Trefftz Co-chain Calculus

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Abstract

• We propose a way to discretize linear stationary or time-harmonic elliptic problems on unbounded domains using co-chain calculus. • Co-chain calculus is a framework that accommodates both

- 1. finite element exterior calculus and
- 2. discrete exterior calculus.

As it relies on volume meshes, it is applied to a bounded domain Ω . • We couple

1. any method that fits co-chain calculus in Ω and

Trefftz Co-chain Calculus

Interface conditions:

$$\begin{cases} \mathbf{t} \left(\star_{\alpha} \mathrm{d} u |_{\Omega} \right) = \mathbf{t} \left(\star_{\alpha} \mathrm{d} u |_{\Omega_{\mathsf{T}}} \right) & \text{on } \Gamma \\ \mathbf{t} u |_{\Omega} = \mathbf{t} u |_{\Omega_{\mathsf{T}}} \end{cases}$$

Seek $u \in \Lambda^{l-1}(\Omega), v \in \mathcal{T}(\Omega_{\mathsf{T}})$: $\begin{cases} \int_{\Omega} \left(\star_{\alpha} \mathrm{d}u \wedge \mathrm{d}\omega + \star_{\gamma} u \wedge \omega \right) + (-1)^{l-1} \int_{\Gamma} \mathbf{t} \left(\star_{\alpha} \mathrm{d}v \right) \wedge \mathbf{t} \, \omega = \int_{\Omega} \phi \wedge \omega \\ (-1)^{l-1} \int_{\Gamma} \mathbf{t} \left(\star_{\alpha} \mathrm{d}w \right) \wedge \mathbf{t} \, u - (-1)^{l-1} \int_{\Gamma} \mathbf{t} \left(\star_{\alpha} \mathrm{d}w \right) \wedge \mathbf{t} \, v = 0 \end{cases}$

2. a Trefftz method in the unbounded $\Omega_T := \mathbb{R}^d \setminus \Omega, d \in \mathbb{N}^*$.



- Trefftz basis functions in $\mathcal{T}(\Omega_{T})$ solve the homogeneous equations exactly.
- Compared to BEM, Trefftz methods enjoy the advantages of
 - 1. a simpler assembly, as there are no singular integrals, and
 - 2. exponential convergence.

This entails a small number of DoFs that allows to compute the Schur complement of the final system.

 $\forall \omega \in \Lambda^{l-1}(\Omega), \forall w \in \mathcal{T}(\Omega_{\mathsf{T}})$ $\begin{cases} \left[\left(\mathbf{D}^{l-1} \right)^{\mathsf{H}} \mathbf{M}_{\alpha}^{l} \mathbf{D}^{l-1} + \mathbf{M}_{\gamma}^{l-1} \right] \vec{\mathbf{u}} + (-1)^{l-1} \left(\mathbf{T}_{\Gamma}^{l-1} \right)^{\mathsf{H}} \widetilde{\mathbf{K}}_{m,\Gamma}^{l-1} \mathbf{P}_{\Gamma} \vec{\mathbf{v}} = \widetilde{\mathbf{K}}_{m+1}^{l-1} \vec{\phi} \\ (-1)^{l-1} \mathbf{P}_{\Gamma}^{\mathsf{H}} \left(\widetilde{\mathbf{K}}_{m,\Gamma}^{l-1} \right)^{\mathsf{H}} \mathbf{T}_{\Gamma}^{l-1} \vec{\mathbf{u}} - \mathbf{M}_{\mathsf{T}} \vec{\mathbf{v}} = \mathbf{0} \end{cases}$ $\mathbf{P}_{\Gamma} \in \mathbb{C}^{\widetilde{N}_{m}^{\mathsf{bnd}}, N_{\mathsf{T}}} \mathsf{Projection} \qquad \mathbf{P}_{\Gamma} \vec{\mathbf{v}} \coloneqq \widetilde{\mathbf{T}}_{\Gamma}^{m} \tilde{\mathbf{j}}$ $\vec{\mathbf{v}} \in \mathbb{C}^{N_{\mathsf{T}}}$ Degrees of freedom $N_{\mathsf{T}} \coloneqq \dim \mathcal{T}^n(\Omega_{\mathsf{T}})$ (discrete Trefftz space) $\mathbf{M}_{\mathsf{T}} \in \mathbb{C}^{N_{\mathsf{T}},N_{\mathsf{T}}}$ Boundary energy $(\mathbf{M}_{\mathsf{T}})_{i,i} \coloneqq (-1)^l \int_{\Gamma} \mathbf{t} (\star_{\alpha} \mathrm{d} v_i^n) \wedge \mathbf{t} v_i^n$ Schur complement: $\left[\left(\mathbf{D}^{l-1} \right)^{\mathsf{H}} \mathbf{M}_{\alpha}^{l} \mathbf{D}^{l-1} + \mathbf{M}_{\gamma}^{l-1} + \left(\mathbf{T}_{\Gamma}^{l-1} \right)^{\mathsf{H}} \widetilde{\mathbf{K}}_{m,\Gamma}^{l-1} \mathbf{P}_{\Gamma} \mathbf{M}_{\mathsf{T}}^{-1} \mathbf{P}_{\Gamma}^{\mathsf{H}} \left(\widetilde{\mathbf{K}}_{m,\Gamma}^{l-1} \right)^{\mathsf{H}} \mathbf{T}_{\Gamma}^{l-1} \right] \vec{\mathbf{u}} = \widetilde{\mathbf{K}}_{m+1}^{l-1} \vec{\boldsymbol{\phi}}$

Concrete Example: Eddy Current

- $\mathbf{A}, \mathbf{B}, \mathbf{H}, \mathbf{j}, \mathbf{j}_0 : \mathbb{R}^3 \to \mathbb{C}^3.$
 - 1. Equilibrium equations:

$$\left\{ egin{array}{ll}
abla imes \mathbf{A} = \mathbf{B} \
abla imes \mathbf{H} = \mathbf{j} + \mathbf{j} \end{array}
ight.$$

Co-chain Calculus

 $u \in \Lambda^{l-1}(\mathbb{R}^n), \, \boldsymbol{\sigma} \in \Lambda^l(\mathbb{R}^n), \, \mathbf{j} \in \Lambda^m(\mathbb{R}^n), \, \psi, \, \phi \in \Lambda^{m+1}(\mathbb{R}^n).$ 1 Equilibrium equations.

$$\begin{cases} \mathrm{d}u = (-1)^l \boldsymbol{\sigma} \\ \mathrm{d}\mathbf{j} = \psi - \phi \end{cases}$$

2. Constitutive equations:

$$\left\{ egin{array}{ll} \mathbf{j} = \star_lpha oldsymbol{\sigma} \ \psi = \star_\gamma u \end{array}
ight.$$

 $\star_{\alpha}, \star_{\gamma}$ are Hodge operators, i.e. linear maps of l-forms into m-forms. Primary elimination:

$$(-1)^{l-1} \mathrm{d} \left(\star_{\alpha} \mathrm{d} u \right) + \star_{\gamma} u = \phi$$

 $\forall \eta \in \Lambda^{l-1}(\Omega),$

$$\int_{\Omega} (\star_{\alpha} \mathrm{d}u \wedge \mathrm{d}\eta + \star_{\gamma} u \wedge \eta) + (-1)^{l-1} \int_{\Gamma} \mathbf{t} (\star_{\alpha} \mathrm{d}u) \wedge \mathbf{t} \eta = \int_{\Omega} \phi \wedge \eta$$

Choosing primary and secondary meshes \mathcal{M}, \mathcal{M} (can be unrelated),

2. Constitutive equations:

 $\begin{cases} \mathbf{H} = \nu \mathbf{B} \\ \mathbf{j} = -i\omega\sigma \mathbf{A} \end{cases}$

Using the Finite Element Method:

Seek $\mathbf{A} \in \mathbf{H}(\mathbf{curl}, \Omega), \quad \mathbf{v} \in \boldsymbol{\mathcal{T}}(\Omega_{\mathsf{T}})$: $\int_{\Omega} \left[\nu \left(\nabla \times \mathbf{A} \right) \cdot \left(\nabla \times \mathbf{A} \right) + \imath \omega \sigma \mathbf{A} \cdot \mathbf{A} \right] \mathrm{d}\mathbf{x} + \int_{\Gamma} \nu \, \mathbf{n} \times \left(\nabla \times \mathbf{v} \right) \cdot \mathbf{A} \, \mathrm{d}S = \int_{\Omega} \mathbf{j}_0 \cdot \mathbf{A} \, \mathrm{d}\mathbf{x} \right]$ $\int_{\Gamma} \nu \, \mathbf{n} \times (\nabla \times \mathbf{A}) \cdot \mathbf{w} \, \mathrm{d}S - \int_{\Gamma} \nu \, \mathbf{n} \times (\nabla \times \mathbf{v}) \cdot \mathbf{w} \, \mathrm{d}S = 0$ $\forall \Lambda \in \mathbf{H}(\mathbf{curl}, \Omega), \, \forall \mathbf{w} \in \mathcal{T}(\Omega_{\mathsf{T}})$ Using the Cell Method: $\begin{cases} \left(\mathbf{C}^\mathsf{T} \mathbf{M}_{\nu} \mathbf{C} + \imath \omega \mathbf{M}_{\sigma} \right) \vec{\mathbf{a}} + \widetilde{\mathbf{C}}_{\Gamma} \mathbf{P}_{\Gamma} \vec{\mathbf{v}} = \vec{\mathbf{j}}_0 \\ \mathbf{P}_{\Gamma}^\mathsf{T} \widetilde{\mathbf{C}}_{\Gamma}^\mathsf{T} \vec{\mathbf{a}} - \mathbf{M}_\mathsf{T} \vec{\mathbf{v}} = \mathbf{0} \end{cases}$

Numerical Example Photonic Nanojet: 1. $r_{\bullet} = 1$ 2. $\epsilon_{\bullet} = 2.5281 \epsilon_0$ 3. $\mu_{\bullet} = \mu_0$

we end up with the discrete system:

 $\left[\left(\mathbf{D}^{l-1} \right)^{\mathsf{H}} \mathbf{M}_{\alpha}^{l} \mathbf{D}^{l-1} + \mathbf{M}_{\gamma}^{l-1} \right] \vec{\mathbf{u}} + (-1)^{l-1} \left(\mathbf{T}_{\Gamma}^{l-1} \right)^{\mathsf{H}} \widetilde{\mathbf{K}}_{m,\Gamma}^{l-1} \widetilde{\mathbf{T}}_{\Gamma}^{m} \widetilde{\mathbf{j}} = \widetilde{\mathbf{K}}_{m+1}^{l-1} \vec{\phi}$

 $\mathbf{D}^{l-1} \in \{-1, 0, 1\}^{N_l, N_{l-1}}$ Exterior derivative $\mathbf{M}_{\alpha}^{l} \in \mathbb{C}^{N_{l},N_{l}}, \, \mathbf{M}_{\gamma}^{l-1} \in \mathbb{C}^{N_{l-1},N_{l-1}}$ Mass¹ $ec{\mathbf{u}} \in \mathbb{C}^{N_{l-1}}, \, \widetilde{\mathbf{j}} \in \mathbb{C}^{\widetilde{N}_m}, \, ec{oldsymbol{\phi}} \in \mathbb{C}^{\widetilde{N}_{m+1}}$ Degrees of freedom² $\mathbf{T}_{\Gamma}^{l-1} \in \{0,1\}^{N_{l-1}^{\mathsf{bnd}},N_{l-1}}, \ \widetilde{\mathbf{T}}_{\Gamma}^{m} \in \{0,1\}^{N_{m}^{\mathsf{bnd}},N_{m}}$ Trace $\widetilde{\mathbf{K}}_{m,\Gamma}^{l-1} \in \mathbb{C}^{N_{l-1}^{\mathrm{bnd}},\widetilde{N}_m^{\mathrm{bnd}}}, \ \widetilde{\mathbf{K}}_{m+1}^{l-1} \in \mathbb{C}^{N_{l-1}^{\mathrm{bnd}},\widetilde{N}_{m+1}^{\mathrm{bnd}}}$ Pairing³

1. Square, Hermitian, and positive-definite matrices. 2. Related to integrals of u, j, ϕ over entities of \mathcal{M} or \mathcal{M} . 3. Discrete representative of the \wedge -product.

4. $\omega = 23.56 \cdot 10^8 \text{ rad s}^{-1}$ 5. $\epsilon_0 = 8.85 \cdot 10^{-12} \,\mathrm{F \, m^{-1}}$ 6. $\mu_0 = 4\pi \cdot 10^{-7} \,\mathrm{H\,s^{-1}}$

- 1. Radius
- 2. Permittivity inside
- 3. Permeability inside
- 4. Angular frequency
- 5. Permittivity outside
- 6. Permeability outside



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