple computation shows that $\widetilde{M} D^{\top}=0$. Hence, without loss of generality, we assume

$$
\begin{equation*}
M D^{\top}=0 \tag{1.7e}
\end{equation*}
$$

so that the following identity holds:

$$
\begin{equation*}
B M^{\top}=B_{1} M^{\top} \tag{1.8}
\end{equation*}
$$

Eventually, notice that neither $A$ nor $F$ is assumed to be stable.

[^1]From the point of view of the motivating $\mathcal{H}_{\infty}$ estimation problem, we mention that assumptions (1.7a) and (1.7d) do not impair loss of generality. In fact if (1.7a) is not satisfied we may resort to a Silverman transfomation (see [6]) that, employing a spectral interactor matrix, yields an equivalent problem in which $D D^{\top}>0$. If (1.7d) is not satisfied, we may perform a preliminary feedback transformation, as described in [5], and obtain an equivalent problem, in which (1.7d) is satisfied.
1.1. $\boldsymbol{J}$-spectral factorization. For a real transfer matrix function $G(z)$, we define $G^{\Upsilon}(z):=G^{\top}\left(z^{-1}\right)$. Let

$$
H(z):=\left[\begin{array}{cc}
H_{1}(z) & 0  \tag{1.9}\\
H_{2}(z) & I_{l}
\end{array}\right]=\left[\begin{array}{c}
C \\
L
\end{array}\right](z I-A)^{-1}[B \mid 0]+\left[\begin{array}{cc}
D & 0 \\
M & I_{l}
\end{array}\right]
$$

and $J_{i, j}(\gamma):=\operatorname{diag}\left\{I_{i},-\gamma^{2} I_{j}\right\}$. (We shall simply write $J$ and $I$, instead of $J_{i, j}(\gamma)$ and $I_{i}$, when there is no risk of confusion.) The transfer function

$$
\Psi(z):=H(z) J_{m, l}(\gamma) H^{\sim}(z)=\left[\begin{array}{cc}
\Psi_{11}(z) & \Psi_{12}(z)  \tag{1.10}\\
\Psi_{12} \sim(z) & \Psi_{22}(z)
\end{array}\right]
$$

with

$$
\begin{equation*}
\Psi_{11}(z):=H_{1}(z) H_{1}^{\sim}(z), \quad \Psi_{12}(z):=H_{1}(z) H_{2}^{\sim}(z), \quad \Psi_{22}(z):=H_{2}(z) H_{2}^{\sim}(z)-\gamma^{2} I \tag{1.11}
\end{equation*}
$$

is called the $J$-spectrum of the system. A square $J$-spectral factor of $\Psi(z)$ is a square matrix function $\Omega(z)$ such that

$$
\begin{equation*}
\Omega(z) J_{p, l}(\gamma) \widetilde{\Omega^{\prime}(z)=\Psi(z) . . . . ~} \tag{1.12}
\end{equation*}
$$

A square $J$-spectral factor is said to be minimum phase if all its zeros lay in the open unit disk. ${ }^{2}$ See [10] and references therein for further details on $J$-spectral factorization.

The estimation problem is strictly related to the existence of a $J$-spectral factor. Precisely, the existence of an estimator with prescribed attenuation level $\gamma$ is equivalent to the existence of a minimum phase $J$-spectral factor of $\Psi(z)$ having a realization with the same state matrix $A$ and the same output matrix $\left[C^{\top} \mid L^{\top}\right]^{\top}$ of the realization (1.9) of $H(z)$; see [3]. Once obtained such a spectral factor, a simple constructive procedure (described in [3]) furnishes a class of estimators $F(z)$ satisfying (1.3).

### 1.2. Algebraic Riccati equation. Let

$$
H_{0}:=\left[\begin{array}{l}
C  \tag{1.13}\\
L
\end{array}\right], \quad J_{0}:=\left[\begin{array}{c}
D \\
M
\end{array}\right], \quad J_{1}:=J_{0} J_{0}^{\top}-\left[\begin{array}{cc}
0 & 0 \\
0 & \gamma^{2} I
\end{array}\right]=\left[\begin{array}{cc}
D D^{\top} & D M^{\top} \\
M D^{\top} & M M^{\top}-\gamma^{2} I
\end{array}\right]
$$

and consider the algebraic Riccati equation (ARE)
(1.14) $\Delta=A \Delta A^{\top}+B B^{\top}-\left(A \Delta H_{0}^{\top}+B J_{0}^{\top}\right)\left(J_{1}+H_{0} \Delta H_{0}^{\top}\right)^{-1}\left(H_{0} \Delta A^{\top}+J_{0} B^{\top}\right)$.

[^2]It is well known (see, e.g., [21]) that the existence of a proper stable filter $F(z)$ satisfying (1.3) is equivalent to the existence of a symmetric, positive definite, stabilizing, feasible solution of (1.14), i.e., a solution $\Delta_{f}=\Delta_{f}^{\top}>0$ such that

$$
\begin{equation*}
A_{f}:=A-\left(A \Delta_{f} H_{0}^{\top}+B J_{0}^{\top}\right)\left(J_{1}+H_{0} \Delta_{f} H_{0}^{\top}\right)^{-1} H_{0} \tag{1.15}
\end{equation*}
$$

is stable and $J_{1}+H_{0} \Delta_{f} H_{0}^{\top}$ has the same inertia of $J_{p, l}(\gamma)$. In passing, notice that as long as $\left(D D^{\top}+C \Delta_{f} C^{\top}\right)>0$, the latter condition (feasibility) may be written in the form

$$
\begin{equation*}
\gamma^{2} I-M M^{\top}-L \Delta_{f} L^{\top}+L \Delta_{f} C^{\top}\left(D D^{\top}+C \Delta_{f} C^{\top}\right)^{-1} C \Delta_{f} L^{\top}>0 \tag{1.16}
\end{equation*}
$$

If we now define

$$
\begin{equation*}
\gamma_{f}:=\inf \{\gamma: \text { there exists a proper stable filter } F(z) \text { satisfying }(1.3)\} \tag{1.17}
\end{equation*}
$$

it is clear that (1.14) admits a feasible symmetric positive semidefinite stabilizing solution if and only if $\gamma>\gamma_{f}$.

If this is the case, there exists a square $J$-spectral factor of $\Psi(z)$ having all zeros in the open unit disk and satisfying the stability conditions described in [3, Theorem 3.1], ensuring the existence of a proper stable filter $F(z)$.

On the other hand, it is also well known that there exists a value $\gamma_{0}$ such that for $\gamma<\gamma_{0}$, the ARE (1.14) does not admit any symmetric solution satisfying the feasibility condition and hence $\Psi(z)$ does not admit any minimal square $J$-spectral factor, namely, a $J$-spectral factor having the last possible McMillan degree. The constant $\gamma_{0}$ has the following interpretation from the $J$-spectral factorization point of view. A necessary condition for the existence of a square spectral factor is that $\Psi\left(e^{j \omega}\right)$ has $p$ positive eigenvalues and $l$ negative eigenvalues for all $\omega \in(-\pi, \pi]$. Since under the present assumptions $\Psi_{11}\left(e^{j \omega}\right)>0$ for all $\omega \in(-\pi, \pi]$, this condition is equivalent to the negative definiteness of the Schur complement

$$
\begin{equation*}
S(z):=\Psi_{22}(z)-\Psi_{12} \sim(z)\left[\Psi_{11}(z)\right]^{-1} \Psi_{12}(z) \tag{1.18}
\end{equation*}
$$

on the unit circle, i.e., to

$$
\begin{equation*}
S\left(e^{j \omega}\right)<0, \quad \omega \in(-\pi, \pi] \tag{1.19}
\end{equation*}
$$

By employing (1.11), we may write $S(z)$ as $S(z)=W(z)-\gamma^{2} I$ with

$$
\begin{align*}
W(z) & :=H_{2}(z) H_{2}^{\sim}(z)-H_{2}(z) H_{1}^{\sim}(z)\left[H_{1}(z) H_{1}^{\sim}(z)\right]^{-1} H_{1}(z) H_{2}^{\sim}(z) \\
& =H_{2}(z)\left[I-H_{1}^{\sim}(z)\left[H_{1}(z) H_{1}^{\sim}(z)\right]^{-1} H_{1}(z)\right] H_{2}^{\sim}(z) . \tag{1.20}
\end{align*}
$$

It is now easy to see that (1.19) is satified for $\gamma$ sufficiently large: $\gamma_{0}$ is the infimum of the values of $\gamma$ for which (1.19) is satisfied, i.e., $\gamma_{0}^{2}$ is the $\mathcal{L}_{\infty}$ norm of $W(z)$ :

$$
\begin{equation*}
\gamma_{0}^{2}=\|W(z)\|_{\infty} \tag{1.21}
\end{equation*}
$$

For a general overview on the connections between ARE and spectral and $J$-spectral factorization see $[11,12,17,1,15]$ and references therein.

Remark 1.1. From the estimation problem point of view, the constant $\gamma_{0}$ is the infimum of the values of $\gamma$ for which there exists a fixed-lag smoother (for some preview horizon $N$ ) achieving the attenuation level $\gamma$. To see this, first recall that under
the detectability assumption of the pair $(A, C)$, without any loss of generality, one can assume from the very beginning the stability of $H_{1}(z)$ and $H_{2}(z)$ and transform the given problem in a unilateral model matching problem in a new stable free filter parameter [21]. Now, let $V_{1}(z)$ be a square spectral factor of $H_{1}(z)$. A simple computation shows that

$$
\begin{equation*}
\left\|F(z) H_{1}(z)-H_{2}(z)\right\|_{\infty}^{2}=\left\|F(z) V_{1}(z)-H_{2}(z) H_{1}^{\sim}(z) V_{1}^{\sim}(z)^{-1}\right\|_{\infty}^{2}+\|W(z)\|_{\infty} \tag{1.22}
\end{equation*}
$$

Hence, the attenuation level cannot be lower than $\gamma_{0}$ and this infimum value is achieved by the acausal estimator

$$
\begin{equation*}
F_{\infty}(z)=H_{2}(z) H_{1}^{\sim}(z)\left(H_{1}(z) H_{1}^{\sim}(z)\right)^{-1}=\sum_{i=-\infty}^{\infty} F_{i} z^{i} \tag{1.23}
\end{equation*}
$$

Now the question arises of how to construct a stable fixed-lag smoother when an attenuation level $\gamma>\gamma_{0}$ is prescribed. This problem can be solved directly from expression (1.22) by taking a finite length expansion of the anticausal part of $F_{\infty}(z)$ given by (1.23). Indeed, let

$$
F_{N}(z)=F_{\infty}(z)-\sum_{i=N+1}^{\infty} F_{i} z^{i}
$$

If $F_{N}(z)$ is the ( $N$ steps acausal, yet stable) estimator so constructed, then

$$
\begin{aligned}
\left\|F_{N}(z) H_{1}(z)-H_{2}(z)\right\|_{\infty}^{2} & \leq \gamma_{0}^{2}+\left\|\sum_{i=N+1}^{\infty} F_{i} z^{i} V_{1}(z)\right\|_{\infty}^{2} \\
& \leq \gamma_{0}^{2}+\left(\sum_{i=N+1}^{\infty}\left\|F_{i}\right\|\right)^{2}\left\|H_{1}(z)\right\|_{\infty}^{2}
\end{aligned}
$$

can be rendered arbitrarily close to $\gamma_{0}^{2}$ by selecting a sufficiently large $N$. The estimator can be realized from the relevant factor following the results in [3] or directly from the solution of a suitable Riccati equation as described in [20].
1.3. Contribution of the paper. In this paper, we investigate the existence of a feasible stabilizing solution of (1.14) and the existence and the construction of a minimum phase square $J$-spectral factor $\Omega(z)$ for values of $\gamma$ in the interval $\left(\gamma_{0}, \gamma_{f}\right]$. Clearly, for such values of $\gamma$ the filter cannot exist, but there exists a fixed-lag smoother associated with a certain preview horizon length $N$ achieving the desired attenuation level. As shown in [5], the computation of the stabilizing solution $\Delta$ of (1.14) and of the corresponding $J$-spectral factor $\Omega(z)$ are crucial steps to obtain efficiently the minimum-lag smoothing filter achieving the desired attenuation level. Notably, the solution $\Delta$ can be directly used to initialize an iterative algorithm to work out a minimum-lag central smoother; see [13, 2].

The contribution of the present paper is to prove that for $\gamma \in\left(\gamma_{0}, \gamma_{f}\right]$, equation (1.14) still admits a stabilizing feasible solution except for a finite number (at most $2[n+\operatorname{rank}(M)])$ of values of $\gamma$. This result generalizes a result in [4] where it was assumed that $D D^{\top}=I, M=0$, and $D B^{\top}=0$. While removing the assumption $D D^{\top}=I$ impairs only slightly longer formulas, the presence of the matrix $M$ gives rise to a much more difficult problem. Also, the presence of the matrix $M$ implies that
having assumed (1.7e) we cannot assume $D B^{\top}=0$ (only one of the two conditions (1.7e) and $D B^{\top}=0$ can be assumed without loss of generality) and this adds some other technical difficulties to the problem. On the other hand the estimation problem becomes much more interesting and of practical importance if we consider the presence of the matrix $M$. Indeed, in this case, not only can we attack deconvolution problems for the importance of which we refer to [9] and references therein, but we can, more generally, address the problem of estimating an arbitrary linear combination of input and state vectors; see the discussion in [7] for the practical importance of the latter problem.
1.4. Paper organization. In section 2 we analyze some auxiliary spectral factorization problems and the corresponding AREs. In section 3 we show that for generic values of $\gamma$, the antistabilizing solution of an ARE considered in section 2 is nonsingular. Such preliminary results are employed in section 4 , where we prove our main result on the existence of a stabilizing feasible solution of (1.14). In section 5 we explicitly derive a minimum phase $J$-spectral factor of $\Psi(z)$. In section 6 we analyze some of the issues arising in correspondence of the values of $\gamma$ for which we cannot guarantee the existence of the stabilizing solution of (1.14) and we discuss some peculiar aspects of the $J$-spectral factorization.
2. Preliminary spectral factorizations and auxiliary AREs. In this section we compute a state-space realization of a spectral factor $T(z)$ of $W(z)$ defined in (1.20) and a spectral factorization of $-S(z)=\gamma^{2} I-W(z)$. These factorizations are related to the stabilizing solutions of a pair of coupled AREs. These solutions will be used to compute the solution (and hence to prove constructively its existence) of a third ARE that is strictly related to (1.14). This procedure will be carried over through various steps.
2.1. Computation of a square spectral factor of $\boldsymbol{H}_{1}(\boldsymbol{z}) \boldsymbol{H}_{1}{ }^{\sim}(\boldsymbol{z})$. Consider the following standard filtering ARE:

$$
\begin{equation*}
P=A P A^{\top}+B B^{\top}-\left(A P C^{\top}+B D^{\top}\right)\left(D D^{\top}+C P C^{\top}\right)^{-1}\left(C P A^{\top}+D B^{\top}\right) \tag{2.1}
\end{equation*}
$$

By assumptions (1.7b) and (1.7c), it admits a positive definite stabilizing solution $P_{s}=P_{s}^{\top}>0$. Such a solution corresponds to spectral factorization of $H_{1}(z) H_{1}^{\sim}(z)$. In fact, by defining

$$
\begin{equation*}
D_{1}:=\left(D D^{\top}+C P_{s} C^{\top}\right)^{1 / 2}, \quad G_{1}:=\left(A P_{s} C^{\top}+B D^{\top}\right) D_{1}^{-1} \tag{2.2}
\end{equation*}
$$

it is easy to check (see [12]) that

$$
\begin{equation*}
H_{1}(z) H_{1}^{\sim}(z)=V_{1}(z) V_{1}^{\sim}(z) \tag{2.3}
\end{equation*}
$$

with

$$
\begin{equation*}
V_{1}(z):=C(z I-A)^{-1} G_{1}+D_{1} . \tag{2.4}
\end{equation*}
$$

2.2. Computation of $V_{1}(z)^{-1} H_{1}(z)$. First notice that

$$
\begin{equation*}
V_{1}(z)^{-1}=D_{1}^{-1}-D_{1}^{-1} C(z I-\Gamma)^{-1} G_{1} D_{1}^{-1} \tag{2.5}
\end{equation*}
$$

with

$$
\begin{equation*}
\Gamma:=A-G_{1} D_{1}^{-1} C \tag{2.6}
\end{equation*}
$$

being the closed-loop matrix corresponding to the stabilizing solution $P_{s}$ of the ARE (2.1) (so that all the eigenvalues of $\Gamma$ lay in the open unit disk).

Then we have

$$
\begin{aligned}
V_{2}(z):= & V_{1}(z)^{-1} H_{1}(z) \\
= & D_{1}^{-1} C(z I-A)^{-1} B-D_{1}^{D_{1}^{-1} C(z I-\Gamma)^{-1} G_{1} D_{1}^{-1} D+D_{1}^{-1} D} \\
& -D_{1}^{-1} C(z I-\Gamma)^{-1} \underbrace{G_{1} D_{1}^{-1} C}_{(z I-\Gamma)-(z I-A)}(z I-A)^{-1} B \\
= & D_{1}^{-1} C(z I-\Gamma)^{-1}\left(B-G_{1} D_{1}^{-1} D\right)+D_{1}^{-1} D \\
= & C_{2}(z I-\Gamma)^{-1} G_{2}+D_{2}
\end{aligned}
$$

with

$$
\begin{equation*}
C_{2}:=D_{1}^{-1} C, \quad G_{2}:=B-G_{1} D_{1}^{-1} D, \quad D_{2}:=D_{1}^{-1} D \tag{2.8}
\end{equation*}
$$

2.3. Computation of $\boldsymbol{V}_{2}(z) \boldsymbol{H}_{2}{ }^{\sim}(z)$. We have

$$
\begin{align*}
V_{2}(z) H_{2}^{\sim}(z)= & {\left[C_{2}(z I-\Gamma)^{-1} G_{2}+D_{2}\right]\left[B^{\top}\left(z^{-1} I-A^{\top}\right)^{-1} L^{\top}+M^{\top}\right] } \\
= & C_{2}(z I-\Gamma)^{-1} G_{2} M^{\top}+D_{2} B^{\top}\left(z^{-1} I-A^{\top}\right)^{-1} L^{\top}+D_{2} M^{\top} \\
& +C_{2}(z I-\Gamma)^{-1} G_{2} B^{\top}\left(z^{-1} I-A^{\top}\right)^{-1} L^{\top} \tag{2.9}
\end{align*}
$$

Taking the ARE (2.1) into account, it is easy to check that

$$
\begin{equation*}
G_{2} B^{\top}=P_{s}-\Gamma P_{s} A^{\top} \tag{2.10}
\end{equation*}
$$

which may be easily rewritten in the form

$$
\begin{equation*}
G_{2} B^{\top}=(z I-\Gamma) P_{s}\left(z^{-1} I-A^{\top}\right)+(z I-\Gamma) P_{s} A^{\top}+\Gamma P_{s}\left(z^{-1} I-A^{\top}\right) \tag{2.11}
\end{equation*}
$$

Plugging this expression in (2.9) we get

$$
\begin{equation*}
V_{2}(z) H_{2}^{\sim}(z)=C_{2}(z I-\Gamma)^{-1} G_{3}+G_{1}^{\top}\left(z^{-1} I-A^{\top}\right)^{-1} L^{\top}+D_{3} \tag{2.12}
\end{equation*}
$$

with $G_{3}:=\Gamma P_{s} L^{\top}+G_{2} M^{\top}$ and $D_{3}:=C_{2} P_{s} L^{\top}+D_{2} M^{\top}$ or, taking into account that $D M^{\top}=0$,

$$
\begin{equation*}
G_{3}=\Gamma P_{s} L^{\top}+B M^{\top}, \quad D_{3}=C_{2} P_{s} L^{\top} \tag{2.13}
\end{equation*}
$$

2.4. Computation of $H_{2}(z) H_{1}{ }^{\sim}(z)\left[H_{1}(z) H_{1}{ }^{\sim}(z)\right]^{-1} H_{1}(z) H_{2}{ }^{\sim}(z)$. It is clear that $H_{2}(z) H_{1}^{\sim}(z)\left[H_{1}(z) H_{1}^{\sim}(z)\right]^{-1} H_{1}(z) H_{2}^{\sim}(z)=H_{2}(z) V_{2}^{\sim}(z) V_{2}(z) H_{2}^{\sim}(z)$. We have

$$
\begin{align*}
H_{2}(z) V_{2}^{\sim}(z) V_{2}(z) H_{2}^{\sim}(z)= & D_{3}^{\top} D_{3}+G_{3}^{\top}\left(z^{-1} I-\Gamma^{\top}\right)^{-1} C_{2}^{\top} C_{2}(z I-\Gamma)^{-1} G_{3} \\
& +D_{3}^{\top} C_{2}(z I-\Gamma)^{-1} G_{3}+D_{3}^{\top} G_{1}^{\top}\left(z^{-1} I-A^{\top}\right)^{-1} L^{\top} \\
& +\left[D_{3}^{\top} C_{2}(z I-\Gamma)^{-1} G_{3}+D_{3}^{\top} G_{1}^{\top}\left(z^{-1} I-A^{\top}\right)^{-1} L^{\top}\right]^{\sim} \\
& +L(z I-A)^{-1} G_{1} C_{2}(z I-\Gamma)^{-1} G_{3} \\
& +\left[L(z I-A)^{-1} G_{1} C_{2}(z I-\Gamma)^{-1} G_{3}\right]^{\sim} \\
& +L(z I-A)^{-1} G_{1} G_{1}^{\top}\left(z^{-1} I-A^{\top}\right)^{-1} L^{\top} . \tag{2.14}
\end{align*}
$$

The last three terms of such expression may be expanded as follows: first taking into account that

$$
\begin{equation*}
G_{1} C_{2}=G_{1} D_{1}^{-1} C=A-\Gamma=(z I-\Gamma)-(z I-A) \tag{2.15}
\end{equation*}
$$

we have

$$
\begin{equation*}
L(z I-A)^{-1} G_{1} C_{2}(z I-\Gamma)^{-1} G_{3}=L(z I-A)^{-1} G_{3}-L(z I-\Gamma)^{-1} G_{3} \tag{2.16}
\end{equation*}
$$

Second, taking into account the ARE (2.1), we have

$$
\begin{align*}
G_{1} G_{1}^{\top}= & A P_{s} A^{\top}-P_{s}+B B^{\top}=B B^{\top}-(z I-A) P_{s}\left(z^{-1} I-A^{\top}\right)-(z I-A) P_{s} A^{\top} \\
& -A P_{s}\left(z^{-1} I-A^{\top}\right) \tag{2.17}
\end{align*}
$$

so that

$$
\begin{align*}
& L(z I-A)^{-1} G_{1} G_{1}^{\top}\left(z^{-1} I-A^{\top}\right)^{-1} L^{\top} \\
& =  \tag{2.18}\\
& =L(z I-A)^{-1} B B^{\top}\left(z^{-1} I-A^{\top}\right)^{-1} L^{\top}-L P_{s} L^{\top} \\
& \\
& \quad-L(z I-A)^{-1} A P_{s} L^{\top}-L P_{s} A^{\top}\left(z^{-1} I-A^{\top}\right)^{-1} L^{\top} .
\end{align*}
$$

2.5. Computation of $\boldsymbol{W}(\boldsymbol{z})$. From the definition (1.20) of $W(z)$ it follows that

$$
\begin{equation*}
W(z)=H_{2}(z) H_{2}^{\sim}(z)-H_{2}(z) V_{2}^{\sim}(z) V_{2}(z) H_{2}^{\sim}(z) . \tag{2.19}
\end{equation*}
$$

Moreover, by computing $H_{2}(z) H_{2}^{\sim}(z)$, we get the following expression:

$$
\begin{align*}
H_{2}(z) H_{2} \sim(z)= & L(z I-A)^{-1} B B^{\top}\left(z^{-1} I-A^{\top}\right)^{-1} L^{\top}+M M^{\top} \\
& +L(z I-A)^{-1} B M^{\top}+\left[L(z I-A)^{-1} B M^{\top}\right]^{\sim} \tag{2.20}
\end{align*}
$$

By subtracting (2.14) from (2.20), taking (2.16), (2.18), and (2.19) into account, we get

$$
\begin{align*}
W(z)= & \left(L-D_{3}^{\top} C_{2}\right)(z I-\Gamma)^{-1} G_{3}+\left[\left(L-D_{3}^{\top} C_{2}\right)(z I-\Gamma)^{-1} G_{3}\right]^{\top} \\
& -G_{3}^{\top}\left(z^{-1} I-\Gamma^{\top}\right)^{-1} C_{2}^{\top} C_{2}(z I-\Gamma)^{-1} G_{3}+L P_{s} L^{\top}+M M^{\top}-D_{3}^{\top} D_{3} \\
& +L(z I-A)^{-1}\left(B M^{\top}-G_{1} D_{3}+A P_{s} L^{\top}-G_{3}\right) \\
& +\left[L(z I-A)^{-1}\left(B M^{\top}-G_{1} D_{3}+A P_{s} L^{\top}-G_{3}\right)\right] . \tag{2.21}
\end{align*}
$$

We now prove that $B M^{\top}-G_{1} D_{3}+A P_{s} L^{\top}-G_{3}=0$, so that the last two terms of the latter expression vanish. We have

$$
\begin{align*}
& B M^{\top}-G_{1} D_{3}+A P_{s} L^{\top}-G_{3} \\
& \quad=B M^{\top}-\underbrace{G_{1} C_{2}}_{A-\Gamma} P_{s} L^{\top}+A P_{s} L^{\top}-\Gamma P_{s} L^{\top}-B M^{\top}=0 . \tag{2.22}
\end{align*}
$$

2.6. Spectral factorization of $\boldsymbol{W}(\boldsymbol{z})$. To factorize $W(z)$ we first need to rewrite the term

$$
G_{3}^{\top}\left(z^{-1} I-\Gamma^{\top}\right)^{-1} C_{2}^{\top} C_{2}(z I-\Gamma)^{-1} G_{3}
$$

To this aim we derive a new expression for $C_{2}^{\top} C_{2}$. Standard computations show that the solution $P_{s}$ of the ARE (2.1) satisfies the following identity [15, p. 271]:

$$
\begin{equation*}
P_{s}=F P_{s} F^{\top}-F P_{s} C^{\top}\left(C P_{s} C^{\top}+D D^{\top}\right)^{-1} C P_{s} F^{\top}+B_{1} B_{1}^{\top}, \tag{2.23}
\end{equation*}
$$

where $F$ and $B_{1}$ are defined in (1.6). We also have

$$
\begin{align*}
\Gamma & =A-G_{1} D_{1}^{-1} C=F+\left[B D^{\top}\left(D D^{\top}\right)^{-1} D_{1}^{2}-A P_{s} C^{\top}-B D^{\top}\right] D_{1}^{-1} C_{2} \\
& =F-F P_{s} C^{\top} D_{1}^{-1} C_{2}=F-F P_{s} C_{2}^{\top} C_{2} \\
& =F\left(I-P_{s} C_{2}^{\top} C_{2}\right) \tag{2.24}
\end{align*}
$$

so that (2.23) may be rewritten in the form

$$
\begin{equation*}
P_{s}-B_{1} B_{1}^{\top}=\Gamma P_{s} F^{\top}=F P_{s} \Gamma^{\top} \tag{2.25}
\end{equation*}
$$

From $D D^{\top}>0$ and $P_{s}>0$ it easily follows that $\left(I-P_{s} C_{2}^{\top} C_{2}\right)=P_{s}\left(P_{s}^{-1}-\right.$ $\left.C^{\top}\left(C P_{s} C^{\top}+D D^{\top}\right)^{-1} C\right)$ is nonsingular (in fact, $\left(P_{s}^{-1}-C^{\top}\left(C P_{s} C^{\top}+D D^{\top}\right)^{-1} C\right)>$ 0 ). Thus $\Gamma$, and hence $F P_{s} \Gamma^{\top}$, are nonsingular so that the left-hand side of (2.25) is nonsingular as well, and we have

$$
\begin{equation*}
\left(P_{s}-B_{1} B_{1}^{\top}\right)^{-1}=\Gamma^{-\top} P_{s}^{-1} F^{-1} \tag{2.26}
\end{equation*}
$$

which may be rewritten as

$$
\begin{equation*}
P_{s}^{-1}=\Gamma^{\top}\left(P_{s}-B_{1} B_{1}^{\top}\right)^{-1} F, \tag{2.27}
\end{equation*}
$$

and, by multiplying on the right side by $\left(I-P_{s} C_{2}^{\top} C_{2}\right)$,

$$
\begin{equation*}
P_{s}^{-1}\left(I-P_{s} C_{2}^{\top} C_{2}\right)=\Gamma^{\top}\left(P_{s}-B_{1} B_{1}^{\top}\right)^{-1} \Gamma \tag{2.28}
\end{equation*}
$$

Therefore
(2.29a) $C_{2}^{\top} C_{2}=P_{s}^{-1}-\Gamma^{\top}\left(P_{s}-B_{1} B_{1}^{\top}\right)^{-1} \Gamma$

$$
\begin{align*}
= & P_{s}^{-1}-\Gamma^{\top} P_{s}^{-1} \Gamma-\Gamma^{\top} P_{s}^{-1} B_{1} V^{-1} B_{1}^{\top} P_{s}^{-1} \Gamma  \tag{2.29b}\\
= & \left(z^{-1} I-\Gamma^{\top}\right) P_{s}^{-1}(z I-\Gamma)+\Gamma^{\top} P_{s}^{-1}(z I-\Gamma)+\left(z^{-1} I-\Gamma^{\top}\right) P_{s}^{-1} \Gamma \\
& -\Gamma^{\top} P_{s}^{-1} B_{1} V^{-1} B_{1}^{\top} P_{s}^{-1} \Gamma \tag{2.29c}
\end{align*}
$$

where

$$
\begin{equation*}
V:=I-B_{1}^{\top} P_{s}^{-1} B_{1} \tag{2.30}
\end{equation*}
$$

Notice that
$P_{s}^{-1}-C_{2}^{\top} C_{2}=P_{s}^{-1}-C^{\top}\left(D D^{\top}+C P_{s} C^{\top}\right)^{-1} C=P_{s}^{-1}\left[P_{s}^{-1}+C^{\top}\left(D D^{\top}\right)^{-1} C\right]^{-1} P_{s}^{-1}>0$ (2.31)
so that from (2.29a) it follows that $P_{s}-B_{1} B_{1}^{\top}>0$, which, together with the positive definiteness of $P_{s}$, implies that $V$ is positive definite as well.

Plugging in the expression in the right-hand side of $(2.29 \mathrm{c})$ in place of $C_{2}^{\top} C_{2}$ in the third term of (2.21), we get

$$
\begin{align*}
W(z)= & \left(L-D_{3}^{\top} C_{2}-G_{3}^{\top} P_{s}^{-1} \Gamma\right)(z I-\Gamma)^{-1} G_{3} \\
& +\left[\left(L-D_{3}^{\top} C_{2}-G_{3}^{\top} P_{s}^{-1} \Gamma\right)(z I-\Gamma)^{-1} G_{3}\right] \\
& +G_{3}^{\top}\left(z^{-1} I-\Gamma^{\top}\right)^{-1} \Gamma^{\top} P_{s}^{-1} B_{1} V^{-1} B_{1}^{\top} P_{s}^{-1} \Gamma(z I-\Gamma)^{-1} G_{3}  \tag{2.32}\\
& +L P_{s} L^{\top}+M M^{\top}-D_{3}^{\top} D_{3}-G_{3}^{\top} P_{s}^{-1} G_{3} .
\end{align*}
$$

Now we are ready to prove that the following spectral factorization of $W(z)$ holds:

$$
\begin{equation*}
W(z)=T(z) T^{\sim}(z) \tag{2.33}
\end{equation*}
$$

where

$$
\begin{equation*}
T(z):=\left[-G_{3}^{\top}\left(z^{-1} I-\Gamma^{\top}\right)^{-1} \Gamma^{\top} P_{s}^{-1} B_{1}-G_{3}^{\top} P_{s}^{-1} B_{1}+M\right] V^{-1 / 2} \tag{2.34}
\end{equation*}
$$

To this aim it is sufficient to show that

$$
\begin{align*}
\left(M-G_{3}^{\top} P_{s}^{-1} B_{1}\right) V^{-1} & \left(M-G_{3}^{\top} P_{s}^{-1} B_{1}\right)^{\top} \\
& \quad-\left(L P_{s} L^{\top}+M M^{\top}-D_{3}^{\top} D_{3}-G_{3}^{\top} P_{s}^{-1} G_{3}\right)=0 \tag{2.35}
\end{align*}
$$

and

$$
\begin{equation*}
\Gamma^{\top} P_{s}^{-1} B_{1} V^{-1}\left(M-G_{3}^{\top} P_{s}^{-1} B_{1}\right)^{\top}+\left(L-D_{3}^{\top} C_{2}-G_{3}^{\top} P_{s}^{-1} \Gamma\right)^{\top}=0 \tag{2.36}
\end{equation*}
$$

As for (2.35) notice that

$$
\begin{equation*}
M-G_{3}^{\top} P_{s}^{-1} B_{1}=M-L P_{s} \Gamma^{\top} P_{s}^{-1} B_{1}-M B_{1}^{\top} P_{s}^{-1} B_{1}=M V-L P_{s} \Gamma^{\top} P_{s}^{-1} B_{1} \tag{2.37}
\end{equation*}
$$

so that, taking (1.8) into account, we may expand the left-hand side of (2.35) as

$$
\begin{align*}
& M V M^{\top}-M B_{1}^{\top} P_{s}^{-1} \Gamma P_{s} L^{\top}-\left(M B_{1}^{\top} P_{s}^{-1} \Gamma P_{s} L^{\top}\right)^{\top} \\
& +L P_{s} \Gamma^{\top} P_{s}^{-1} B_{1} V^{-1} B_{1}^{\top} P_{s}^{-1} \Gamma P_{s} L^{\top}-L P_{s} L^{\top}-M M^{\top} \\
& +L P_{s} C_{2}^{\top} C_{2} P_{s} L^{\top}+L P_{s} \Gamma^{\top} P_{s}^{-1} \Gamma P_{s} L^{\top}+M B_{1}^{\top} P_{s}^{-1} B_{1} M^{\top}  \tag{2.38}\\
& +M B_{1}^{\top} P_{s}^{-1} \Gamma P_{s} L^{\top}+\left(M B_{1}^{\top} P_{s}^{-1} \Gamma P_{s} L^{\top}\right)^{\top} \\
& =L P_{s}\left(\Gamma^{\top} P_{s}^{-1} B_{1} V^{-1} B_{1}^{\top} P_{s}^{-1} \Gamma-P_{s}^{-1}+C_{2}^{\top} C_{2}+\Gamma^{\top} P_{s}^{-1} \Gamma\right) P_{s} L^{\top}=0
\end{align*}
$$

where the latter equality follows from the expression (2.29b) for $C_{2}^{\top} C_{2}$.
Similarly, we may expand the left-hand side of (2.36) as

$$
\begin{align*}
& L^{\top}-C_{2}^{\top} C_{2} P_{s} L^{\top}-\Gamma^{\top} P_{s}^{-1} \Gamma P_{s} L^{\top}-\Gamma^{\top} P_{s}^{-1} B_{1} V^{-1} B_{1}^{\top} P_{s}^{-1} \Gamma P_{s} L^{\top} \\
& =\left(P_{s}^{-1}-C_{2}^{\top} C_{2}-\Gamma^{\top} P_{s}^{-1} \Gamma-\Gamma^{\top} P_{s}^{-1} B_{1} V^{-1} B_{1}^{\top} P_{s}^{-1} \Gamma\right) P_{s} L^{\top}=0 \tag{2.39}
\end{align*}
$$

We now rewrite $T(z)$ in a more convenient form. Notice that

$$
\left(z^{-1} I-\Gamma^{\top}\right)^{-1}=-\Gamma^{-\top}-\Gamma^{-\top}\left(z I-\Gamma^{-\top}\right)^{-1} \Gamma^{-\top}
$$

so that

$$
\begin{align*}
T(z) & =\left[G_{3}^{\top} \Gamma^{-\top}\left(z I-\Gamma^{-\top}\right)^{-1} P_{s}^{-1} B_{1}+M\right] V^{-1 / 2} \\
& =\left[G_{3}^{\top} \Gamma^{-\top} P_{s}^{-1}\left(z I-P_{s} \Gamma^{-\top} P_{s}^{-1}\right)^{-1} B_{1}+M\right] V^{-1 / 2} \\
& =L_{1}\left(z I-F_{a}\right)^{-1} B_{1} V^{-1 / 2}+M V^{-1 / 2} \tag{2.40}
\end{align*}
$$

with

$$
\begin{equation*}
L_{1}:=G_{3}^{\top} \Gamma^{-\top} P_{s}^{-1}=L+M B_{1}^{\top} \Gamma^{-\top} P_{s}^{-1}, \quad F_{a}:=P_{s} \Gamma^{-\top} P_{s}^{-1} \tag{2.41}
\end{equation*}
$$

Notice that since $\Gamma$ is stable, $F_{a}$ is antistable. As a direct consequence of (2.26) we have the following relation that will be useful in what follows:

$$
\begin{equation*}
F_{a}=\left(I-B_{1} B_{1}^{\top} P_{s}^{-1}\right)^{-1} F \tag{2.42}
\end{equation*}
$$

2.7. Spectral factorization of $-\boldsymbol{S}(\boldsymbol{z})$. We have

$$
-S(z)=\gamma^{2} I-W(z)=\gamma^{2} I-T(z) T^{\sim}(z)
$$

As we have done for the factorization in section 2.1 consider the following ARE for the spectral factorization of $-S(z)$ :

$$
\begin{equation*}
X=F_{a} X F_{a}^{\top}-B_{1} V^{-1} B_{1}^{\top}-\left(F_{a} X L_{1}^{\top}-B_{1} V^{-1} M^{\top}\right)\left(R+L_{1} X L_{1}^{\top}\right)^{-1}\left(L_{1} X F_{a}^{\top}-M V^{-1} B_{1}^{\top}\right) \tag{2.43}
\end{equation*}
$$

with

$$
\begin{equation*}
R:=\gamma^{2} I-M V^{-1} M^{\top} \tag{2.44}
\end{equation*}
$$

In view of (1.21), we have $\|T(z)\|_{\infty}<\gamma$ for any $\gamma>\gamma_{0}$ so that as a consequence of a generalized version of the discrete-time bounded real lemma (see [8, Theorem 2.1]), (2.43) admits a stabilizing solution, namely, a solution $X_{s}=X_{s}^{\top}$ such that

$$
\begin{equation*}
\Gamma_{a}:=F_{a}-\left(F_{a} X_{s} L_{1}^{\top}-B_{1} V^{-1} M^{\top}\right)\left(R+L_{1} X_{s} L_{1}^{\top}\right)^{-1} L_{1} \tag{2.45}
\end{equation*}
$$

is a stability matrix; moreover $\left(R+L_{1} X_{s} L_{1}^{\top}\right)>0$ and the function

$$
\begin{equation*}
T_{1}(z):=\left[I+L_{1}\left(z I-F_{a}\right)^{-1}\left(F_{a} X_{s} L_{1}^{\top}-B_{1} V^{-1} M^{\top}\right)\left(R+L_{1} X_{s} L_{1}^{\top}\right)^{-1}\right]\left(R+L_{1} X_{s} L_{1}^{\top}\right)^{1 / 2} \tag{2.46}
\end{equation*}
$$

is a square spectral factor of $-S(z)=\gamma^{2} I_{l}-W(z)$, namely, it is a square matrix function such that $-S(z)=T_{1}(z) T_{1}^{\sim}(z)$. Moreover, the numerator matrix of $T_{1}(z)$ is given by (2.45) and hence $\left[T_{1}(z)\right]^{-1}$ is stable.

We now show that $X_{s}$ is positive definite. In fact, since $\left(R+L_{1} X_{s} L_{1}^{\top}\right)>0$, we have

$$
\begin{equation*}
X_{s} \leq F_{a} X_{s} F_{a}^{\top}-B_{1} V^{-1} B_{1}^{\top} \tag{2.47}
\end{equation*}
$$

But $\left(F, B_{1}\right)$ is, by assumption, reachable, $\left(F_{a}, B_{1}\right)$ is such (because $F_{a}$ is obtained from $F$ by state feedback as it is apparent from (2.42)), and then also ( $F_{a}, B_{1} V^{-1 / 2}$ ) is reachable. Therefore, since $F_{a}$ is antistable, a standard Lyapunov argument, shows that $X_{s}>0$.
2.8. The ARE for $P_{s}^{-1}-X_{s}^{-1}$. Define

$$
\begin{equation*}
R_{1}:=I-\frac{M^{\top} M}{\gamma^{2}} \tag{2.48}
\end{equation*}
$$

and

$$
\begin{equation*}
F_{y}:=A-B J_{0}^{\prime} J_{1}^{-1} H_{0} \tag{2.49}
\end{equation*}
$$

where $J_{0}, J_{1}, H_{0}$ are defined in (1.13). Notice that $F_{y}$ is well defined if and only if $J_{1}$ is nonsingular or, equivalently, if and only if $R_{1}$ is nonsingular. Let
(2.50) $\mathcal{G}_{1}:=\left\{\gamma>\gamma_{0}\right.$ : at least one of the matrices $R, R_{1}$, and $F_{y}$, is singular $\}$.

Remark 2.1. Notice that since $F_{y}$ may be written in the form

$$
\begin{equation*}
F_{y}=F-B_{1} M^{\top}\left(M M^{\top}-\gamma^{2} I\right)^{-1} L \tag{2.51}
\end{equation*}
$$

where $F$ is nonsingular, $\mathcal{G}_{1}$ contains a finite number of points. More precisely, it contains, at most, $n+2 \operatorname{rank}(M)$ points.

We consider now the following ARE:

$$
\begin{align*}
Y= & F^{\top} Y F+\left(F^{\top} Y B_{1}+\frac{L^{\top} M}{\gamma^{2}}\right)\left(R_{1}-B_{1}^{\top} Y B_{1}\right)^{-1}\left(B_{1}^{\top} Y F+\frac{M^{\top} L}{\gamma^{2}}\right)  \tag{2.52}\\
& -C^{\top}\left(D D^{\top}\right)^{-1} C+\frac{L^{\top} L}{\gamma^{2}}
\end{align*}
$$

and prove that it admits a solution that can be explicitely computed. To this aim we need a technical lemma.

Lemma 2.1. Let $\gamma>\gamma_{0}$ and $\gamma \notin \mathcal{G}_{1}$ and define

$$
\begin{equation*}
F_{1}:=F_{a}+B_{1} V^{-1} M^{\top} R^{-1} L_{1} \tag{2.53}
\end{equation*}
$$

Then the following relation holds:

$$
\begin{equation*}
X_{s}^{-1}=F_{1}^{\top} X_{s}^{-1}\left(I-B_{1} R_{1}^{-1} B_{1}^{\top}\left(P_{s}^{-1}-X_{s}^{-1}\right)\right)^{-1} F_{y}-L_{1}^{\top} R^{-1} L_{1} \tag{2.54}
\end{equation*}
$$

Proof. First we establish the following useful identity:

$$
\begin{equation*}
F_{y}=\left(I-B_{1} R_{1}^{-1} B_{1}^{\top} P_{s}^{-1}\right) F_{1} . \tag{2.55}
\end{equation*}
$$

To show identity (2.55), we write $F_{1}$ in the form

$$
\begin{aligned}
F_{1} & =\left(I+B_{1} V^{-1} M^{\top} R^{-1} M B_{1}^{\top} P_{s}^{-1}\right) F_{a}+B_{1} V^{-1} M^{\top} R^{-1} L \\
(2.56) & =\left(I+B_{1} V^{-1} M^{\top} R^{-1} M B_{1}^{\top} P_{s}^{-1}\right)\left(I-B_{1} B_{1}^{\top} P_{s}^{-1}\right)^{-1} F+B_{1} V^{-1} M^{\top} R^{-1} L
\end{aligned}
$$

which gives

$$
\begin{aligned}
&\left(I-B_{1} R_{1}^{-1} B_{1}^{\top} P_{s}^{-1}\right) F_{1} \\
&=\left(I-B_{1} R_{1}^{-1} B_{1}^{\top} P_{s}^{-1}\right)\left(I+B_{1} V^{-1} M^{\top} R^{-1} M B_{1}^{\top} P_{s}^{-1}\right)\left(I-B_{1} B_{1}^{\top} P_{s}^{-1}\right)^{-1} F \\
&(2.57) \quad+\left(I-B_{1} R_{1}^{-1} B_{1}^{\top} P_{s}^{-1}\right) B_{1} V^{-1} M^{\top} R^{-1} L
\end{aligned}
$$

so that, taking (2.51) into account, it is sufficient to prove that

$$
\begin{equation*}
\left(I-B_{1} R_{1}^{-1} B_{1}^{\top} P_{s}^{-1}\right)\left(I+B_{1} V^{-1} M^{\top} R^{-1} M B_{1}^{\top} P_{s}^{-1}\right)\left(I-B_{1} B_{1}^{\top} P_{s}^{-1}\right)^{-1}=I \tag{2.58}
\end{equation*}
$$

and

$$
\begin{equation*}
\left(I-B_{1} R_{1}^{-1} B_{1}^{\top} P_{s}^{-1}\right) B_{1} V^{-1} M^{\top} R^{-1}=-B_{1} M^{\top}\left(M M^{\top}-\gamma^{2} I\right)^{-1} \tag{2.59}
\end{equation*}
$$

As for (2.58), we have

$$
\begin{align*}
& \left(I-B_{1} R_{1}^{-1} B_{1}^{\top} P_{s}^{-1}\right)\left(I+B_{1} V^{-1} M^{\top} R^{-1} M B_{1}^{\top} P_{s}^{-1}\right)\left(I-B_{1} B_{1}^{\top} P_{s}^{-1}\right)^{-1} \\
= & \left(I-B_{1} R_{1}^{-1} B_{1}^{\top} P_{s}^{-1}\right)\left(I-B_{1} B_{1}^{\top} P_{s}^{-1}+B_{1}\left(I+V^{-1} M^{\top} R^{-1} M\right) B_{1}^{\top} P_{s}^{-1}\right)\left(I-B_{1} B_{1}^{\top} P_{s}^{-1}\right)^{-1} \\
= & \left(I-B_{1} R_{1}^{-1} B_{1}^{\top} P_{s}^{-1}\right)\left[I-B_{1} B_{1}^{\top} P_{s}^{-1}+B_{1}\left(V-\frac{M^{\top} M}{\gamma^{2}}\right)^{-1} V B_{1}^{\top} P_{s}^{-1}\right]\left(I-B_{1} B_{1}^{\top} P_{s}^{-1}\right)^{-1} \\
= & \left(I-B_{1} R_{1}^{-1} B_{1}^{\top} P_{s}^{-1}\right)\left[I+B_{1}\left(V-\frac{M^{\top} M}{\gamma^{2}}\right)^{-1} V B_{1}^{\top} P_{s}^{-1}\left(I-B_{1} B_{1}^{\top} P_{s}^{-1}\right)^{-1}\right] \\
= & \left(I-B_{1} R_{1}^{-1} B_{1}^{\top} P_{s}^{-1}\right)\left[I+B_{1}\left(V-\frac{M^{\top} M}{\gamma^{2}}\right)^{-1} B_{1}^{\top} P_{s}^{-1}\right] \\
= & \left(I-B_{1} R_{1}^{-1} B_{1}^{\top} P_{s}^{-1}\right)\left(I-B_{1} R_{1}^{-1} B_{1}^{\top} P_{s}^{-1}\right)^{-1}=I . \tag{2.60}
\end{align*}
$$

As for (2.59), we have

$$
\begin{align*}
\left(I-B_{1} R_{1}^{-1} B_{1}^{\top} P_{s}^{-1}\right) B_{1} V^{-1} M^{\top} R^{-1} & =B_{1}\left(I-R_{1}^{-1} B_{1}^{\top} P_{s}^{-1} B_{1}\right) V^{-1} M^{\top} R^{-1} \\
& =B_{1} R_{1}^{-1}(\underbrace{R_{1}-B_{1}^{\top} P_{s}^{-1} B_{1}}_{V-\frac{M^{\top} M}{\gamma^{2}}}) V^{-1} M^{\top} R^{-1} \\
& =B_{1} R_{1}^{-1} M^{\top}\left(I-\frac{M V^{-1} M^{\top}}{\gamma^{2}}\right) R^{-1} \\
& =-B_{1} M^{\top}\left(M M^{\top}-\gamma^{2} I\right)^{-1} . \tag{2.61}
\end{align*}
$$

Since $X_{s}$ is a solution of the ARE (2.43), as a direct consequence of standard equivalence of Riccati equations [15, p. 271], we have the following identity:

$$
\begin{equation*}
X_{s}=F_{1} X_{s} F_{1}^{\top}-F_{1} X_{s} L_{1}^{\top}\left(L_{1} X_{s} L_{1}^{\top}+R\right)^{-1} L_{1} X_{s} F_{1}^{\top}-B_{1}\left(V-\frac{M^{\top} M}{\gamma^{2}}\right)^{-1} B_{1}^{\top} \tag{2.62}
\end{equation*}
$$

with $F_{1}$ being defined in (2.53). The closed-loop matrix may thus be written in the form

$$
\begin{equation*}
\Gamma_{a}=F_{1}-F_{1} X_{s} L_{1}^{\top}\left(L_{1} X_{s} L_{1}^{\top}+R\right)^{-1} L_{1}=F_{1}\left(I+X_{s} L_{1}^{\top} R^{-1} L_{1}\right)^{-1} \tag{2.63}
\end{equation*}
$$

Notice that $V-\frac{M^{\top} M}{\gamma^{2}}$ is nonsingular because $V$ and $R$ are such. In the same way, $X_{s}^{-1}+L_{1}^{\top} R^{-1} L_{1}$, and hence ( $I+X_{s} L_{1}^{\top} R^{-1} L_{1}$ ), are nonsingular because $X_{s}, L_{1} X_{s} L_{1}^{\top}+$ $R$, and $R$ are such. From (2.62) we easily get

$$
\begin{equation*}
X_{s}+B_{1}\left(V-\frac{M^{\top} M}{\gamma^{2}}\right)^{-1} B_{1}^{\top}=F_{1}\left(X_{s}-X_{s} L_{1}^{\top}\left(L_{1} X_{s} L_{1}^{\top}+R\right)^{-1} L_{1} X_{s}\right) F_{1}^{\top} \tag{2.64}
\end{equation*}
$$

As an immediate consequence of (2.55) we have that $F_{1}$ is nonsingular, so that the right-hand side of (2.64) (which may be written in the form $\left.F_{1}\left(I+X_{s} L_{1}^{\top} R^{-1} L_{1}\right)^{-1} F_{1}^{\top}\right)$ is nonsingular. Hence $X_{s}+B_{1}\left(V-\frac{M^{\top} M}{\gamma^{2}}\right)^{-1} B_{1}^{\top}$ is nonsingular as well and then $\frac{M^{\top} M}{\gamma^{2}}-V-B_{1}^{\top} X_{s}^{-1} B_{1}=-\left(R_{1}-B_{1}^{\top}\left(P_{s}^{-1}-X_{s}^{-1}\right) B_{1}\right)$ is such. (Actually, a more detailed analysis allows to conclude that $\left(R_{1}-B_{1}^{\top}\left(P_{s}^{-1}-X_{s}^{-1}\right) B_{1}\right)>0$.) Then we have

$$
\begin{equation*}
\left[X_{s}+B_{1}\left(V-\frac{M^{\top} M}{\gamma^{2}}\right)^{-1} B_{1}^{\top}\right]^{-1}=\left[F_{1}\left(X_{s}-X_{s} L_{1}^{\top}\left(L_{1} X_{s} L_{1}^{\top}+R\right)^{-1} L_{1} X_{s}\right) F_{1}^{\top}\right]^{-1} \tag{2.65}
\end{equation*}
$$

which yields

$$
\begin{equation*}
X_{s}^{-1}+X_{s}^{-1} B_{1}\left(\frac{M^{\top} M}{\gamma^{2}}-V-B_{1}^{\top} X_{s}^{-1} B_{1}\right)^{-1} B_{1}^{\top} X_{s}^{-1}=F_{1}^{-\top}\left(X_{s}^{-1}+L_{1}^{\top} R^{-1} L_{1}\right) F_{1}^{-1} \tag{2.66}
\end{equation*}
$$

and hence
$X_{s}^{-1}=F_{1}^{\top} X_{s}^{-1} F_{1}-F_{1}^{\top} X_{s}^{-1} B_{1}\left(R_{1}-B_{1}^{\top}\left(P_{s}^{-1}-X_{s}^{-1}\right) B_{1}\right)^{-1} B_{1}^{\top} X_{s}^{-1} F_{1}-L_{1}^{\top} R^{-1} L_{1}$

$$
\begin{equation*}
=F_{1}^{\top} X_{s}^{-1}\left(I-B_{1} R_{1}^{-1} B_{1}^{\top}\left(P_{s}^{-1}-X_{s}^{-1}\right)\right)^{-1}\left(I-B_{1} R_{1}^{-1} B_{1}^{\top} P_{s}^{-1}\right) F_{1}-L_{1}^{\top} R^{-1} L_{1} . \tag{2.67a}
\end{equation*}
$$

Finally, by plugging (2.55) in (2.67b) we obtain (2.54). $\quad$.
Proposition 2.1. Let $\gamma>\gamma_{0}$ and $\gamma \notin \mathcal{G}_{1}$. Then the ARE (2.52) admits the (necessarily unique) antistabilizing solution. In fact, such solution is given by ${ }^{3}$

$$
\begin{equation*}
Y_{a}:=P_{s}^{-1}-X_{s}^{-1} \tag{2.68}
\end{equation*}
$$

Proof. We have to show that when $Y=Y_{a},(2.52)$ is an identity and that all the eigenvalues of

$$
\begin{equation*}
\Gamma_{y}:=F-B_{1}\left(B_{1}^{\top} Y_{a} B_{1}-R_{1}\right)^{-1}\left(B_{1}^{\top} Y_{a} F+\frac{M^{\top} L}{\gamma^{2}}\right) \tag{2.69}
\end{equation*}
$$

lay in $\{z \in \mathbb{C}:|z|>1\}$.
We first prove that $Y_{a}$ is a solution of the ARE

$$
\begin{equation*}
Y=F_{y}^{\top} Y F_{y}+F_{y}^{\top} Y B_{1}\left(R_{1}-B_{1}^{\top} Y B_{1}\right)^{-1} B_{1}^{\top} Y F_{y}-H_{0}^{\top} J_{1}^{-1} H_{0} \tag{2.70}
\end{equation*}
$$

From (2.23) we get

$$
\begin{equation*}
\left(P_{s}-B_{1} B_{1}^{\top}\right)^{-1}=F^{-\top}\left(P_{s}-P_{s} C^{\top}\left(C P_{s} C^{\top}+D D^{\top}\right)^{-1} C P_{s}\right)^{-1} F^{-1} \tag{2.71}
\end{equation*}
$$

which, using the same procedure that lead to (2.67a), yields

$$
\begin{equation*}
P_{s}^{-1}=F^{\top} P_{s}^{-1} F-F^{\top} P_{s}^{-1} B_{1}\left(B_{1}^{\top} P_{s}^{-1} B_{1}-I\right)^{-1} B_{1}^{\top} P_{s}^{-1} F-C^{\top}\left(D D^{\top}\right)^{-1} C \tag{2.72}
\end{equation*}
$$

which, taking (2.42) into account, may be also written in the form

$$
\begin{equation*}
P_{s}^{-1}=F_{a}^{\top} P_{s}^{-1} F-C^{\top}\left(D D^{\top}\right)^{-1} C \tag{2.73}
\end{equation*}
$$

and, finally, in the form

$$
\begin{equation*}
P_{s}^{-1}=F_{1}^{\top} P_{s}^{-1} F_{y}+F_{a}^{\top} P_{s}^{-1}\left(F-F_{y}\right)-\left(F_{1}-F_{a}\right)^{\top} P_{s}^{-1} F_{y}-C^{\top}\left(D D^{\top}\right)^{-1} C . \tag{2.74}
\end{equation*}
$$

By subtracting (2.54) from (2.74) we now get

$$
\begin{equation*}
Y_{a}=F_{1}^{\top}\left[P_{s}^{-1}-X_{s}^{-1}\left(I-B_{1} R_{1}^{-1} B_{1}^{\top} Y_{a}\right)^{-1}\right] F_{y}+R_{2} \tag{2.75}
\end{equation*}
$$

with
(2.76) $R_{2}:=F_{a}^{\top} P_{s}^{-1}\left(F-F_{y}\right)-\left(F_{1}-F_{a}\right)^{\top} P_{s}^{-1} F_{y}-C^{\top}\left(D D^{\top}\right)^{-1} C+L_{1}^{\top} R^{-1} L_{1}$.

We may rewrite (2.75) as follows:

$$
\begin{align*}
Y_{a} & =F_{1}^{\top}\left[\left(Y_{a}-P_{s}^{-1} B_{1} R_{1}^{-1} B_{1}^{\top} Y_{a}\right)\left(I-B_{1} R_{1}^{-1} B_{1}^{\top} Y_{a}\right)^{-1}\right] F_{y}+R_{2} \\
& =F_{1}^{\top}\left(I-P_{s}^{-1} B_{1} R_{1}^{-1} B_{1}^{\top}\right) Y_{a}\left(I-B_{1} R_{1}^{-1} B_{1}^{\top} Y_{a}\right)^{-1} F_{y}+R_{2} \\
& =F_{y}^{\top} Y_{a}\left(I-B_{1} R_{1}^{-1} B_{1}^{\top} Y_{a}\right)^{-1} F_{y}+R_{2} \\
& =F_{y}^{\top} Y_{a} F_{y}+F_{y}^{\top} Y_{a} B_{1}\left(R_{1}-B_{1}^{\top} Y_{a} B_{1}\right)^{-1} B_{1}^{\top} Y_{a} F_{y}+R_{2}, \tag{2.77}
\end{align*}
$$

where we have taken (2.55) into account.

[^3]To prove that $Y_{a}$ is a solution of the $\operatorname{ARE}$ (2.70), it remains only to show that

$$
\begin{equation*}
R_{2}=-H_{0}^{\top} J_{1}^{-1} H_{0} \tag{2.78}
\end{equation*}
$$

Indeed, taking into account (2.51), (2.53), and the first of (2.41), we have

$$
\begin{align*}
& R_{2}+H_{0}^{\top} J_{1}^{-1} H_{0} \\
= & F_{a}^{\top} P_{s}^{-1}\left(F-F_{y}\right)-\left(F_{1}-F_{a}\right)^{\top} P_{s}^{-1} F_{y}+L_{1}^{\top} R^{-1} L_{1}+L^{\top}\left(M M^{\top}-\gamma^{2} I\right)^{-1} L \\
= & F_{a}^{\top} P_{s}^{-1} B_{1} M^{\top}\left(M M^{\top}-\gamma^{2} I\right)^{-1} L-L_{1}^{\top} R^{-1} M V^{-1} B_{1}^{\top} P_{s}^{-1} F \\
& +L_{1}^{\top} R^{-1} M \underbrace{V^{-1} B_{1}^{\top} P_{s}^{-1} B_{1}}_{V^{-1}-I} M^{\top}\left(M M^{\top}-\gamma^{2} I\right)^{-1} L \\
& +L_{1}^{\top} R^{-1} L_{1}+L^{\top}\left(M M^{\top}-\gamma^{2} I\right)^{-1} L \\
= & F_{a}^{\top} P_{s}^{-1} B_{1} M^{\top}\left(M M^{\top}-\gamma^{2} I\right)^{-1} L-L_{1}^{\top} R^{-1} M V^{-1} B_{1}^{\top} P_{s}^{-1} F \\
& +L_{1}^{\top} R^{-1} \underbrace{M V^{-1} M^{\top}}_{-R+\gamma^{2} I}\left(M M^{\top}-\gamma^{2} I\right)^{-1} L \\
& -L_{1}^{\top} R^{-1} \underbrace{M}_{\left(M M^{\top}\right.}\left(M M^{\top}-\gamma^{2} I\right)^{-1} L \\
& +L_{1}^{\top} R^{-1} L_{1}+L^{\top}\left(M M^{\top}-\gamma^{2} I\right)^{-1} L \\
= & F_{a}^{\top} P_{s}^{-1} B_{1} M^{\top}\left(M M^{\top}-\gamma^{2} I\right)^{-1} L-L_{1}^{\top} R^{-1} M V^{-1} B_{1}^{\top} P_{s}^{-1} F \\
& +\underbrace{\left(M M^{\top}-\gamma^{2} I\right)^{-1} L+L_{1}^{\top} R^{-1} \underbrace{\left(L_{1}-L\right.}_{P_{1}^{\top} P_{s}^{-1} F_{a}})}_{\left.M^{\top}-L_{1}^{\top}\right)} \\
= & L_{1}^{\top} F_{a}^{\top} P_{s}^{-1} M\left(B_{1} M^{\top}\right. \\
= & L_{1}^{\top} R^{-1} M(B_{1}^{\top} P_{s}^{-1} F_{a}-V_{s}^{-1} F_{a}-B_{1}^{\top} P_{1}^{\top} P_{s}^{P_{s}^{-1}} \underbrace{\left.\left(I-B_{1} B_{1}^{\top} P_{s}^{-1}\right)^{-1} F\right)}_{F_{a}})=0 .  \tag{2.79}\\
79) &
\end{align*}
$$

We now prove that $Y_{a}$ is a solution of (2.52) as well. In fact, we have

$$
\begin{aligned}
F_{y} & =F-B_{1} M^{\top}\left(M M^{\top}-\gamma^{2} I\right)^{-1} L=F+B_{1} M^{\top}\left(\frac{I}{\gamma^{2}}+\frac{M R_{1}^{-1} M^{\top}}{\gamma^{4}}\right) L \\
(2.80) & =F+B_{1} R_{1}^{-1} \frac{M^{\top} L}{\gamma^{2}}
\end{aligned}
$$

and

$$
\begin{align*}
H_{0}^{\top} J_{1}^{-1} H_{0} & =C^{\top}\left(D D^{\top}\right)^{-1} C+L^{\top}\left(M M^{\top}-\gamma^{2} I\right)^{-1} L \\
& =C^{\top}\left(D D^{\top}\right)^{-1} C-\frac{L^{\top} L}{\gamma^{2}}-\frac{L^{\top} M}{\gamma^{2}} R_{1}^{-1} \frac{M^{\top} L}{\gamma^{2}} \tag{2.81}
\end{align*}
$$

so that, by using standard manipulations [15, p. 271], (2.70) may be rewritten in the form (2.52).

It remains to prove that the solution $Y_{a}$ is indeed antistabilizing, i.e., that all the eigenvalues of $\Gamma_{y}$ lay in $\{z \in \mathbb{C}:|z|>1\}$. Notice that $\Gamma_{y}$ may be written in the form

$$
\begin{equation*}
\Gamma_{y}=F_{y}+B_{1}\left(R_{1}-B_{1}^{\top} Y_{a} B_{1}\right)^{-1} B_{1}^{\top} Y_{a} F_{y}=\left(I-B_{1} R_{1}^{-1} B_{1}^{\top} Y_{a}\right)^{-1} F_{y} \tag{2.82}
\end{equation*}
$$

Moreover, we rewrite (2.54) in the form

$$
\begin{equation*}
X_{s}^{-1}+L_{1}^{\top} R^{-1} L_{1}=F_{1}^{\top} X_{s}^{-1} \Gamma_{y} \tag{2.83}
\end{equation*}
$$

which, together with (2.63) yields

$$
\begin{equation*}
I=X_{s}\left(I+L_{1}^{\top} R^{-1} L_{1} X_{s}\right)^{-1} F_{1}^{\top} X_{s}^{-1} \Gamma_{y}=X_{s} \Gamma_{a}^{\top} X_{s}^{-1} \Gamma_{y} \tag{2.84}
\end{equation*}
$$

so that $\Gamma_{y}$ is clearly antistable.
3. A monotonicity result. In this section we prove that the values of $\gamma$ for which $Y_{a}(\gamma)$ is singular are finitely many. This will be usefull in order to prove existence of the stabilizing solution of (1.14).

Proposition 3.1. Let $Y_{a}(\gamma)$ be as in (2.68). The set

$$
\begin{equation*}
\mathcal{G}_{2}:=\left\{\gamma>\gamma_{0}: Y_{a}(\gamma) \text { is singular }\right\} \tag{3.1}
\end{equation*}
$$

contains, at most, $n$ points.
Proof. The solution $X_{s}$ of (2.43) is a continuous function of $\gamma$ and its first derivative with respect to $\gamma^{2}$ exists and is continuous; see, e.g., [15, Theorem 14.2.2]. Since $P_{s}$ does not depend on $\gamma$, the derivative $\frac{d Y_{a}}{d \gamma^{2}}$ exists and is continuos as well. Consider an open set $\mathcal{I}=(a, b)$ such that $a \geq \gamma_{0}, b>a$ (possibly $b=\infty$ ) and $\mathcal{I} \cap \mathcal{G}_{1}=\emptyset$. (Notice that the set $\left\{\gamma \in \mathbb{R}: \gamma>\gamma_{0}\right\}$ may be written as the union of a finite number of sets of this form and of a finite number of isolated points.) For $\gamma \in \mathcal{I}$, we may take derivatives in both sides of (2.52) and, by defining

$$
\begin{equation*}
Y_{\gamma}:=\frac{d Y_{a}}{d \gamma^{2}} \tag{3.2}
\end{equation*}
$$

we get

$$
\begin{aligned}
Y_{\gamma}= & F^{\top} Y_{\gamma} F-\left(F^{\top} Y_{\gamma} B_{1}-\frac{L^{\top} M}{\gamma^{4}}\right)\left(R_{1}-B_{1}^{\top} Y_{a} B_{1}\right)^{-1}\left(B_{1}^{\top} Y_{a} F+\frac{M^{\top} L}{\gamma^{2}}\right) \\
- & \left(F^{\top} Y_{a} B_{1}+\frac{L^{\top} M}{\gamma^{2}}\right)\left(R_{1}-B_{1}^{\top} Y_{a} B_{1}\right)^{-1}\left(B_{1}^{\top} Y_{\gamma} F-\frac{M^{\top} L}{\gamma^{4}}\right) \\
- & \left(F^{\top} Y_{a} B_{1}+\frac{L^{\top} M}{\gamma^{2}}\right)\left(R_{1}-B_{1}^{\top} Y_{a} B_{1}\right)^{-1} \frac{M^{\top} M}{\gamma^{4}} \\
& \left(R_{1}-B_{1}^{\top} Y_{a} B_{1}\right)^{-1}\left(B_{1}^{\top} Y_{a} F+\frac{M^{\top} L}{\gamma^{2}}\right) \\
+ & \left(F^{\top} Y_{a} B_{1}+\frac{L^{\top} M}{\gamma^{2}}\right)\left(R_{1}-B_{1}^{\top} Y_{a} B_{1}\right)^{-1} B_{1}^{\top} Y_{\gamma} B_{1} \\
& \left(R_{1}-B_{1}^{\top} Y_{a} B_{1}\right)^{-1}\left(B_{1}^{\top} Y_{a} F+\frac{M^{\top} L}{\gamma^{2}}\right) \\
- & \frac{L^{\top} L}{\gamma^{4}}
\end{aligned}
$$

and, taking (2.69) into account,

$$
\begin{equation*}
Y_{\gamma}=\Gamma_{y}^{\top} Y_{\gamma} \Gamma_{y}-\frac{L_{2}^{\top} L_{2}}{\gamma^{4}} \tag{3.4}
\end{equation*}
$$

with

$$
\begin{equation*}
L_{2}:=L-M\left(R_{1}-B_{1}^{\top} Y_{a} B_{1}\right)^{-1}\left(B_{1}^{\top} Y_{a} F+\frac{M^{\top} L}{\gamma^{2}}\right) \tag{3.5}
\end{equation*}
$$

Taking (2.82) into account, we may rewrite (2.70) (with $Y+Y_{a}$ ) in the form

$$
\begin{equation*}
Y_{a}=\Gamma_{y}^{\top} Y_{a} \Gamma_{y}-\Gamma_{y}^{\top} Y_{a} B_{1} R_{1}^{-1} B_{1}^{\top} Y_{a} \Gamma_{y}-H_{0}^{\top} J_{1}^{-1} H_{0} \tag{3.6}
\end{equation*}
$$

Define

$$
\begin{equation*}
Z:=Y_{a}+\gamma^{2} Y_{\gamma} \tag{3.7}
\end{equation*}
$$

By adding (3.4) multiplied by $\gamma^{2}$ to (3.6), we get

$$
\begin{equation*}
Z=\Gamma_{y}^{\top} Z \Gamma_{y}-\Gamma_{y}^{\top} Y_{a} B_{1} R_{1}^{-1} B_{1}^{\top} Y_{a} \Gamma_{y}-H_{0}^{\top} J_{1}^{-1} H_{0}-\frac{L_{2}^{\top} L_{2}}{\gamma^{2}} \tag{3.8}
\end{equation*}
$$

By taking into account the definitions (2.69) and (3.5) of $\Gamma_{y}$ and $L_{2}$, respectively, and identity (2.81), it is not difficult to see that
(3.9) $-\Gamma_{y}^{\top} Y_{a} B_{1} R_{1}^{-1} B_{1}^{\top} Y_{a} \Gamma_{y}-H_{0}^{\top} J_{1}^{-1} H_{0}-\frac{L_{2}^{\top} L_{2}}{\gamma^{2}}=-C_{1}^{\top} C_{1}-C^{\top}\left(D D^{\top}\right)^{-1} C$
with

$$
\begin{equation*}
C_{1}:=\left(R_{1}-B_{1}^{\top} Y_{a} B_{1}\right)^{-1}\left(B_{1}^{\top} Y_{a} F+\frac{M^{\top} L}{\gamma^{2}}\right) \tag{3.10}
\end{equation*}
$$

Therefore, $Z$ satisfies

$$
\begin{equation*}
Z=\Gamma_{y}^{\top} Z \Gamma_{y}-C_{1}^{\top} C_{1}-C^{\top}\left(D D^{\top}\right)^{-1} C \tag{3.11}
\end{equation*}
$$

Notice that since $\Gamma_{y}$ is antistable, (3.4) and (3.11) imply that

$$
\begin{align*}
& Y_{\gamma} \geq 0  \tag{3.12a}\\
& Z \geq 0 \tag{3.12b}
\end{align*}
$$

Moreover, by multiplying (3.4) and (3.11) on the left side by $\Gamma_{y}^{-\top}$ and on the right side by $\Gamma_{y}^{-1}$, it is easy to see that both $\operatorname{ker} Y_{\gamma}$ and $\operatorname{ker} Z$ are invariant for $\Gamma_{y}^{-1}$ and hence for $\Gamma_{y}$. Then

$$
\begin{equation*}
\operatorname{ker} Y_{\gamma} \subseteq \operatorname{ker} L_{2} \tag{3.13}
\end{equation*}
$$

and

$$
\operatorname{ker} Z \subseteq \operatorname{ker}\left[\begin{array}{c}
C_{1}  \tag{3.14}\\
C
\end{array}\right]
$$

Therefore,

$$
\mathcal{K}:=\operatorname{ker}\left[\begin{array}{l}
Y_{a}  \tag{3.15}\\
Y_{\gamma}
\end{array}\right]=\operatorname{ker}\left[\begin{array}{c}
Z \\
Y_{\gamma}
\end{array}\right]
$$

is $\Gamma_{y}$-invariant and satisfies

$$
\mathcal{K} \subseteq \operatorname{ker}\left[\begin{array}{c}
C  \tag{3.16}\\
C_{1} \\
L_{2}
\end{array}\right]
$$

We now prove that $\mathcal{K}=\{0\}$. To this end it is sufficient to show that

$$
\operatorname{ker}\left[\begin{array}{c}
\Gamma_{y}-\lambda I  \tag{3.17}\\
C \\
C_{1} \\
L_{2}
\end{array}\right]=\{0\} \quad \forall \lambda \in \mathbb{C} \text {. }
$$

Assume by contradiction that there exist $v \neq 0$ and $\lambda \in \mathbb{C}$ such that

$$
\left[\begin{array}{c}
\Gamma_{y}-\lambda I  \tag{3.18}\\
C \\
C_{1} \\
L_{2}
\end{array}\right] v=0
$$

Then, since $\Gamma_{y}$ is antistable, $|\lambda|>1$. Moreover, the definitions (1.6) of $F$ and (2.69) of $\Gamma_{y}$ yield

$$
\begin{equation*}
\Gamma_{y}=A-B D^{\top}\left(D D^{\top}\right)^{-1} C-B_{1} C_{1} \tag{3.19}
\end{equation*}
$$

so that $\Gamma_{y} v=A v$ and we get

$$
\left[\begin{array}{c}
A-\lambda I  \tag{3.20}\\
C
\end{array}\right] v=0, \quad|\lambda|>1,
$$

which is in contradiction with detectability of the pair $(A, C)$. Therefore $\mathcal{K}=0$ and then, when one eigenvalue of $Y_{a}$ (that is a continuous function of $Y_{a}$ and hence of $\gamma)$ is zero, its derivative with respect to $\gamma$ is positive. Thus, if $Y_{a}(\bar{\gamma})$ is singular and (counting with multiplicity) has, say, $k_{+}$positive, $k_{0}$ zero, and $k_{-}$negative eigenvalues, then there exists a positive value $\delta$ such that $Y_{a}(\gamma)$ has $k_{+}$positive and $k_{0}+k_{-}$negative eigenvalues for $\gamma \in(\bar{\gamma}-\delta, \bar{\gamma})$ and $k_{+}+k_{0}$ positive and $k_{-}$negative eigenvalues for $\gamma \in(\bar{\gamma}, \bar{\gamma}+\delta)$. Clearly, this may happen for, at most, $n$ different values of $\gamma$.

So far we have assumed that $\gamma \in \mathcal{I}$. As already observed, the set $\left\{\gamma \in \mathbb{R}: \gamma>\gamma_{0}\right\}$ may be written as the union of a finite number of sets of the same type of $\mathcal{I}$ and of a finite number of values of $\gamma$. In correspondence of such values, $Y_{a}$ remains a continuous function of $\gamma$ so that we can extend the conclusion to the whole set $\left\{\gamma \in \mathbb{R}: \gamma>\gamma_{0}\right\}$.
4. Existence of a stabilizing solution of the ARE (1.14). Next we show that when $Y_{a}$ is nonsingular, then $Y_{a}^{-1}$ is the (necessarily unique) stabilizing solution of (1.14), so that, in view of Propositions 2.1 and 3.1 , we get that, except for a finite number of values of $\gamma>\gamma_{0}$, the ARE (1.14) admits the stabilizing solution. We shall prove in the next section that such solution is also feasible.

Consider the set $\mathcal{G}_{1} \cup \mathcal{G}_{2}$ and observe that it contains, at most, $2[n+\operatorname{rank}(M)]$ points. Then, the set of regular values of $\gamma$ defined as

$$
\begin{equation*}
\mathcal{G}_{r}:=\left\{\gamma>\gamma_{0}: \gamma \notin \mathcal{G}_{1} \cup \mathcal{G}_{2}\right\} \tag{4.1}
\end{equation*}
$$

is generic in $\left\{\gamma \in \mathbb{R}: \gamma>\gamma_{0}\right\}$.
The following theorem is our main result.
Theorem 4.1. Let $\gamma \in \mathcal{G}_{r}$ and let $Y_{a}(\gamma)$ be the corresponding antistabilizing solution of (2.52). Then (1.14) admits a unique symmetric stabilizing solution. Such solution is given by

$$
\begin{equation*}
\Delta_{s}=Y_{a}^{-1} \tag{4.2}
\end{equation*}
$$

Proof. As shown in Proposition 2.1, $Y_{a}$ is a solution of (2.70) so that we have

$$
\begin{align*}
Y_{a}+H_{0}^{\top} J_{1}^{-1} H_{0} & =F_{y}^{\top} Y_{a} F_{y}+F_{y}^{\top} Y_{a} B_{1}\left(R_{1}-B_{1}^{\top} Y_{a} B_{1}\right)^{-1} B_{1}^{\top} Y_{a} F_{y} \\
& =F_{y}^{\top}\left(Y_{a}+Y_{a} B_{1}\left(R_{1}-B_{1}^{\top} Y_{a} B_{1}\right)^{-1} B_{1}^{\top} Y_{a}\right) F_{y} \tag{4.3}
\end{align*}
$$

Since $\gamma \in \mathcal{G}_{r}, Y_{a}, R_{1}$ and $\left(R_{1}-B_{1}^{\top} Y_{a} B_{1}\right)$ are nonsingular so that $Y_{a}+Y_{a} B_{1}\left(R_{1}-\right.$ $\left.B_{1}^{\top} Y_{a} B_{1}\right)^{-1} B_{1}^{\top} Y_{a}$ is such. Moreover, $F_{y}$ is nonsingular so that the left-hand side of (4.3) is nonsingular and the same procedure that led to (2.67a) gives
(4.4) $Y_{a}^{-1}=F_{y} Y_{a}^{-1} F_{y}^{\top}-F_{y} Y_{a}^{-1} H_{0}^{\top}\left(J_{1}+H_{0} Y_{a}^{-1} H_{0}^{\top}\right)^{-1} H_{0} Y_{a}^{-1} F_{y}^{\top}+B_{1} R_{1}^{-1} B_{1}^{\top}$,
which, taking (4.2) into account, yields, after some standard manipulations [15, p. 271]

$$
\begin{equation*}
\Delta_{s}=A \Delta_{s} A^{\top}+B B^{\top}-\left(A \Delta_{s} H_{0}^{\top}+B J_{0}^{\top}\right)\left(J_{1}+H_{0} \Delta_{s} H_{0}^{\top}\right)^{-1}\left(H_{0} \Delta_{s} A^{\top}+J_{0} B^{\top}\right) \tag{4.5}
\end{equation*}
$$

It remains to show that
$\Gamma_{s}:=A-\left(A \Delta_{s} H_{0}^{\top}+B J_{0}^{\top}\right)\left(J_{1}+H_{0} \Delta_{s} H_{0}^{\top}\right)^{-1} H_{0}=F_{y}-F_{y} \Delta_{s} H_{0}^{\top}\left(J_{1}+H_{0} \Delta_{s} H_{0}^{\top}\right)^{-1} H_{0}$
is stable. To this aim we rewrite (4.4) in the form

$$
\begin{equation*}
Y_{a}^{-1}-B_{1} R_{1}^{-1} B_{1}^{\top}=\Gamma_{s} Y_{a}^{-1} F_{y}^{\top} \tag{4.7}
\end{equation*}
$$

which, taking (2.82) into account, yields

$$
\begin{equation*}
I=\Gamma_{s} Y_{a}^{-1} F_{y}^{\top}\left(Y_{a}^{-1}-B_{1} R_{1}^{-1} B_{1}^{\top}\right)^{-1}=\Gamma_{s} Y_{a}^{-1} \Gamma_{y}^{\top} Y_{a} \tag{4.8}
\end{equation*}
$$

so that, since $\Gamma_{y}$ is antistable, $\Gamma_{s}$ is stable.
Remark 4.1. It is worth noticing that the computation of the stabilizing solution of the Riccati equation (1.14) can be numerically obtained by resorting to standard routines available in most control packages without the need of computing the solution the auxiliary Riccati equations (2.1) and (2.43). Indeed, those Riccati equations were only instrumental to the purpose of proving the existence of the stabilizing solution of (1.14).
5. Construction of the minimum phase $J$-spectral factor and $\mathcal{H}_{\infty}$ estimator design. The following important theorem gives a constructive procedure to obtain, from the solution $\Delta_{s}$, a minimum-phase $J$-spectral factor of $\Psi(z)$ having a realization with the same state matrix $A$ and the same output matrix $H_{0}$ of the realization (1.9) of $H(z)$.

Theorem 5.1. Let

$$
\begin{equation*}
\gamma \in \mathcal{G}_{r} \tag{5.1}
\end{equation*}
$$

and let $\Delta_{s}$ be the corresponding stabilizing solution of equation (1.14). Then,

1. The solution $\Delta_{s}$ is feasible, i.e., $J_{1}+H_{0} \Delta_{s} H_{0}^{\top}$ has the same inertia of $J_{p, l}(\gamma)$ and hence there exists a nonsingular matrix $\Lambda$ such that

$$
\begin{equation*}
\Lambda J_{p, l}(\gamma) \Lambda^{\top}=J_{1}+H_{0} \Delta_{s} H_{0}^{\top} \tag{5.2}
\end{equation*}
$$

2. The transfer matrix

$$
\begin{equation*}
\Omega_{s}(z):=H_{0}(z I-A)^{-1}\left(A \Delta_{s} H_{0}^{\top}+B J_{0}^{\top}\right)\left(J_{1}+H_{0} \Delta_{s} H_{0}^{\top}\right)^{-1} \Lambda+\Lambda \tag{5.3}
\end{equation*}
$$

is a square $J$-spectral factor of the $J$-spectral density $\Psi(z)$.
3. All the zeros of $\Omega_{s}(z)$ lay in the open unit disk, i.e., $\Omega_{s}(z)^{-1}$ is stable.

Proof. The proof of this theorem may be obtained following the same lines of Theorem 4.1 in [4].

We recall from [3] that there is a simple procedure that, from $\Omega_{s}(z)$, furnishes a realization of an $\mathcal{H}_{\infty}$ smoothing filter with attenuation level $\gamma$.
6. The critical values of $\gamma$ : A peculiar feature of the $J$-spectral factorization. In this section, we analyze the ARE (1.14) and the associated $J$-spectral factorization problem in the case when $\gamma \in \mathcal{G}_{s}:=\mathcal{G}_{1} \cup \mathcal{G}_{2}$. To this aim we write the set $\mathcal{G}_{s}$ as the union of disjoint sets $\mathcal{G}_{s}=\mathcal{G}_{1}^{\prime} \cup \mathcal{G}_{2}$ with $\mathcal{G}_{1}^{\prime}:=\mathcal{G}_{1} \cap \overline{\mathcal{G}_{2}}$ (usually $\mathcal{G}_{1}^{\prime}=\mathcal{G}_{1}$ ) and we consider the cases $\gamma \in \mathcal{G}_{1}^{\prime}$ and $\gamma \in \mathcal{G}_{2}$ separately.

For $\gamma \in \mathcal{G}_{1}^{\prime}, Y_{a}(\gamma)$ is still nonsingular and we can define $\Delta_{s}:=Y_{a}(\gamma)^{-1}$. In all the manifold examples that we have worked out, it turns out that, even for $\gamma \in$ $\mathcal{G}_{1}^{\prime}, \Delta_{s}$ defined in this way continues to be the stabilizing solution of (1.14). This leads us to conjecture that Theorem 4.1 holds (possibly with the Moore-Penrose pseudoinverse in place of the inverse in the $\operatorname{ARE}$ (1.14)) for all $\gamma \in \mathcal{G}_{r} \cup \mathcal{G}_{1}^{\prime}$. In view of continutity, to prove this generalization it would be sufficient to show that for $\gamma \in \mathcal{G}_{1}^{\prime}$, ( $J_{1}+H_{0} \Delta_{s} H_{0}^{\top}$ ) is invertible or (using the pseudoinverse in place of the inverse) that $\left(A \Delta_{s} H_{0}^{\top}+B J_{0}^{\top}\right)\left(J_{1}+H_{0} \Delta_{s} H_{0}^{\top}\right)^{\sharp}\left(H_{0} \Delta_{s} A^{\top}+J_{0} B^{\top}\right)$ is a continuous function of $\gamma$ for each $\gamma \in \mathcal{G}_{1}^{\prime}$.

Much more interesting is the behavior associated with $\gamma \in \mathcal{G}_{2}$. In this case, as shown in the following example, the ARE does not admit a stabilizing solution nor does the $J$-spectral density admit a minimum phase $J$-spectral factor having a realization with state matrix equal to $A$. This is a peculiar (and in the authors' opinion, rather counterintuitive) feature of the $J$-spectral factorization. In fact, in the standard (positive) spectral factorization this phenomenon cannot occur.

Let us consider the following very simple example: $A=2, C=L=1, B=[1 \mid 0]$, $D=[0 \mid 1]$, and $M=[0 \mid 0]$. In this case $W(z)$ defined in (1.20) is given by

$$
\begin{equation*}
W(z)=\frac{1}{2\left(3-z-z^{-1}\right)} \tag{6.1}
\end{equation*}
$$

so that

$$
\begin{equation*}
\gamma_{0}^{2}=\sup _{\vartheta}\left|\frac{1}{2(3-2 \cos (\vartheta))}\right|=\frac{1}{2} \tag{6.2}
\end{equation*}
$$

Moreover, the stabilizing solution of the ARE (2.1) is easily computed to be $P_{s}=$ $2+\sqrt{5}$ and the ARE (2.43) assumes the form

$$
\begin{equation*}
X^{2}+\left(\frac{3+\sqrt{5}-\gamma^{2}(10+6 \sqrt{5})}{4}\right) X+\gamma^{2}\left(\frac{3+\sqrt{5}}{4}\right)=0 \tag{6.3}
\end{equation*}
$$

whose stabilizing solution $X_{s}(\gamma)$ equals $P_{s}$ for $\gamma=1$. Moreover, $X_{s}(\gamma)>P_{s}$ for $\gamma>1$ and $X_{s}(\gamma)<P_{s}$ for $\gamma<1$. Thus, in this case, $\mathcal{G}_{2}=\{1\}$. The ARE (1.14) assumes the form

$$
\begin{equation*}
3 \Delta+1-\frac{4 \Delta^{2}}{\Delta\left(1-\gamma^{2}\right)-\gamma^{2}}\left(1-\gamma^{2}\right)=0 \tag{6.4}
\end{equation*}
$$

For $\gamma<1$, (6.4) admits a negative stabilizing solution $\Delta_{s}$ that tends to $-\infty$ as $\gamma \rightarrow 1_{-}$. For $\gamma>1$, (6.4) admits a positive stabilizing solution $\Delta_{s}$ that tends to $+\infty$ as $\gamma \rightarrow 1_{+}$.

For $\gamma=1$, the order of such equation collapses and the stabilizing solution does not exist any more: the only solution is $\Delta_{a}=-1 / 3$ and the corresponding closed-loop matrix is $A-0=A=2$. Notice that $\Delta_{a}=-1 / 3$ is feasible. In fact

$$
J_{1}+H_{0} \Delta_{a} H_{0}^{\top}=\frac{1}{3}\left[\begin{array}{cc}
2 & -1  \tag{6.5}\\
-1 & -4
\end{array}\right]
$$

has a positive and a negative eigenvalue so that there exists a matrix $\Lambda$ such that ${ }^{4}$

$$
\begin{equation*}
J_{1}+H_{0} \Delta_{a} H_{0}^{\top}=\Lambda J_{1,1}(1) \Lambda^{\top} \tag{6.6}
\end{equation*}
$$

The same computation used in the proof of Theorem 5.1 shows that

$$
\begin{align*}
\Omega_{a}(z) & :=H_{0}(z I-A)^{-1}\left(A \Delta_{a} H_{0}^{\top}+B J_{0}^{\top}\right)\left(J_{1}+H_{0} \Delta_{a} H_{0}^{\top}\right)^{-1} \Lambda+\Lambda \\
& =\left[\begin{array}{l}
1 \\
1
\end{array}\right](z-2)^{-1}\left[\left.-\frac{2}{3} \right\rvert\, \frac{2}{3}\right] \Lambda+\Lambda \tag{6.7}
\end{align*}
$$

is a $J$-spectral factor of the $J$-spectrum $\Psi(z)$ that, in this case, is given by

$$
\Psi(z)=\left[\begin{array}{cc}
\frac{1}{(z-2)\left(z^{-1}-2\right)}+1 & \frac{1}{(z-2)\left(z^{-1}-2\right)}  \tag{6.8}\\
\frac{1}{(z-2)\left(z^{-1}-2\right)} & \frac{1}{(z-2)\left(z^{-1}-2\right)}-1
\end{array}\right] .
$$

The numerator matrix of $\Omega_{a}(z)$ is $2-\left[\left.-\frac{2}{3} \right\rvert\, \frac{2}{3}\right] \Lambda \Lambda^{-1}\left[\begin{array}{l}1 \\ 1\end{array}\right]=2$ : the unique zero of $\Omega_{a}(z)$ lies outside the closed unit disk. (Such an $\Omega_{a}(z)$ is said to be a maximum phase $J$-spectral factor.) The set of all $J$-spectral factors of $\Psi(z)$ having a realization with state matrix $A=2$ may be obtained as follows. We set

$$
\Omega(z)=\left[\begin{array}{l}
h_{1}  \tag{6.9}\\
h_{2}
\end{array}\right](z-2)^{-1}\left[g_{1} \mid g_{2}\right]+\left[\begin{array}{ll}
d_{11} & d_{12} \\
d_{21} & d_{22}
\end{array}\right]
$$

and impose

$$
\begin{equation*}
\Omega(z) J_{1,1}(1) \Omega(z)=\Psi(z) . \tag{6.10}
\end{equation*}
$$

The corresponding solutions, up to a noninteresting change of basis in the state space, may be parametrizied as follows:

$$
\Omega(z)=\left[\begin{array}{l}
1  \tag{6.11}\\
1
\end{array}\right](z-2)^{-1}[g \mid \pm g]+\left[\begin{array}{cc}
\frac{-2-3 g^{2}}{6 g} & \pm \frac{2-3 g^{2}}{6 g} \\
\frac{-1+3 g^{2}}{3 g} & \pm \frac{1+3 g^{2}}{3 g}
\end{array}\right], \quad g \in \mathbb{R} \backslash\{0\}
$$

which is readily seen to be equal to the right-hand side of (6.7) as $\Lambda$ varies among the solutions of (6.6). In conclusion, we have produced a maximum-phase $J$-spectral

[^4](where the $\pm$ signs are either both + or both - ).
factor and have shown that the corresponding minimum-phase one (i.e., a minimumphase $J$-spectral factor having a realization with the same state-space matrix) does not exist. This is a very peculiar behavior that has no counterpart in the regular spectral factorization.

Notice that to obtain (6.11) we have imposed the McMillan degree ${ }^{5}$ of $\Omega(z)$ to be equal to 1 or, equivalently, we have restricted our search to minimal $J$-spectral factors, namely, $J$-spectral factors having the last possible McMillan degree. (Such McMillan degree is clearly equal to one-half of the McMillan degree of the $J$-spectrum $\Psi(z)$.$) See [16] for a discussion on the minimality of spectral factors. It is interesting$ to observe that extending the search to nonminimal $J$-spectral factors (i.e., $J$-spectral factors having larger McMillan degree), a minimum-phase $J$-spectral factor having a unique pole in $z=2$ does exist. Indeed, it is not difficult to check that

$$
\Omega_{m}(z)=\left(z I_{2}-2 I_{2}\right)^{-1}\left[\begin{array}{cc}
\frac{-11}{2 \sqrt{2}} & \frac{-3}{2 \sqrt{2}}  \tag{6.12}\\
-\frac{1}{\sqrt{2}} & \frac{3}{\sqrt{2}}
\end{array}\right]+\left[\begin{array}{cc}
\frac{-5 \sqrt{2}}{\frac{3}{3}} & \frac{-\sqrt{2}}{3} \\
\frac{-1}{3 \sqrt{2}} & \frac{7}{3 \sqrt{2}}
\end{array}\right]
$$

is a minimal realization of a $J$-spectral factor of $\Psi(z)$. As $H(z), \Omega_{m}(z)$ has a unique pole in $z=2$ and the corresponding numerator matrix

$$
2 I_{2}-\left[\begin{array}{cc}
\frac{-11}{2 \sqrt{2}} & \frac{-3}{2 \sqrt{2}}  \tag{6.13}\\
-\frac{1}{\sqrt{2}} & \frac{3}{\sqrt{2}}
\end{array}\right]\left[\begin{array}{cc}
\frac{-5 \sqrt{2}}{\frac{3}{3 \sqrt{2}}} & \frac{-\sqrt{2}}{3} \\
\frac{7}{3 \sqrt{2}}
\end{array}\right]^{-1}=\left[\begin{array}{cc}
\frac{1}{3} & \frac{1}{6} \\
-\frac{1}{6} & \frac{2}{3}
\end{array}\right]
$$

has a unique eigenvalue (of multiplicity 2) in $z=1 / 2$. Again, this fact has no counterpart in classical spectral factorization.

Eventually, notice that these discrepancies between $J$-spectral factorization and classical spectral factorization are associated to very particular $J$-spectra. In fact they may occur only for a finite number of values of $\gamma$.
6.1. Example with $\gamma \in \mathcal{G}_{r}$. In the following we consider the previous example in the case when $\gamma_{0}<\gamma<\gamma_{f}$ and $\gamma \in \mathcal{G}_{r}$. We design a smoothing filter starting from the stabilizing solution $\Delta_{s}$ of (1.14). Let $\gamma^{2}=3 / 4$. The stabilizing solution $\Delta_{s}$ is then given by $-4-\sqrt{13}$ and the corresponding $J$-spectral factor is given by

$$
\Omega_{s}(z)=\left[\begin{array}{l}
1  \tag{6.14}\\
1
\end{array}\right](z-2)^{-1}[2.942 \mid 6.374]+\left[\begin{array}{cc}
-0.3124 & 2.99 \\
0.2785 & 3.353
\end{array}\right]
$$

From $\Omega_{s}(z)$, by following the procedure described in [3], we easily get the following transfer function of a 1-step lag smoother,

$$
\begin{equation*}
S_{1}(z)=z\left[0.2063(z-0.02824)^{-1} 2.958+0.1906\right] \tag{6.15}
\end{equation*}
$$

and it is easy to check that

$$
\begin{equation*}
\left\|S_{1}(z) H_{1}(z)-H_{2}(z)\right\|=0.8383<\gamma \simeq 0.8660 \tag{6.16}
\end{equation*}
$$

7. Conclusions. In this paper a general J-spectral factorization problem was considered and its relation with the existence of the stabilizing solution of the associated Riccati equation was investigated. The stabilizing solution depends on a positive

[^5]parameter which represents the prescribed attenuation level for the underlying estimation problem. We have shown that the stabilizing solution of the ARE still exists (except for a finite number of values of $\gamma$ ) as long as a fixed-lag acausal estimator (smoother) does. A few aspects of the $J$-spectral factorization problem and the properties of its solutions are discussed in correspondence to the (finite number of) values of $\gamma$ for which the stabilizing solution of the ARE does not exist.

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[^1]:    ${ }^{1}$ We denote by $M^{\top}$ the transpose of a matrix $M$.

[^2]:    ${ }^{2}$ Notice that in some papers and books the definition of minimum phase requires the stability of the spectral factor. Here, since neither $A$ nor $F$ is assumed to be "stable," we are not interested in stable $J$-spectral factors and we adopt the definition of phase minimality given above.

[^3]:    ${ }^{3}$ We shall denote such a solution by $Y_{a}(\gamma)$ when we want to stress its dependence upon $\gamma$.

[^4]:    ${ }^{4}$ The set of solutions of (6.6) may be parametrized as

    $$
    \Lambda=\left[\begin{array}{cc}
    \frac{-2-3 g^{2}}{6 g} & \pm \frac{2-3 g^{2}}{6 g} \\
    \frac{-1+3 g^{2}}{3 g} & \pm \frac{1+3 g^{2}}{3 g}
    \end{array}\right], g \in \mathbb{R} \backslash\{0\}
    $$

[^5]:    ${ }^{5}$ We recall that the McMillan degree of a rational proper matrix function $P(z)$ is the state-space dimension of a minimal realization of $P(z)$; see [14] for more details.

