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T. Aaltonen *et al.* (CDF Collaboration)

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Measurements of Direct CP Violating Asymmetries in Charmless Decays of Strange Bottom Mesons and Bottom Baryons

T. Aaltonen,²¹ B. Álvarez González^{v,9}, S. Amerio,⁴¹ D. Amidei,³² A. Anastassov,³⁶ A. Annovi,¹⁷ J. Antos,¹²
 G. Apollinari,¹⁵ J.A. Appel,¹⁵ A. Apresyan,⁴⁶ T. Arisawa,⁵⁶ A. Artikov,¹³ J. Asaadi,⁵¹ W. Ashmanskas,¹⁵
 B. Auerbach,⁵⁹ A. Aurisano,⁵¹ F. Azfar,⁴⁰ W. Badgett,¹⁵ A. Barbaro-Galtieri,²⁶ V.E. Barnes,⁴⁶ B.A. Barnett,²³
 P. Barria^{cc,44} P. Bartos,¹² M. Bauce^{aa,41} G. Bauer,³⁰ F. Bedeschi,⁴⁴ D. Beecher,²⁸ S. Behari,²³ G. Bellettini^{bb,44}
 J. Bellinger,⁵⁸ D. Benjamin,¹⁴ A. Beretvas,¹⁵ A. Bhatti,⁴⁸ M. Binkley^{*,15} D. Bisello^{aa,41} I. Bizjak^{gg,28} K.R. Bland,⁵
 B. Blumenfeld,²³ A. Bocci,¹⁴ A. Bodek,⁴⁷ D. Bortoletto,⁴⁶ J. Boudreau,⁴⁵ A. Boveia,¹¹ B. Brau^{a,15} L. Brigliadori^{z,6}
 A. Brisuda,¹² C. Bromberg,³³ E. Brucken,²¹ M. Bucciantonio^{bb,44} J. Budagov,¹³ H.S. Budd,⁴⁷ S. Budd,²²
 K. Burkett,¹⁵ G. Busetto^{aa,41} P. Bussey,¹⁹ A. Buzatu,³¹ C. Calancha,²⁹ S. Camarda,⁴ M. Campanelli,³³
 M. Campbell,³² F. Canelli^{12,15} A. Canepa,⁴³ B. Carls,²² D. Carlsmith,⁵⁸ R. Carosi,⁴⁴ S. Carrillo^{k,16} S. Carron,¹⁵
 B. Casal,⁹ M. Casarsa,¹⁵ A. Castro^{z,6} P. Catastini,¹⁵ D. Cauz,⁵² V. Cavaliere^{cc,44} M. Cavalli-Sforza,⁴ A. Cerri^{f,26}
 L. Cerrito^{q,28} Y.C. Chen,¹ M. Chertok,⁷ G. Chiarelli,⁴⁴ G. Chlachidze,¹⁵ F. Chlebana,¹⁵ K. Cho,²⁵
 D. Chokheli,¹³ J.P. Chou,²⁰ W.H. Chung,⁵⁸ Y.S. Chung,⁴⁷ C.I. Ciobanu,⁴² M.A. Ciocci^{cc,44} A. Clark,¹⁸
 G. Compostella^{aa,41} M.E. Convery,¹⁵ J. Conway,⁷ M. Corbo,⁴² M. Cordelli,¹⁷ C.A. Cox,⁷ D.J. Cox,⁷ F. Crescioli^{bb,44}
 C. Cuenca Almenar,⁵⁹ J. Cuevas^{v,9} R. Culbertson,¹⁵ D. Dagenhart,¹⁵ N. d'Ascenzo^{t,42} M. Datta,¹⁵ P. de Barbaro,⁴⁷
 S. De Cecco,⁴⁹ G. De Lorenzo,⁴ M. Dell'Orso^{bb,44} C. Deluca,⁴ L. Demortier,⁴⁸ J. Deng^{c,14} M. Deninno,⁶
 F. Devoto,²¹ M. d'Errico^{aa,41} A. Di Canto^{bb,44} B. Di Ruzza,⁴⁴ J.R. Dittmann,⁵ M. D'Onofrio,²⁷ S. Donati^{bb,44}
 P. Dong,¹⁵ M. Dorigo,⁵² T. Dorigo,⁴¹ K. Ebina,⁵⁶ A. Elagin,⁵¹ A. Eppig,³² R. Erbacher,⁷ D. Errede,²² S. Errede,²²
 N. Ershaidat^{y,42} R. Eusebi,⁵¹ H.C. Fang,²⁶ S. Farrington,⁴⁰ M. Feindt,²⁴ J.P. Fernandez,²⁹ C. Ferrazza^{dd,44}
 R. Field,¹⁶ G. Flanagan^{r,46} R. Forrest,⁷ M.J. Frank,⁵ M. Franklin,²⁰ J.C. Freeman,¹⁵ Y. Funakoshi,⁵⁶
 I. Furic,¹⁶ M. Gallinaro,⁴⁸ J. Galyardt,¹⁰ J.E. Garcia,¹⁸ A.F. Garfinkel,⁴⁶ P. Garosi^{cc,44} H. Gerberich,²²
 E. Gerchtein,¹⁵ S. Giagu^{ee,49} V. Giakoumopoulou,³ P. Giannetti,⁴⁴ K. Gibson,⁴⁵ C.M. Ginsburg,¹⁵ N. Giokaris,³
 P. Giromini,¹⁷ M. Giunta,⁴⁴ G. Giurgiu,²³ V. Glagolev,¹³ D. Glenzinski,¹⁵ M. Gold,³⁵ D. Goldin,⁵¹
 N. Goldschmidt,¹⁶ A. Golossanov,¹⁵ G. Gomez,⁹ G. Gomez-Ceballos,³⁰ M. Goncharov,³⁰ O. González,²⁹
 I. Gorelov,³⁵ A.T. Goshaw,¹⁴ K. Goulianos,⁴⁸ A. Gresele,⁴¹ S. Grinstein,⁴ C. Grosso-Pilcher,¹¹ R.C. Group,⁵⁵
 J. Guimaraes da Costa,²⁰ Z. Gunay-Unalan,³³ C. Haber,²⁶ S.R. Hahn,¹⁵ E. Halkiadakis,⁵⁰ A. Hamaguchi,³⁹
 J.Y. Han,⁴⁷ F. Happacher,¹⁷ K. Hara,⁵³ D. Hare,⁵⁰ M. Hare,⁵⁴ R.F. Harr,⁵⁷ K. Hatakeyama,⁵ C. Hays,⁴⁰ M. Heck,²⁴
 J. Heinrich,⁴³ M. Herndon,⁵⁸ S. Hewamanage,⁵ D. Hidas,⁵⁰ A. Hocker,¹⁵ W. Hopkins^{g,15} D. Horn,²⁴ S. Hou,¹
 R.E. Hughes,³⁷ M. Hurwitz,¹¹ U. Husemann,⁵⁹ N. Hussain,³¹ M. Hussein,³³ J. Huston,³³ G. Introzzi,⁴⁴ M. Iori^{ee,49}
 A. Ivanov^{o,7} E. James,¹⁵ D. Jang,¹⁰ B. Jayatilaka,¹⁴ E.J. Jeon,²⁵ M.K. Jha,⁶ S. Jindariani,¹⁵ W. Johnson,⁷
 M. Jones,⁴⁶ K.K. Joo,²⁵ S.Y. Jun,¹⁰ T.R. Junk,¹⁵ T. Kamon,⁵¹ P.E. Karchin,⁵⁷ Y. Kato^{n,39} W. Ketchum,¹¹
 J. Keung,⁴³ V. Khotilovich,⁵¹ B. Kilminster,¹⁵ D.H. Kim,²⁵ H.S. Kim,²⁵ H.W. Kim,²⁵ J.E. Kim,²⁵ M.J. Kim,¹⁷
 S.B. Kim,²⁵ S.H. Kim,⁵³ Y.K. Kim,¹¹ N. Kimura,⁵⁶ M. Kirby,¹⁵ S. Klimentenko,¹⁶ K. Kondo,⁵⁶ D.J. Kong,²⁵
 J. Konigsberg,¹⁶ A.V. Kotwal,¹⁴ M. Kreps,²⁴ J. Kroll,⁴³ D. Krop,¹¹ N. Krumnack^{l,5} M. Kruse,¹⁴ V. Krutelyov^{d,51}
 T. Kuhr,²⁴ M. Kurata,⁵³ S. Kwang,¹¹ A.T. Laasanen,⁴⁶ S. Lami,⁴⁴ S. Lammel,¹⁵ M. Lancaster,²⁸ R.L. Lander,⁷
 K. Lannon^{u,37} A. Lath,⁵⁰ G. Latino^{cc,44} I. Lazzizzera,⁴¹ T. LeCompte,² E. Lee,⁵¹ H.S. Lee,¹¹ J.S. Lee,²⁵
 S.W. Lee^{w,51} S. Leo^{bb,44} S. Leone,⁴⁴ J.D. Lewis,¹⁵ C.-J. Lin,²⁶ J. Linacre,⁴⁰ M. Lindgren,¹⁵ E. Lipeles,⁴³ A. Lister,¹⁸
 D.O. Litvintsev,¹⁵ C. Liu,⁴⁵ Q. Liu,⁴⁶ T. Liu,¹⁵ S. Lockwitz,⁵⁹ N.S. Lockyer,⁴³ A. Loginov,⁵⁹ D. Lucchesi^{aa,41}
 J. Lueck,²⁴ P. Lujan,²⁶ P. Lukens,¹⁵ G. Lungu,⁴⁸ J. Lys,²⁶ R. Lysak,¹² R. Madrak,¹⁵ K. Maeshima,¹⁵ K. Makhoul,³⁰
 P. Maksimovic,²³ S. Malik,⁴⁸ G. Manca^{b,27} A. Manousakis-Katsikakis,³ F. Margaroli,⁴⁶ C. Marino,²⁴ M. Martínez,⁴
 R. Martínez-Ballarín,²⁹ P. Mastrandrea,⁴⁹ M. Mathis,²³ M.E. Mattson,⁵⁷ P. Mazzanti,⁶ K.S. McFarland,⁴⁷
 P. McIntyre,⁵¹ R. McNulty^{i,27} A. Mehta,²⁷ P. Mehtala,²¹ A. Menzione,⁴⁴ C. Mesropian,⁴⁸ T. Miao,¹⁵ D. Mietlicki,³²
 A. Mitra,¹ H. Miyake,⁵³ S. Moed,²⁰ N. Moggi,⁶ M.N. Mondragon^{k,15} C.S. Moon,²⁵ R. Moore,¹⁵ M.J. Morello,¹⁵
 J. Morlock,²⁴ P. Movilla Fernandez,¹⁵ A. Mukherjee,¹⁵ Th. Muller,²⁴ P. Murat,¹⁵ M. Mussini^{z,6} J. Nachtman^{m,15}
 Y. Nagai,⁵³ J. Naganoma,⁵⁶ I. Nakano,³⁸ A. Napier,⁵⁴ J. Nett,⁵¹ C. Neu,⁵⁵ M.S. Neubauer,²² J. Nielsen^{e,26}
 L. Nodulman,² O. Norriella,²² E. Nurse,²⁸ L. Oakes,⁴⁰ S.H. Oh,¹⁴ Y.D. Oh,²⁵ I. Oksuzian,⁵⁵ T. Okusawa,³⁹
 R. Orava,²¹ L. Ortolan,⁴ S. Pagan Griso^{aa,41} C. Pagliarone,⁵² E. Palencia^{f,9} V. Papadimitriou,¹⁵ A.A. Paramonov,²
 J. Patrick,¹⁵ G. Pauletta^{ff,52} M. Paulini,¹⁰ C. Paus,³⁰ D.E. Pellett,⁷ A. Penzo,⁵² T.J. Phillips,¹⁴ G. Piacentino,⁴⁴
 E. Pianori,⁴³ J. Pilot,³⁷ K. Pitts,²² C. Plager,⁸ L. Pondrom,⁵⁸ K. Potamianos,⁴⁶ O. Poukhov^{*,13} F. Prokoshin^{x,13}

A. Pronko,¹⁵ F. Ptohos^h,¹⁷ E. Pueschel,¹⁰ G. Punzi^{bb},⁴⁴ J. Pursley,⁵⁸ A. Rahaman,⁴⁵ V. Ramakrishnan,⁵⁸ N. Ranjan,⁴⁶ I. Redondo,²⁹ P. Renton,⁴⁰ M. Rescigno,⁴⁹ F. Rimondi^z,⁶ L. Ristori⁴⁵,¹⁵ A. Robson,¹⁹ T. Rodrigo,⁹ T. Rodriguez,⁴³ E. Rogers,²² S. Rolli,⁵⁴ R. Roser,¹⁵ M. Rossi,⁵² F. Rubbo,¹⁵ F. Ruffini^{cc},⁴⁴ A. Ruiz,⁹ J. Russ,¹⁰ V. Rusu,¹⁵ A. Safonov,⁵¹ W.K. Sakumoto,⁴⁷ Y. Sakurai,⁵⁶ L. Santi^{ff},⁵² L. Sartori,⁴⁴ K. Sato,⁵³ V. Saveliev^t,⁴² A. Savoy-Navarro,⁴² P. Schlabach,¹⁵ A. Schmidt,²⁴ E.E. Schmidt,¹⁵ M.P. Schmidt^{*},⁵⁹ M. Schmitt,³⁶ T. Schwarz,⁷ L. Scodellaro,⁹ A. Scribano^{cc},⁴⁴ F. Scuri,⁴⁴ A. Sedov,⁴⁶ S. Seidel,³⁵ Y. Seiya,³⁹ A. Semenov,¹³ F. Sforza^{bb},⁴⁴ A. Sfyrla,²² S.Z. Shalhout,⁷ T. Shears,²⁷ P.F. Shepard,⁴⁵ M. Shimojima^s,⁵³ S. Shiraishi,¹¹ M. Shochet,¹¹ I. Shreyber,³⁴ A. Simonenko,¹³ P. Sinervo,³¹ A. Sissakian^{*},¹³ K. Sliwa,⁵⁴ J.R. Smith,⁷ F.D. Snider,¹⁵ A. Soha,¹⁵ S. Somalwar,⁵⁰ V. Sorin,⁴ P. Squillacioti,¹⁵ M. Stancari,¹⁵ M. Stanitzki,⁵⁹ R. St. Denis,¹⁹ B. Stelzer,³¹ O. Stelzer-Chilton,³¹ D. Stentz,³⁶ J. Strologas,³⁵ G.L. Strycker,³² Y. Sudo,⁵³ A. Sukhanov,¹⁶ I. Suslov,¹³ K. Takemasa,⁵³ Y. Takeuchi,⁵³ J. Tang,¹¹ M. Tecchio,³² P.K. Teng,¹ J. Thom^g,¹⁵ J. Thome,¹⁰ G.A. Thompson,²² E. Thomson,⁴³ P. Ttito-Guzmán,²⁹ S. Tkaczyk,¹⁵ D. Toback,⁵¹ S. Tokar,¹² K. Tollefson,³³ T. Tomura,⁵³ D. Tonelli,¹⁵ S. Torre,¹⁷ D. Torretta,¹⁵ P. Totaro^{ff},⁵² M. Trovato^{dd},⁴⁴ Y. Tu,⁴³ F. Ukegawa,⁵³ S. Uozumi,²⁵ A. Varganov,³² F. Vázquez^k,¹⁶ G. Velez,¹⁵ C. Vellidis,³ M. Vidal,²⁹ I. Vila,⁹ R. Vilar,⁹ J. Vizán,⁹ M. Vogel,³⁵ G. Volpi^{bb},⁴⁴ P. Wagner,⁴³ R.L. Wagner,¹⁵ T. Wakisaka,³⁹ R. Wallny,⁸ S.M. Wang,¹ A. Warburton,³¹ D. Waters,²⁸ M. Weinberger,⁵¹ W.C. Wester III,¹⁵ B. Whitehouse,⁵⁴ D. Whiteson^c,⁴³ A.B. Wicklund,² E. Wicklund,¹⁵ S. Wilbur,¹¹ F. Wick,²⁴ H.H. Williams,⁴³ J.S. Wilson,³⁷ P. Wilson,¹⁵ B.L. Winer,³⁷ P. Wittich^g,¹⁵ S. Wolbers,¹⁵ H. Wolfe,³⁷ T. Wright,³² X. Wu,¹⁸ Z. Wu,⁵ K. Yamamoto,³⁹ J. Yamaoka,¹⁴ T. Yang,¹⁵ U.K. Yang^p,¹¹ Y.C. Yang,²⁵ W.-M. Yao,²⁶ G.P. Yeh,¹⁵ K. Yi^m,¹⁵ J. Yoh,¹⁵ K. Yorita,⁵⁶ T. Yoshida^j,³⁹ G.B. Yu,¹⁴ I. Yu,²⁵ S.S. Yu,¹⁵ J.C. Yun,¹⁵ A. Zanetti,⁵² Y. Zeng,¹⁴ and S. Zucchelli^{z6}

(CDF Collaboration[†])

¹*Institute of Physics, Academia Sinica, Taipei, Taiwan 11529, Republic of China*

²*Argonne National Laboratory, Argonne, Illinois 60439, USA*

³*University of Athens, 157 71 Athens, Greece*

⁴*Institut de Física d'Altes Energies, ICREA, Universitat Autònoma de Barcelona, E-08193, Bellaterra (Barcelona), Spain*

⁵*Baylor University, Waco, Texas 76798, USA*

⁶*Istituto Nazionale di Fisica Nucleare Bologna, ^zUniversity of Bologna, I-40127 Bologna, Italy*

⁷*University of California, Davis, Davis, California 95616, USA*

⁸*University of California, Los Angeles, Los Angeles, California 90024, USA*

⁹*Instituto de Física de Cantabria, CSIC-University of Cantabria, 39005 Santander, Spain*

¹⁰*Carnegie Mellon University, Pittsburgh, Pennsylvania 15213, USA*

¹¹*Enrico Fermi Institute, University of Chicago, Chicago, Illinois 60637, USA*

¹²*Comenius University, 842 48 Bratislava, Slovakia; Institute of Experimental Physics, 040 01 Kosice, Slovakia*

¹³*Joint Institute for Nuclear Research, RU-141980 Dubna, Russia*

¹⁴*Duke University, Durham, North Carolina 27708, USA*

¹⁵*Fermi National Accelerator Laboratory, Batavia, Illinois 60510, USA*

¹⁶*University of Florida, Gainesville, Florida 32611, USA*

¹⁷*Laboratori Nazionali di Frascati, Istituto Nazionale di Fisica Nucleare, I-00044 Frascati, Italy*

¹⁸*University of Geneva, CH-1211 Geneva 4, Switzerland*

¹⁹*Glasgow University, Glasgow G12 8QQ, United Kingdom*

²⁰*Harvard University, Cambridge, Massachusetts 02138, USA*

²¹*Division of High Energy Physics, Department of Physics, University of Helsinki and Helsinki Institute of Physics, FIN-00014, Helsinki, Finland*

²²*University of Illinois, Urbana, Illinois 61801, USA*

²³*The Johns Hopkins University, Baltimore, Maryland 21218, USA*

²⁴*Institut für Experimentelle Kernphysik, Karlsruhe Institute of Technology, D-76131 Karlsruhe, Germany*

²⁵*Center for High Energy Physics: Kyungpook National University,*

Daegu 702-701, Korea; Seoul National University, Seoul 151-742,

Korea; Sungkyunkwan University, Suwon 440-746,

Korea; Korea Institute of Science and Technology Information, Daejeon 305-806, Korea; Chonnam National University, Gwangju 500-757, Korea; Chonbuk National University, Jeonju 561-756, Korea

²⁶*Ernest Orlando Lawrence Berkeley National Laboratory, Berkeley, California 94720, USA*

²⁷*University of Liverpool, Liverpool L69 7ZE, United Kingdom*

²⁸*University College London, London WC1E 6BT, United Kingdom*

²⁹*Centro de Investigaciones Energeticas Medioambientales y Tecnológicas, E-28040 Madrid, Spain*

³⁰*Massachusetts Institute of Technology, Cambridge, Massachusetts 02139, USA*

- ³¹*Institute of Particle Physics: McGill University, Montréal, Québec, Canada H3A 2T8; Simon Fraser University, Burnaby, British Columbia, Canada V5A 1S6; University of Toronto, Toronto, Ontario, Canada M5S 1A7; and TRIUMF, Vancouver, British Columbia, Canada V6T 2A3*
- ³²*University of Michigan, Ann Arbor, Michigan 48109, USA*
- ³³*Michigan State University, East Lansing, Michigan 48824, USA*
- ³⁴*Institution for Theoretical and Experimental Physics, ITEP, Moscow 117259, Russia*
- ³⁵*University of New Mexico, Albuquerque, New Mexico 87131, USA*
- ³⁶*Northwestern University, Evanston, Illinois 60208, USA*
- ³⁷*The Ohio State University, Columbus, Ohio 43210, USA*
- ³⁸*Okayama University, Okayama 700-8530, Japan*
- ³⁹*Osaka City University, Osaka 588, Japan*
- ⁴⁰*University of Oxford, Oxford OX1 3RH, United Kingdom*
- ⁴¹*Istituto Nazionale di Fisica Nucleare, Sezione di Padova-Trento, ^{aa}University of Padova, I-35131 Padova, Italy*
- ⁴²*LPNHE, Université Pierre et Marie Curie/IN2P3-CNRS, UMR7585, Paris, F-75252 France*
- ⁴³*University of Pennsylvania, Philadelphia, Pennsylvania 19104, USA*
- ⁴⁴*Istituto Nazionale di Fisica Nucleare Pisa, ^{bb}University of Pisa, ^{cc}University of Siena and ^{dd}Scuola Normale Superiore, I-56127 Pisa, Italy*
- ⁴⁵*University of Pittsburgh, Pittsburgh, Pennsylvania 15260, USA*
- ⁴⁶*Purdue University, West Lafayette, Indiana 47907, USA*
- ⁴⁷*University of Rochester, Rochester, New York 14627, USA*
- ⁴⁸*The Rockefeller University, New York, New York 10065, USA*
- ⁴⁹*Istituto Nazionale di Fisica Nucleare, Sezione di Roma 1, ^{ee}Sapienza Università di Roma, I-00185 Roma, Italy*
- ⁵⁰*Rutgers University, Piscataway, New Jersey 08855, USA*
- ⁵¹*Texas A&M University, College Station, Texas 77843, USA*
- ⁵²*Istituto Nazionale di Fisica Nucleare Trieste/Udine, I-34100 Trieste, ^{ff}University of Trieste/Udine, I-33100 Udine, Italy*
- ⁵³*University of Tsukuba, Tsukuba, Ibaraki 305, Japan*
- ⁵⁴*Tufts University, Medford, Massachusetts 02155, USA*
- ⁵⁵*University of Virginia, Charlottesville, VA 22906, USA*
- ⁵⁶*Waseda University, Tokyo 169, Japan*
- ⁵⁷*Wayne State University, Detroit, Michigan 48201, USA*
- ⁵⁸*University of Wisconsin, Madison, Wisconsin 53706, USA*
- ⁵⁹*Yale University, New Haven, Connecticut 06520, USA*

We report measurements of direct CP -violating asymmetries in charmless decays of neutral bottom hadrons to pairs of charged hadrons with the upgraded Collider Detector at the Fermilab Tevatron. Using a data sample corresponding to 1 fb^{-1} of integrated luminosity, we obtain the first measurements of direct CP violation in bottom strange mesons, $A_{CP}(B_s^0 \rightarrow K^- \pi^+) = +0.39 \pm 0.15 \text{ (stat)} \pm 0.08 \text{ (syst)}$, and bottom baryons, $A_{CP}(\Lambda_b^0 \rightarrow p \pi^-) = +0.03 \pm 0.17 \text{ (stat)} \pm 0.05 \text{ (syst)}$ and $A_{CP}(\Lambda_b^0 \rightarrow p K^-) = +0.37 \pm 0.17 \text{ (stat)} \pm 0.03 \text{ (syst)}$. In addition, we measure CP violation in $B^0 \rightarrow K^+ \pi^-$ decays with 3.5σ significance, $A_{CP}(B^0 \rightarrow K^+ \pi^-) = -0.086 \pm 0.023 \text{ (stat)} \pm 0.009 \text{ (syst)}$, in agreement with the current world average. Measurements of branching fractions of $B_s^0 \rightarrow K^+ K^-$ and $B^0 \rightarrow \pi^+ \pi^-$ decays are also updated.

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Noninvariance of the fundamental interactions under

*Deceased

†With visitors from ^aUniversity of Massachusetts Amherst, Amherst, Massachusetts 01003, ^bIstituto Nazionale di Fisica Nucleare, Sezione di Cagliari, 09042 Monserrato (Cagliari), Italy, ^cUniversity of California Irvine, Irvine, CA 92697, ^dUniversity of California Santa Barbara, Santa Barbara, CA 93106 ^eUniversity of California Santa Cruz, Santa Cruz, CA 95064, ^fCERN, CH-1211 Geneva, Switzerland, ^gCornell University, Ithaca, NY 14853, ^hUniversity of Cyprus, Nicosia CY-1678, Cyprus, ⁱUniversity College Dublin, Dublin 4, Ireland, ^jUniversity of Fukui, Fukui City, Fukui Prefecture, Japan 910-0017, ^kUniversidad Iberoamericana, Mexico D.F., Mexico, ^lIowa State University, Ames, IA 50011, ^mUniversity of Iowa, Iowa City, IA 52242, ⁿKinki University,

Higashi-Osaka City, Japan 577-8502, ^oKansas State University, Manhattan, KS 66506, ^pUniversity of Manchester, Manchester M13 9PL, England, ^qQueen Mary, University of London, London, E1 4NS, England, ^rMuons, Inc., Batavia, IL 60510, ^sNagasaki Institute of Applied Science, Nagasaki, Japan, ^tNational Research Nuclear University, Moscow, Russia, ^uUniversity of Notre Dame, Notre Dame, IN 46556, ^vUniversidad de Oviedo, E-33007 Oviedo, Spain, ^wTexas Tech University, Lubbock, TX 79609, ^xUniversidad Tecnica Federico Santa Maria, 110v Valparaiso, Chile, ^yYarmouk University, Irbid 211-63, Jordan, ^{gg}On leave from J. Stefan Institute, Ljubljana, Slovenia,

the combined symmetry transformation of charge conjugation and parity inversion (CP violation) is an established experimental fact. The vast majority of experimental data are well described by the standard model (SM), and have supported the success of the Cabibbo-Kobayashi-Maskawa (CKM) [1] theory of quark-flavor dynamics. However, additional sources of CP violation are required to explain the matter–anti-matter asymmetry of the Universe in standard big-bang cosmology. This would have profound consequences on our understanding of fundamental interactions.

Violation of CP is *direct* if the partial decay-width (Γ) of a particle into a final state differs from the width of the corresponding antiparticle into the CP -conjugate final state. In recent times, the pattern of direct CP violation in charmless mesonic decays of B mesons has shown some unanticipated discrepancies from expectations. Under standard assumptions of isospin symmetry and smallness of contributions from higher-order processes, similar CP asymmetries are predicted for $B^0 \rightarrow K^+\pi^-$ and $B^+ \rightarrow K^+\pi^0$ decays [2, 3]. However, experimental data show a significant discrepancy [4], which has prompted intense experimental and theoretical activity. Several simple extensions of the standard model could accommodate the discrepancy [5], but uncertainty on the contribution of higher-order SM amplitudes has prevented a firm conclusion [6]. The violation of CP symmetry in charmless modes remains, therefore, a very interesting subject of study. Rich samples of bottom-flavored hadrons of all types from the Tevatron offer the opportunity to explore new territory in the field of B_s^0 mesons and b -flavored baryons. Additional information coming from different decays yields further constraints on the possible explanations of previous findings, and may possibly reveal new deviations from expectations.

Specifically, measurements of direct CP violation in $B_s^0 \rightarrow K^-\pi^+$ decays have been proposed as a nearly model-independent test for the presence of non-SM physics [7, 8]. The relationships between charged-current quark couplings in the SM predict a well-defined hierarchy between direct CP violation in $B^0 \rightarrow K^+\pi^-$ and $B_s^0 \rightarrow K^-\pi^+$ decays, yielding a significant asymmetry for the latter, of about 40%. This large effect allows easier experimental investigation and any discrepancy may indicate contributions from non-SM amplitudes.

Supplementary information could come from CP violation in bottom baryons, an effect which has not been measured so far. Interest in charmless b -baryon decays is prompted by branching fractions recently observed being larger than expected [9–11]. Asymmetries up to about 10% are predicted for $\Lambda_b^0 \rightarrow pK^-$ and $\Lambda_b^0 \rightarrow p\pi^-$ decays in the SM [10, 12], and are accessible with current CDF event samples.

In this letter we report the first measurement of direct CP violation in decays of bottom strange mesons and bottom baryons. We use 1 fb^{-1} of $\bar{p}p$ collisions

at $\sqrt{s} = 1.96 \text{ TeV}$, collected by the upgraded Collider Detector (CDF II) at the Tevatron. The CP -violating asymmetries are measured in the recently established [11] $B_s^0 \rightarrow K^-\pi^+$, $\Lambda_b^0 \rightarrow p\pi^-$ and $\Lambda_b^0 \rightarrow pK^-$ decays [13]. We also update our previous measurements [14] of asymmetry in the $B^0 \rightarrow K^+\pi^-$ decay, and branching fractions of $B^0 \rightarrow \pi^+\pi^-$ and $B_s^0 \rightarrow K^+K^-$ decays.

The CDF II detector is described in detail in Ref. [15] with the detector sub-systems relevant for this analysis discussed in Ref. [14]. The data are collected by a three-level trigger system. At level 1, tracks are reconstructed in the transverse plane. Two opposite-charge particles are required, with reconstructed transverse momenta $p_{T1}, p_{T2} > 2 \text{ GeV}/c$, the scalar sum $p_{T1} + p_{T2} > 5.5 \text{ GeV}/c$, and an azimuthal opening-angle $\Delta\phi < 135^\circ$. At level 2, tracks are combined with silicon hits and their impact parameter d (transverse distance of closest approach to the beam line) is determined with $45 \mu\text{m}$ resolution (including the beam spread) and required to be $0.1 < d < 1.0 \text{ mm}$. A tighter opening-angle requirement, $20^\circ < \Delta\phi < 135^\circ$, is also applied. Each track pair is then used to form a B candidate, which is required to have an impact parameter $d_B < 140 \mu\text{m}$ and to have travelled a distance $L_T > 200 \mu\text{m}$ in the transverse plane. At level 3, a cluster of computers confirms the selection with a full event reconstruction.

The offline selection is based on a more accurate determination of the same quantities used in the trigger, with the addition of requirements on two other observables: the isolation (I_B) of the B candidate [17], and the quality of the three-dimensional fit (χ^2 with 1 d.o.f.) of the decay vertex of the B candidate [11]. Asymmetries in the rarer $B_s^0 \rightarrow K^-\pi^+$ and Λ_b^0 decays are measured using the selection in Ref. [11]. For the measurement of the $B^0 \rightarrow K^+\pi^-$ asymmetry, instead, the selection is optimized by minimizing the expected variance of the measurement, evaluated by performing the full analysis on a set of simulated samples obtained with varied selection criteria [18]. This procedure yields the final selection: $I_B > 0.5$, $\chi^2 < 7$, $d > 100 \mu\text{m}$, $d_B < 80 \mu\text{m}$, and $L_T > 300 \mu\text{m}$. Only one B candidate per event is found after this selection, and a mass ($m_{\pi\pi}$) is assigned to each, using a nominal charged-pion mass assignment for both decay products. The resulting mass distribution is shown in Fig. 1. A large peak is visible, dominated by the overlapping contributions of the $B^0 \rightarrow K^+\pi^-$, $B^0 \rightarrow \pi^+\pi^-$, and $B_s^0 \rightarrow K^+K^-$ decays [14]. Signals for $B_s^0 \rightarrow K^-\pi^+$, $\Lambda_b^0 \rightarrow p\pi^-$, and $\Lambda_b^0 \rightarrow pK^-$ modes populate masses higher than the main peak (5.33 – $5.55 \text{ GeV}/c^2$) [11]. Backgrounds include mis-reconstructed multi-body b -hadron decays (physics background) and random pairs of particles (combinatorial background).

We incorporate kinematic and particle identification information in an unbinned likelihood fit [11, 14] to determine the fraction of each mode and the charge asymmetries, uncorrected for instrumental effects, $\tilde{A}_{CP} =$

$[N_{b \rightarrow f} - N_{\bar{b} \rightarrow \bar{f}}]/[N_{b \rightarrow f} + N_{\bar{b} \rightarrow \bar{f}}]$ of the flavor-specific decays $B^0 \rightarrow K^+\pi^-$, $B_s^0 \rightarrow K^-\pi^+$, and $\Lambda_b^0 \rightarrow p\pi^-$, pK^- . For each channel, $N_{b \rightarrow f}$ ($N_{\bar{b} \rightarrow \bar{f}}$) is the reconstructed number of decays of hadrons containing the b (\bar{b}) quark into the final state f (\bar{f}). The decay flavor is inferred from the charges of final state particles assuming equal numbers of b and \bar{b} quarks at production (dominated by the strong interaction). Any effect from CP violation in b -meson flavor mixing is assumed negligible [19].

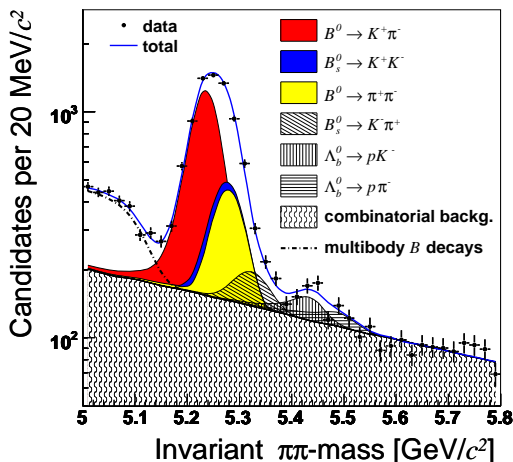


FIG. 1: Mass distribution of the 13 502 reconstructed candidates. The charged-pion mass is assigned to both tracks. The total projection and projections of each signal and background component of the likelihood fit are overlaid on the data distribution. Signals and multi-body B backgrounds are shown stacked on the combinatorial background component.

The whole kinematic information is summarized by three loosely correlated observables [11]: the mass $m_{\pi\pi}$; the signed momentum imbalance $\alpha = (1 - p_1/p_2) \times q_1$, where p_1 (p_2) is the lower (higher) of the particle momenta, and q_1 is the sign of the charge of the particle of momentum p_1 ; and the scalar sum of particle momenta $p_{tot} = p_1 + p_2$. Particle identification relies on measurement of the specific ionization (dE/dx) in the drift chamber. For charged kaons and pions the dE/dx response was calibrated with a sample of 1.5×10^6 $D^{*+} \rightarrow D^0\pi^+$ decays, using the charge of the pion from D^{*+} decay to identify the products of the Cabibbo-favored D^0 decay. For protons we used 124 000 $\Lambda \rightarrow p\pi^-$ decays, where the kinematics and the momentum threshold of the trigger allow unambiguous identification of the decay products [18, 20]. Identification information for each particle is summarized by a single observable in our fit (“kaonness”), defined as $\kappa = (dE/dx - dE/dx_\pi)/(dE/dx_K - dE/dx_\pi)$, where dE/dx is the observed response, and $dE/dx_{\pi(K)}$ is the average responses expected for pions (kaons). The separation between $K^+\pi^-$ or $p\pi^-$ final states and their charge-conjugates is in excess of 2.1σ (Fig. 2). Although a lower dE/dx separation is avail-

able between pK^- and $\bar{p}K^+$, due to similar ionization rates of protons and kaons, sufficient discrimination is achieved from their greater kinematics differences. The background model allows for independent contributions of positively and negatively charged pions, kaons, protons, and electrons, whose fractions are determined by the fit. Muons are indistinguishable from pions with the available 10% fractional dE/dx resolution and are therefore incorporated into the pion component.

The signal yields from the fit (Table I) are corrected for different detection efficiencies to determine the physical asymmetries, $A_{CP}(b \rightarrow f)$, defined as

$$\frac{\mathcal{B}(b \rightarrow f) - \mathcal{B}(\bar{b} \rightarrow \bar{f})}{\mathcal{B}(b \rightarrow f) + \mathcal{B}(\bar{b} \rightarrow \bar{f})} = \frac{N_{b \rightarrow f} - c_f N_{\bar{b} \rightarrow \bar{f}}}{N_{b \rightarrow f} + c_f N_{\bar{b} \rightarrow \bar{f}}}, \quad (1)$$

where $c_f = \varepsilon(f)/\varepsilon(\bar{f})$ is the ratio between the efficiencies for triggering and reconstructing the final state f with respect to the state \bar{f} . The c_f factors correct for detector-induced charge asymmetries, and are extracted from control samples in data. Simulation is only used to account for small differences between the kinematics of $B \rightarrow h^+h'^-$ decays and control signals. The corrections for $f = K^+\pi^-$ are extracted from a sample of about 700 000 $D^0 \rightarrow K^-\pi^+$ decays, reconstructed in the same data set. By imposing the same offline selection to the D^0 decays we obtain $K^\mp\pi^\pm$ final states in a similar kinematic region as our signals (see Fig. 2). We assume that $K^+\pi^-$ and $K^-\pi^+$ final states from charm decays are produced in equal numbers at the Tevatron, because production is dominated by the strong interaction and, compared to the detector effects to be corrected, the possible CP -violating asymmetry in $D^0 \rightarrow K^-\pi^+$ decays is tiny ($< 10^{-3}$) as predicted by the SM [21] and confirmed by current experimental determinations [22]. We also checked that possible asymmetries in D^0 meson yields induced by CP violation in $B \rightarrow DX$ decays are small and can be neglected [18]. Therefore, any asymmetry between observed numbers of reconstructed $K^-\pi^+$ and $K^+\pi^-$ charm decays can be ascribed to detector-induced effects and used to extract the desired correction factors. The ratio $N_{D^0 \rightarrow K^+\pi^-}/N_{D^0 \rightarrow K^-\pi^+}$ is measured with the same fit used for the signal. The dE/dx information is not used because kinematics alone is sufficient to provide an excellent separation in charm decays, as shown in Fig. 2. We checked separately that dE/dx information does not introduce additional charge asymmetries [18]. We find $c_{K^-\pi^+} = 0.9871 \pm 0.0027$, which is consistent and more precise than a previous estimate based on simulation [23]. For the $\Lambda_b^0 \rightarrow p\pi^-$ asymmetry, the factor $c_{p\pi^-} = 1.0145 \pm 0.0075$ is extracted using a similar strategy applied to a control sample of $\Lambda \rightarrow p\pi^-$ decays [20]. This factor is dominated by the different interaction probability of protons and antiprotons with detector material. In the measurement of CP violation in $\Lambda_b^0 \rightarrow pK^-$ decays, instrumental charge-asymmetries

TABLE I: Raw signal yields determined by the fit and final results. The first uncertainty is statistical, the second is systematic. Absolute branching fractions are derived by normalizing to the known value $\mathcal{B}(B^0 \rightarrow K^+\pi^-) = (19.4 \pm 0.6) \times 10^{-6}$, and assuming the average value at high energy for the production fraction $f_s/f_d = 0.282 \pm 0.038$ [19].

Mode	$N_{b \rightarrow f}$	$N_{\bar{b} \rightarrow \bar{f}}$	$A_{CP}(b \rightarrow f)$ (%)	Relative \mathcal{B}	Absolute $\mathcal{B}(10^{-6})$
$B^0 \rightarrow K^+\pi^-$	1836 ± 61	2209 ± 64	$-8.6 \pm 2.3 \pm 0.9$	–	–
$B_s^0 \rightarrow K^-\pi^+$	160 ± 26	70 ± 22	$+39 \pm 15 \pm 8$	–	–
$\Lambda_b^0 \rightarrow pK^-$	80 ± 14	36 ± 11	$+37 \pm 17 \pm 3$	–	–
$\Lambda_b^0 \rightarrow p\pi^-$	40 ± 10	38 ± 9	$+3 \pm 17 \pm 5$	–	–
$B^0 \rightarrow \pi^+\pi^-$	1121 ± 63	–	–	$\frac{\mathcal{B}(B^0 \rightarrow \pi^+\pi^-)}{\mathcal{B}(B^0 \rightarrow K^+\pi^-)} = 0.259 \pm 0.017 \pm 0.016$	$5.02 \pm 0.33 \pm 0.35$
$B_s^0 \rightarrow K^+K^-$	1307 ± 64	–	–	$\frac{f_s}{f_d} \frac{\mathcal{B}(B_s^0 \rightarrow K^+K^-)}{\mathcal{B}(B^0 \rightarrow K^+\pi^-)} = 0.347 \pm 0.020 \pm 0.021$	$23.9 \pm 1.4 \pm 3.6$

induced in both kaons and protons are relevant. The c_{pK^-} factor is extracted by combining the previous ones and assuming the trigger and reconstruction efficiency for two particles factorizes as the product of the single-particle efficiencies. Corrections are also applied for the

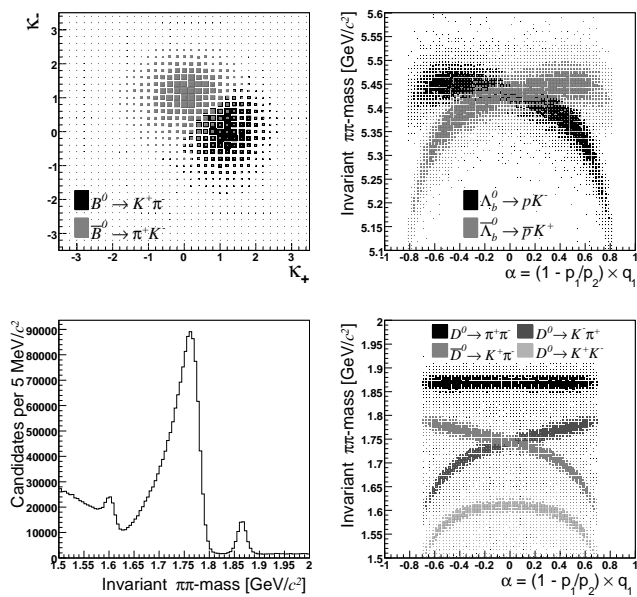


FIG. 2: Joint kaonness distribution for the positive (abscissa) and negative (ordinate) final state particles in $B^0 \rightarrow K^+\pi^-$ decays as determined from the calibration data of charm decays (top left). Dipion mass as a function of α for simulated $\Lambda_b^0 \rightarrow pK^-$ decays (top right). Mass of $D^0 \rightarrow h^+h'^-$ candidates with pion assignment to both final state particles (bottom left). Same quantity as a function of α for simulated $D^0 \rightarrow h^+h'^-$ decays (bottom right)

branching ratio measurements. These corrections do not exceed 7% and account for differences in trigger and reconstruction efficiency between channels due to different lifetimes and kinematics (from simulation), and isolation properties (from control samples of fully reconstructed $B^0 \rightarrow J/\psi K^*(892)^0$ and $B_s^0 \rightarrow J/\psi\phi$ decays).

The dominant contributions to the systematic uncertainties on the asymmetry measurements come from the uncertainty on the dE/dx calibration and parameteriza-

tion, the uncertainty on the combinatorial background model, and the uncertainty on b -hadron masses. Smaller contributions are assigned for the uncertainty on the global mass scale and the c_f corrections. The uncertainty on the dE/dx model dominates also the systematic uncertainty for the branching ratio measurements, for which the mass scale (in the $B^0 \rightarrow \pi^+\pi^-$ case) and the uncertainty on the difference in isolation efficiency between B^0 and B_s^0 mesons ($B_s^0 \rightarrow K^+K^-$) also play a role. The results are reported in Table I. We report 3.5σ evidence of CP violation in $B^0 \rightarrow K^+\pi^-$ decays. The observed asymmetry is consistent, and of comparable accuracy, with current results from asymmetric e^+e^- colliders [4]. It is also consistent with the result in Ref. [14] and supersedes it. The $B_s^0 \rightarrow K^-\pi^+$ result is the first measurement of direct CP violation in bottom strange mesons. It differs by 2.3σ from zero and it is consistent with recent theoretical predictions [3, 24]. It allows the first experimental verification of the model-independent test proposed in Ref. [8]. Under the assumption of equal B^0 and B_s^0 lifetimes, using the measurement of the $B_s^0 \rightarrow K^-\pi^+$ branching ratio [11] and known values for the branching ratio and CP -violating asymmetry in $B^0 \rightarrow K^+\pi^-$ decays, and the b -quark fragmentation probabilities [19], we obtain $R = [\Gamma(B^0 \rightarrow K^+\pi^-) - \Gamma(\bar{B}^0 \rightarrow K^-\pi^+)] / [\Gamma(\bar{B}^0 \rightarrow K^+\pi^-) - \Gamma(B_s^0 \rightarrow K^-\pi^+)] = 0.85 \pm 0.42$ (stat) ± 0.13 (syst), which is consistent with the standard prediction, $R^{\text{SM}} = 1$ [8]. The first measurement of CP violation in bottom baryons is also reported. The observed asymmetry in the $\Lambda_b^0 \rightarrow pK^-$ decay is 2.1σ from zero. The $\Lambda_b^0 \rightarrow p\pi^-$ result is consistent with zero. However, the limited experimental precision does not allow a conclusive discrimination between the standard model prediction (8%) and much suppressed values ($\approx 0.3\%$) expected in R -parity violating supersymmetric scenarios [12].

Table I includes also improved measurements of $B_s^0 \rightarrow K^+K^-$ and $B^0 \rightarrow \pi^+\pi^-$ CP -averaged branching fractions, using the $B^0 \rightarrow K^+\pi^-$ channel as a reference. Results are consistent with previous CDF measurements [14] and supersede them. The $B_s^0 \rightarrow K^+K^-$ result is the most precise to date and consistent with recent theoretical predictions [3, 24–26]. Theory uncertain-

ties, which are significantly larger than the experimental ones, prevent sensible discrimination between models. The present measurement of $\mathcal{B}(B^0 \rightarrow \pi^+\pi^-)$ agrees with measurements at e^+e^- colliders [27] with comparable accuracy. The dominant systematic uncertainties are limited by the finite size of control samples and should decrease in future extensions of the measurements.

In conclusion, we have measured CP -violating asymmetries in charmless B^0 , B_s^0 , and Λ_b^0 decays into pairs of charged hadrons reconstructed in CDF data. We report the first measurement of direct CP violation in bottom strange mesons, the first measurement of CP violation in bottom baryons, evidence for CP violation in $B^0 \rightarrow K^+\pi^-$ decays, and updated measurements of the $B_s^0 \rightarrow K^+K^-$ and $B^0 \rightarrow \pi^+\pi^-$ branching fractions.

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